Jets, tidal disruption and lenses

Plan and reviews

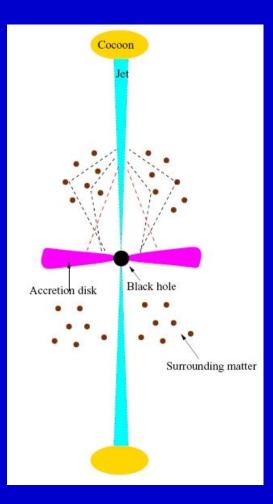
<u>Plan</u>

- 1. Jets: AGNs and close binary systems
- 2. Tidal distruption of stars by SMBHs
- 3. Spectral lines and lensing

<u>Reviews</u>

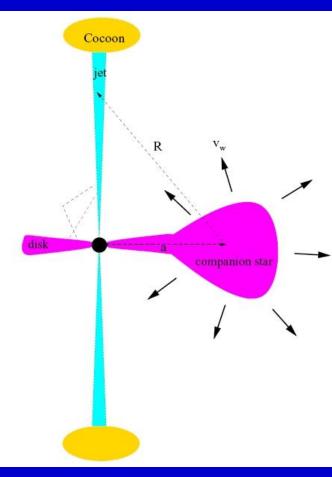
- astro-ph/0611521 High-Energy Aspects of Astrophysical Jets
- astro-ph/0306429 Extreme blazars
- astro-ph/0312545 AGN Unification: An Update
- astro-ph/0212065 Fluorescent iron lines as a probe of astrophysical black hole systems
- arXiv: 1104.0006 AGN jets
- astro-ph/0406319 Astrophysical Jets and Outflows
- arXiv: 1309.4772 Blazars: jets and discs
- arXiv: 1707.07134 AGNs of different types

Jets in AGNs and close binaries



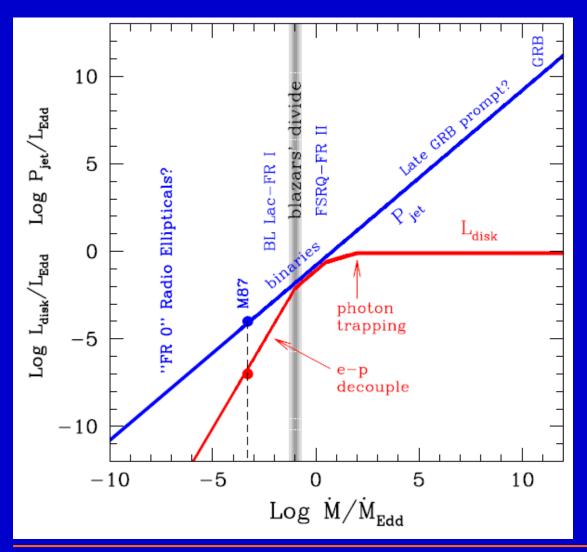
AGN: M_{BH} =10⁸-10⁹ M_0 L<~ L_{Edd} ~10⁴²-10⁴⁷ erg/s < few Mpc Γ ~5-50 Δt ~ hours-years

CBS: $M_{BH} \sim 10 M_0$ $L < \sim L_{Edd} \sim 10^{37} \cdot 10^{40} \text{ erg/s}$ $\sim pc$ $\Gamma \sim 1 \cdot 10$ $\Delta t \sim days$



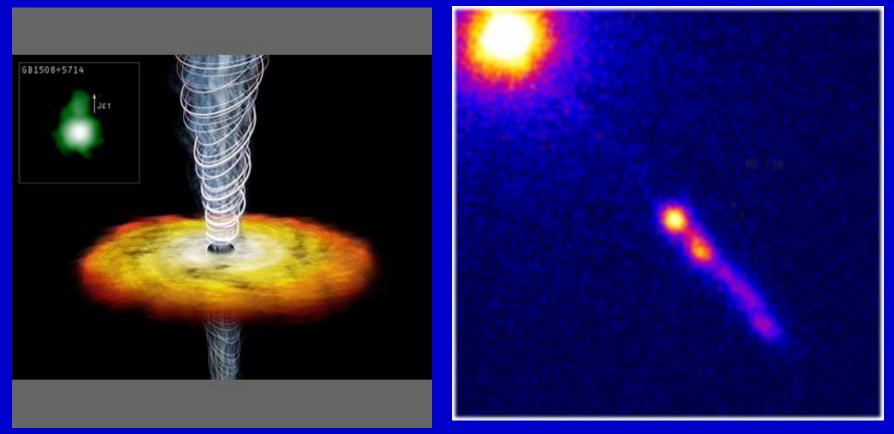
see astro-ph/0611521

All jets in one plot



Close-by and far-away jets

1% of SMBH are active. 10% out of them launch relativistic jets. Jets are not magnetically dominated.

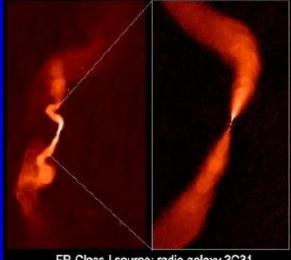


GB1508+5714 z=4.30

3C273

See a review in 1104.0006

Classification of AGN radio jets



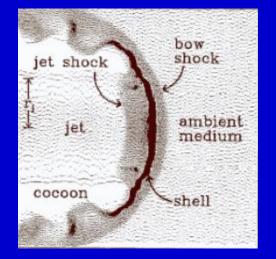
FR Class I source: radio galaxy 3C31



FR Class II source: quasar 3C175

astro-ph/0406319

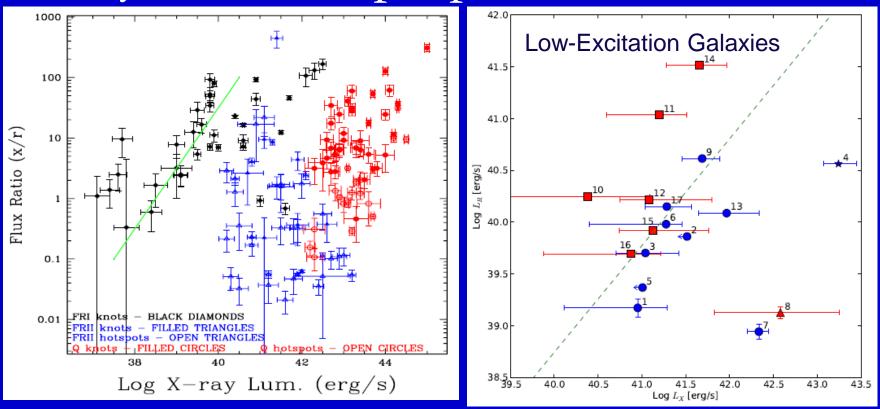
FR I. Two-sided jets. Jets dominate in the emission. Usually are found in rich clusters.



FR II. One-sided jets. Radio lobes dominate over jets. Mostly isolated galaxies or poor groups.

See a review on radio galaxies in arXiv: 1101.0837

X-ray and radio properties



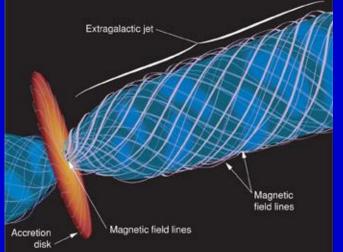
1003.0976

Magnetic field in a jet



Observations of M87 tell us that the magnetic field in the jet is mostly parallel to the jet axis, but in the emission regions ("knots") it becomes perpendicular (see astro-ph/0406319).

The same structure is observed in several jets with radio lobes.

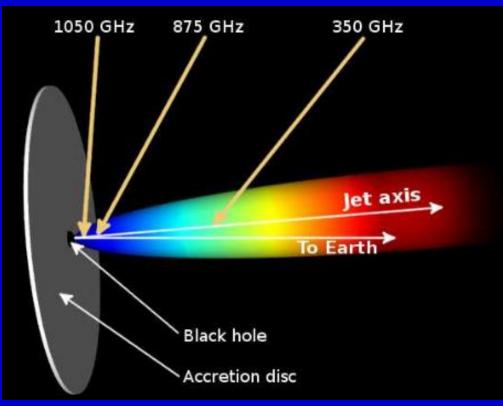


New RadioAstron data give some new insights.

Magnetic field in the jet

Due to modern high resolution observations new important results on the magnetic field in jets are obtained.

ALMA 0.01 pc scale

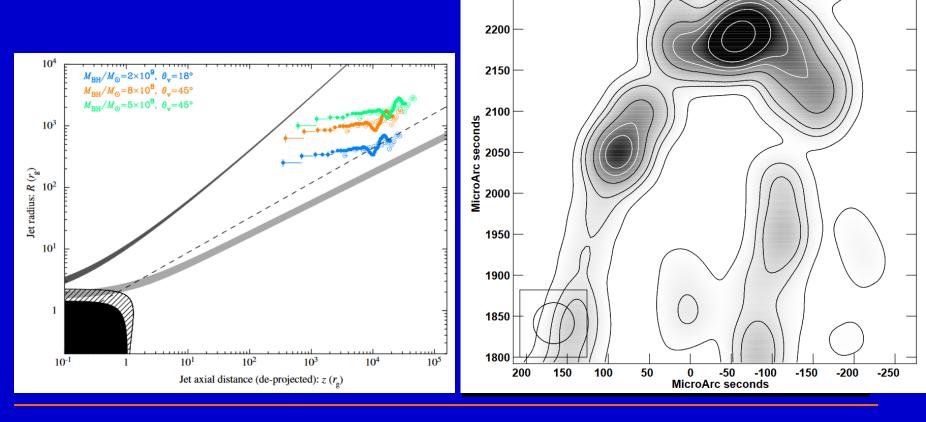


Magnetic fields of at least tens of Gauss (and possibly considerably higher) on scales of the order of light days (0.01 pc) from the black hole.

3C84 jet

Jet followed down to >~100 R_{sh} BH mass ~2 10⁹

Jet from the disc, not from ergosphere?



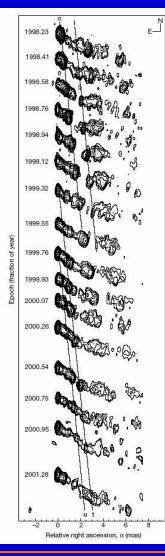
2.5

2.0

2250

Core

Blobs in jets





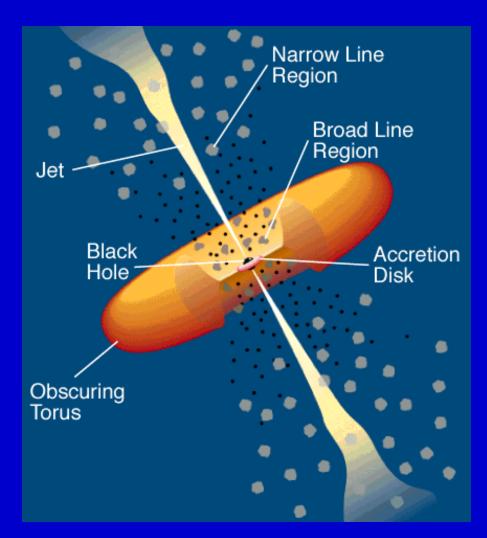
It is believed that bright features in AGN jets can be results of the Kelvin-Helmholtz instability. This instability leads to a spiral structure formation in a jet. (see, for example, astro-ph/0103379)

3C 120

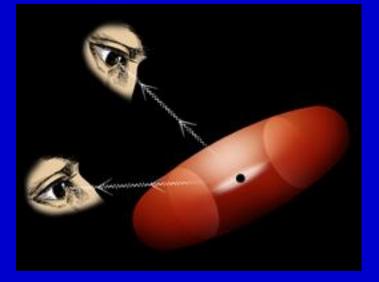
However, in the case of 3C 120 the blobs appearence is due to processes in the disc. Dips in X-rays (related to the disc) appear before blobs ejection (Marscher et al. 2002).

Marscher, A.P., et al., NATURE Vol 417 p. 625

Blazars



If a jet is pointing towards us, then we see a *blaza*r.



A review on blazar jets 1010.2856

Blazars at very high energies

Blazars are powerful gamma-ray sources. The most powerful of them have equivalent isotropic luminosity 10⁴⁹ erg/s.

Collimation $\theta^2/2 \sim 10^{-2} - 10^{-3}$. θ – jet opening angle. EGRET detected 66 (+27) sources of this type. New breakthrough is expected after the launch of GLAST.

Several sources have been detected in the TeV range by ground-based gamma-ray telescopes. All of them, except M87, are BL Lacs at z<0.2 (more precisely, to high-frequency-peaked BL Lac – HBL).

Observations show that often (but not always) after a gamma-ray bursts few weeks or months later a burst happens also in the radio band.

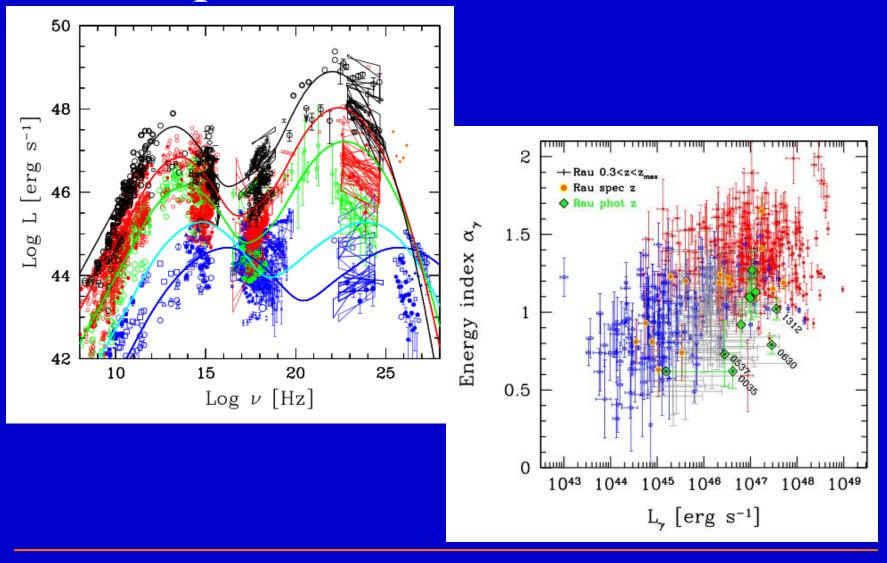
(see astro-ph/0611521)



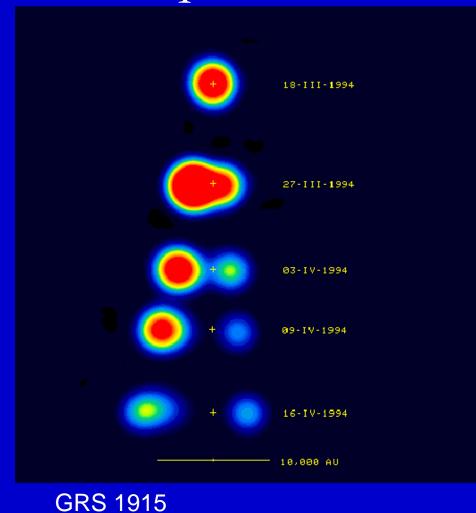
Fermi

1309.4772

Blazar spectra

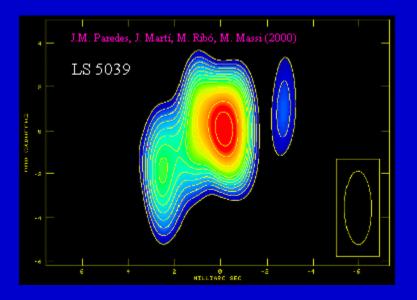


Microquasars



The correlation between X-ray and synchrotron (i.e. between disc and jet emission) is observed.

Microquasars jets in radio



LS 5039/RX J1826.2-1450 – is a galactic massive X-ray binary.

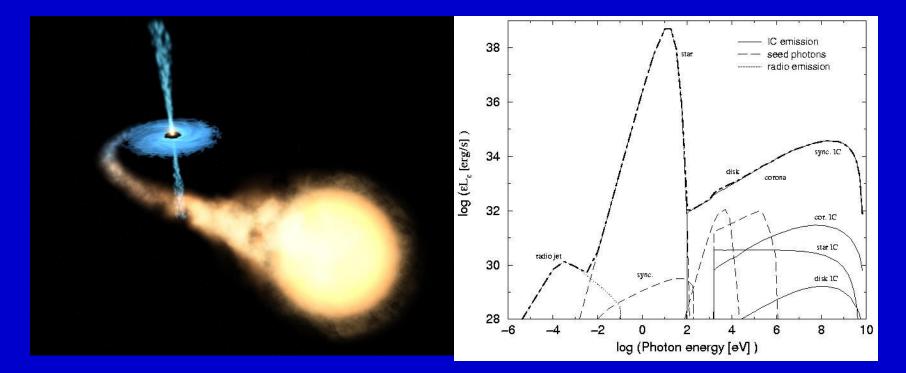
The jet length is ~1000 a.e.

Probably, the source was observed by EGRET as 3EG J1824-1514.

(Many examples of VLBI radio jets from different sources can be found at the web-site http://www.evlbi.org/gallery/images.html)

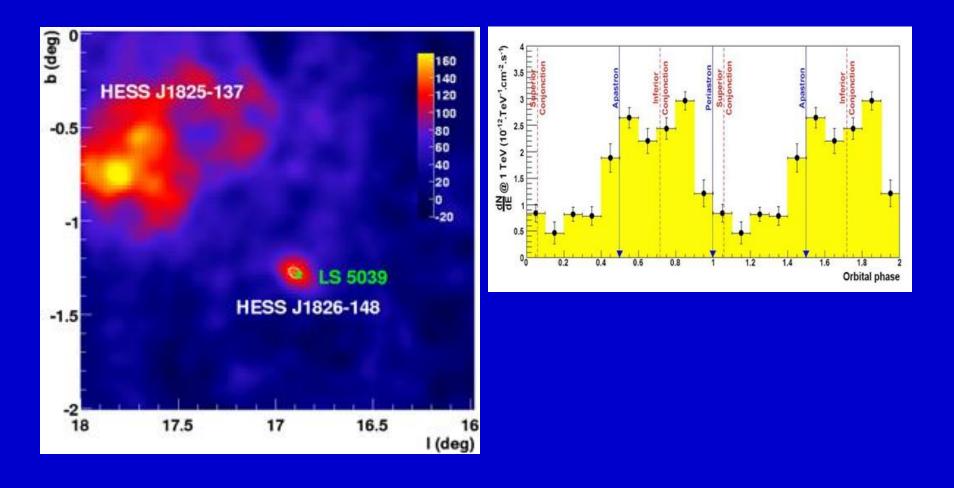
The role of a donor

An important difference between the microquasars case and AGNs is related to the existence of a donor-star. Especially, if it is a giant, then the star can inject matter and photons into the jet.



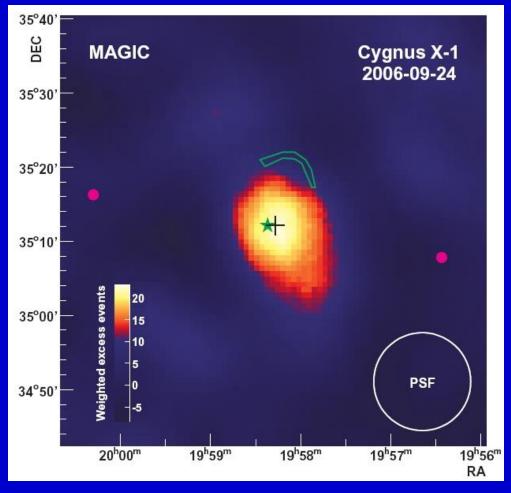
see Paredes astro-ph/0412057

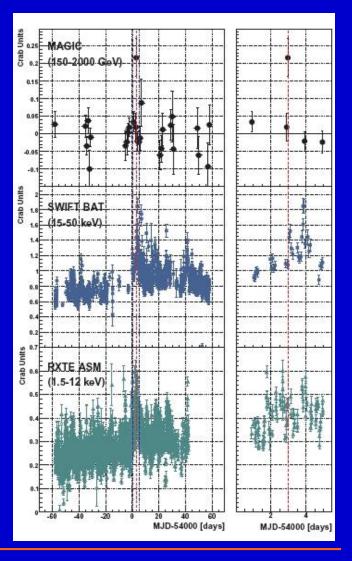
Microquasars in gamma-rays: TeV range



F. Aharonian et al.

TeV emission from Cyg X-1

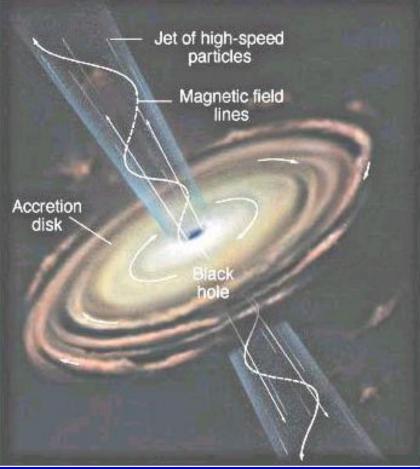




arxiv:0706.1505

See a review on jets in binaries in 1407.3674

Jet models

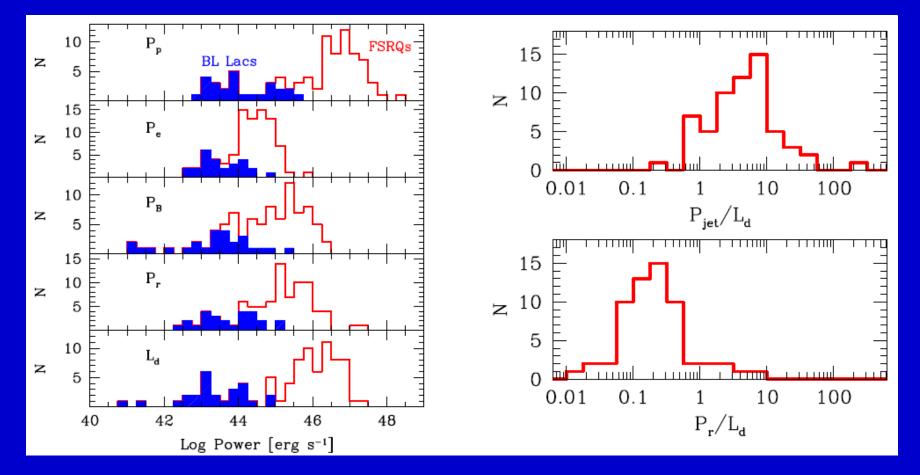


(the table is from astro-ph/0611521)

Table 1. sources.	Characteristic parameters of relativistic j			
	L_j (erg/s)	Г	Δt	
GRB	10^{47} - 10^{50}	$10^2 - 10^3$	millisec - min	
AGN	$10^{42} - 10^{47}$	5 - 50	hours - years	
MQ	$10^{37} - 10^{40}$	1 - 10	days	
\mathbf{GF}	$10^{43} - 10^{46}$	1	seconds	

In all models jets are related to discs. Velocity at the base of a jet is about the parabolic (escape) velocity.

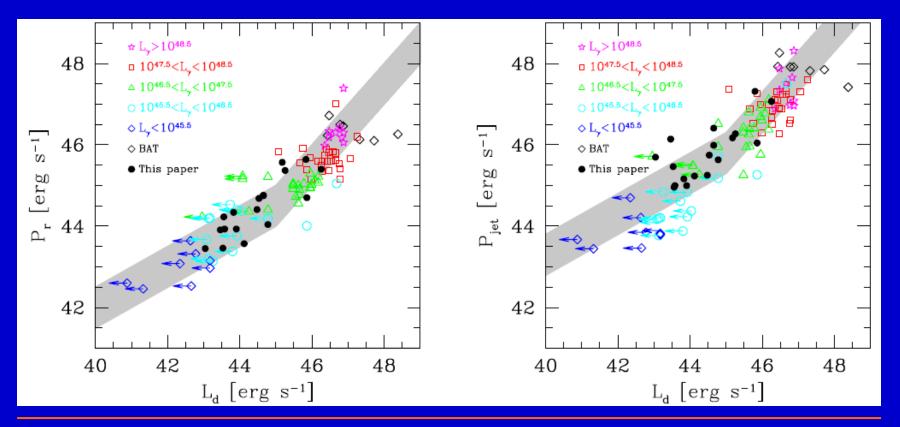
Jet power



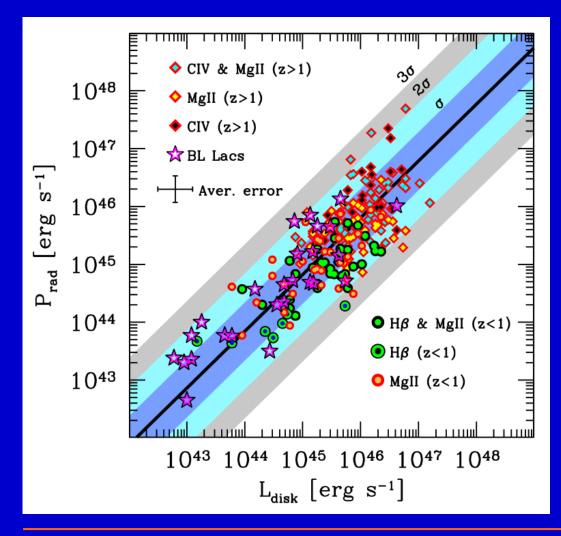
 $P_{jet} > 2P_r \sim 2L/\Gamma^2$

Jet and disc

 $\begin{array}{l} \mathsf{L}_{\mathsf{d}} - \mathsf{disc \, luminosity.} \\ \mathsf{P}_{\mathsf{jet}} = \mathsf{P}_{\mathsf{B}} + \mathsf{P}_{\mathsf{P}} + \mathsf{P}_{\mathsf{e}} \\ \mathsf{In \, the \, gray \, stripes \, P_{\mathsf{jet}}} \sim \mathsf{Mdot \, and} \\ \mathsf{at \, low \, accretion \, rates \, \mathsf{L}_{\mathsf{d}}} \sim \mathsf{Mdot}^2, \, \mathsf{at \, large} \, \text{-} \, \mathsf{L}_{\mathsf{d}} \sim \mathsf{Mdot}. \end{array}$

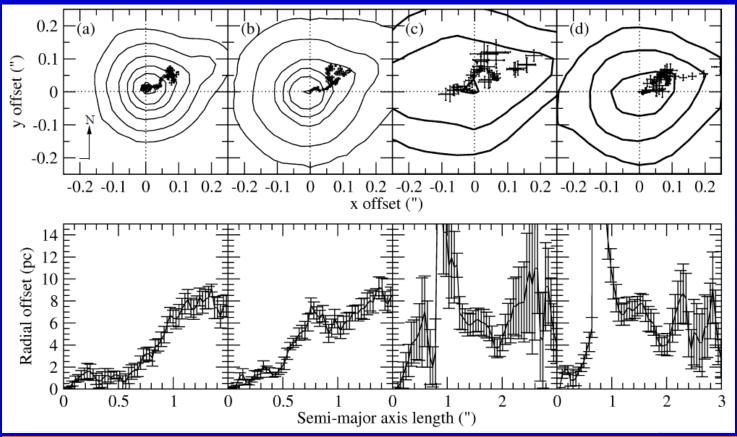


Jets are more powerful than discs



Displaced SMBH in M87

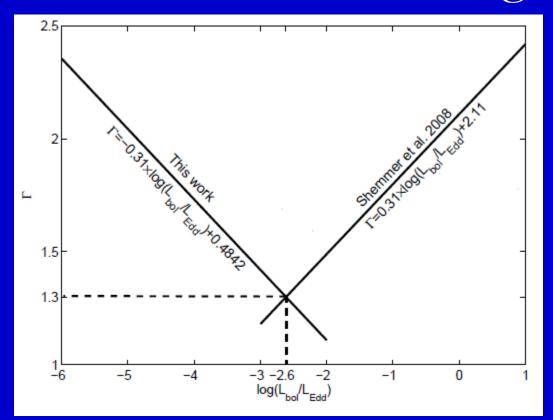
Projected displacement of 6.8+/-0.8 pc consistent with the jet axis displaced in the counter-jet direction



1005.2173

Other explanations are possible

Different accretion regimes in AGNs



Anticorrelation for low-luminosity AGNs (LINERS).

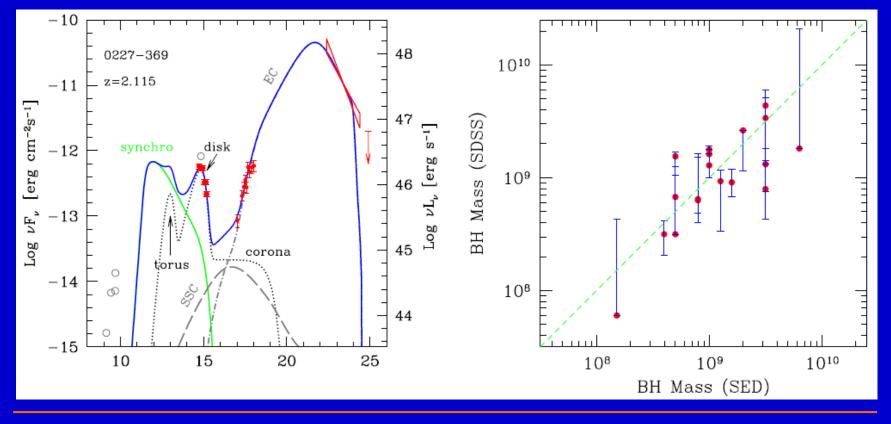
Correlation for luminous AGNs.

In the critical point the accretion regime can be changed: from a standard thin accretion disc to RIAF (radiatively inefficient accretion flow).

Γ- photon index

BH mass determination

Accretion disc contribution is visible in opt-UV. It allows to estimate the BH mass. It can be compared with emission lines estimates. Bars correspond to different lines used.



BH mass and jet properties in M87

Table 3. Predicted jet and BH parameters for $M = 6.5 \times 10^9 M_{\odot}$.

$\sigma_{ m M}$	$\Gamma heta_j$	ϕ	a_*	$\frac{W_j}{(10^{42} \text{ erg/s})}$
(1)	(2)	(3)	(4)	(5)
10	0.062	3.2	0.093	1.9
20	0.063	3.3	0.144	4.7

Estimates are in correspondence with the EHT data on the BH mass.

If the mass is assumed to be known, then the initial magnetization σ_M can be determined.

Tidal disruption

The Hills limit: 3 10⁸ solar masses. A BH disrupts stars.

After a disruption in $(t_0 - t_D) \sim 1.1 M_8^{1/2} \text{ yr}$ happens a burst with
the temperature $T_{\text{eff}} \approx (L_{\text{Edd}}/4\pi\sigma R_T^2)^{1/4} = 3.7 \times 10^5 M_8^{1/12} \text{ K}$ The maximum accretion rate $\dot{M}_{\text{max}} \sim 0.14 M_8^{1/2} M_{\odot} \text{ yr}^{-1}$ This rate corresponds to the moment $(t_{\text{max}} - t_D) \sim 1.5 (t_0 - t_D)$ Then the rate can be described as $\dot{M}(t) = 0.3 M_8 [(t - t_D)/(t_0 - t_D)]^{-5/3} M_{\odot}$

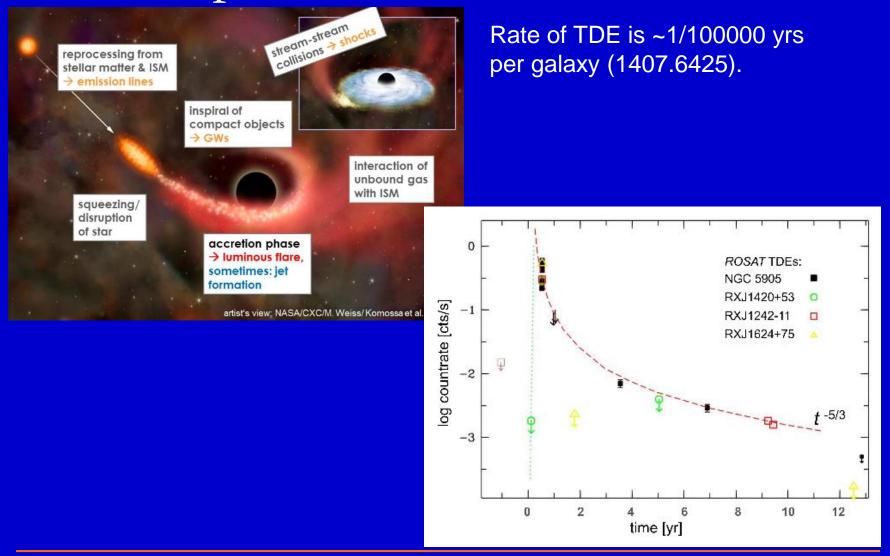
For a BH with M $< 10^7 M_0$ the luminosity at maximum is:

$$L_{\rm flare} \geq \eta \dot{M}_{\rm Edd} c^2 > 1.3 \times 10^{45} M_7 \ {\rm ergs} \ {\rm s}^{-1}$$

(astro-ph/0402497)

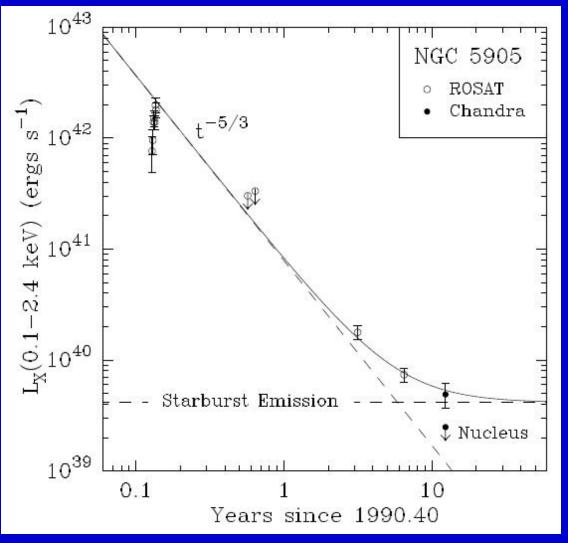
See a review in 1505.01093

General picture



See a review in 1505.01093

A burst in NGC 5905



The decay was well described by the relation:

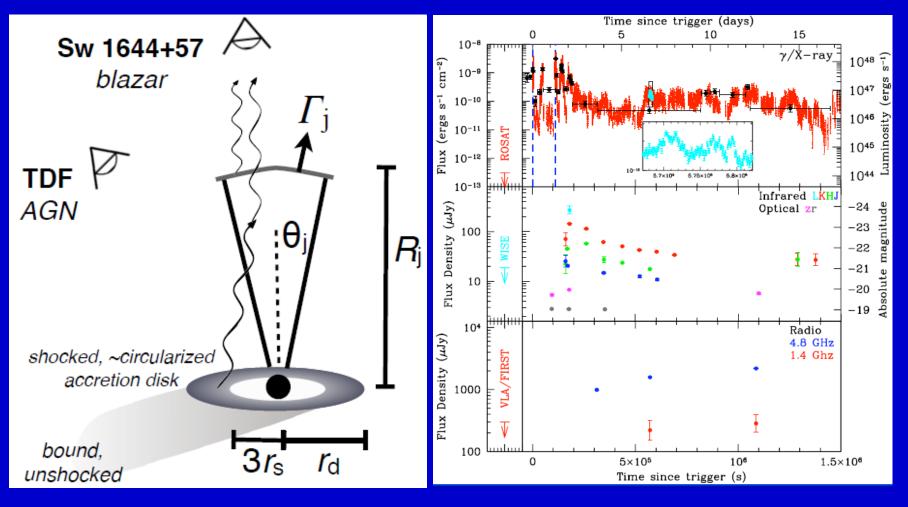
$$(t-t_D)^{-5/3}$$

Two other bursts discovered by ROSAT and observed by HST and Chandra:

RX J1624.9+7554 RX J1242.6-1119A

astro-ph/0402497

High-energy transient Swift J164449.3+573451

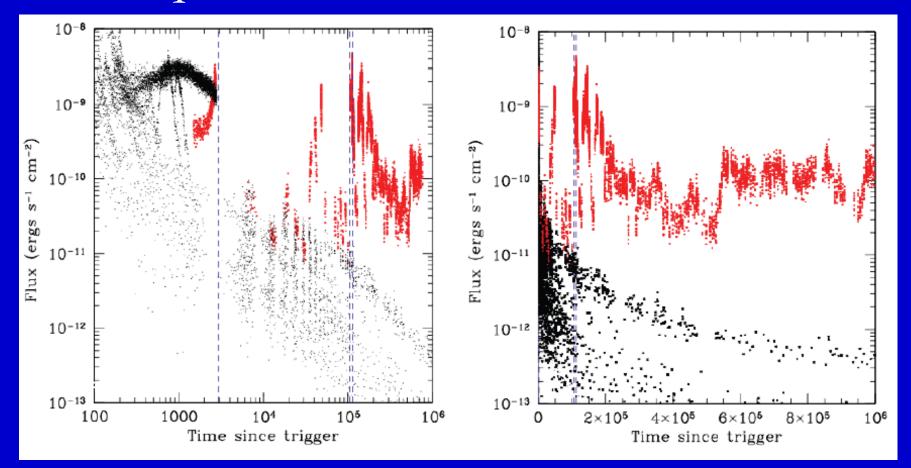


1104.3257

1104.3356

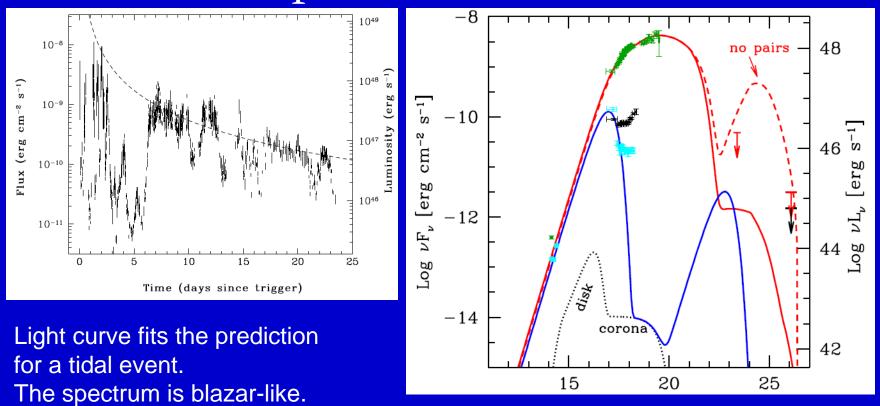
Also known as GRB 110328A

A unique event

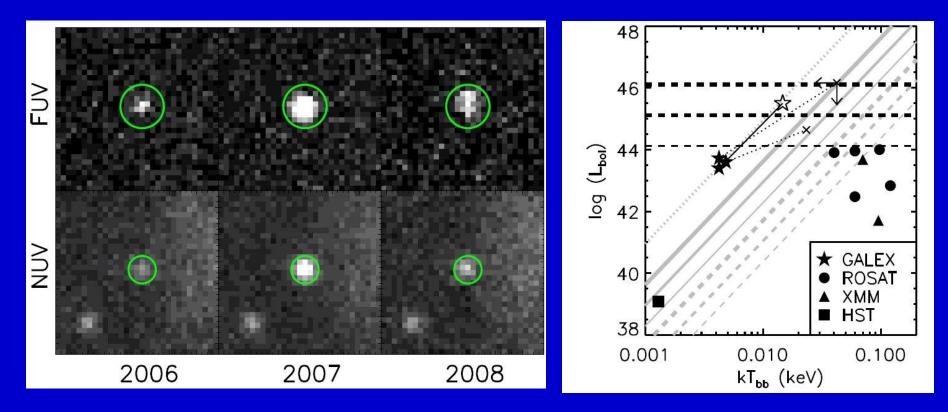


The transient does not looks similar to GRBs, SN or any other type of event

A tidal disruption event



Optical observations of tidal disruptions



In optics one can observe events far from the horizon. Future surveys (like Pan-STARRS) can discover 20-30 events per year.

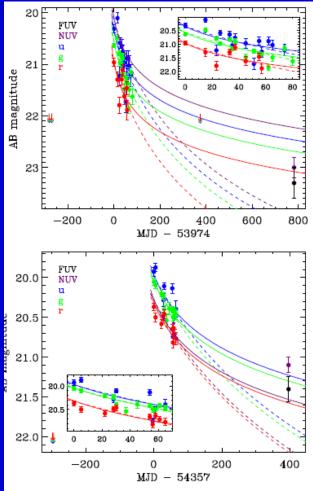
arxiv:0904.1596

Two more examples of optical flares due to tidal disruption events

SDSS data

Atypical flares

Rate estimates: ~10⁻⁵ per year per galaxy or slightly more



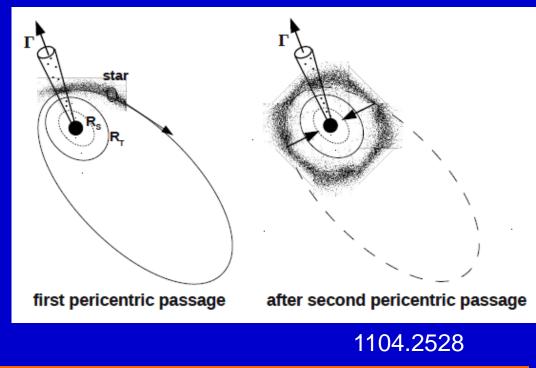
1009.1627

Dashed lines: -5/3 Solid lines: -5/9

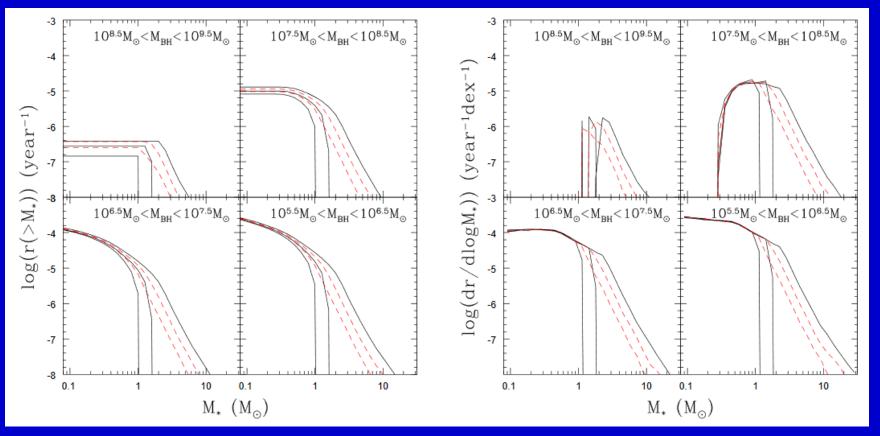
Theoretical models

- X-rays. 1105.2060
- Radio. 1104.4105

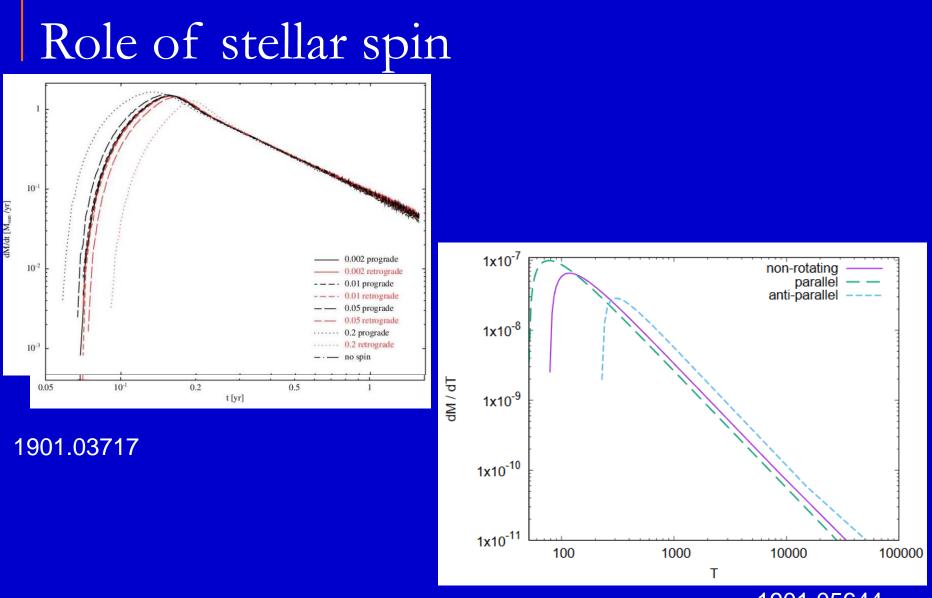
Flows of hot X-ray emitting gas have been identified after one of tidal disruption events. 1510.06348



Stellar masses



For not too massive BHs the main contribution is due to M-dwarfs (0.3Msolar). For massive BHs contribution from massive and evolved stars grows.



TDE and binary SMBHs

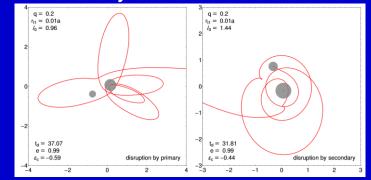
It has been predicted that after a TDE in a system of close SMBH binary

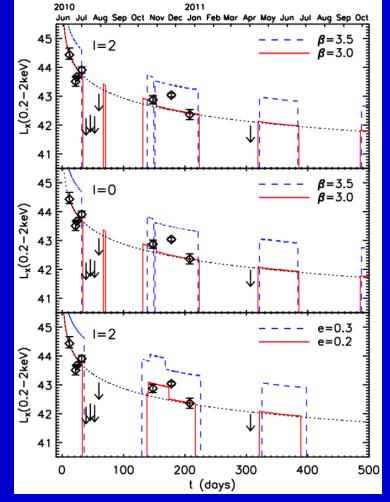
there might be particular drops in the light curve. Such phenomena was observed.

SDSS J120136.02+300305.5 XMM-Newton observations

Masses $\sim 10^7$ and 10^6 solar masses, orbital separation $\sim 10^{-3}$ pc Eccentricity ~ 0.3

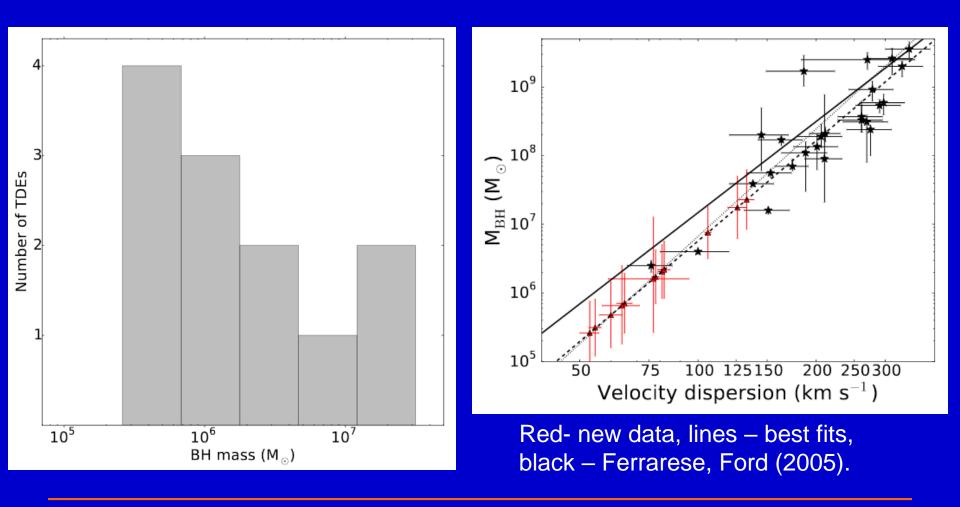
TDEs in supermassive BH binaries were recently modeled in 1802.07850.



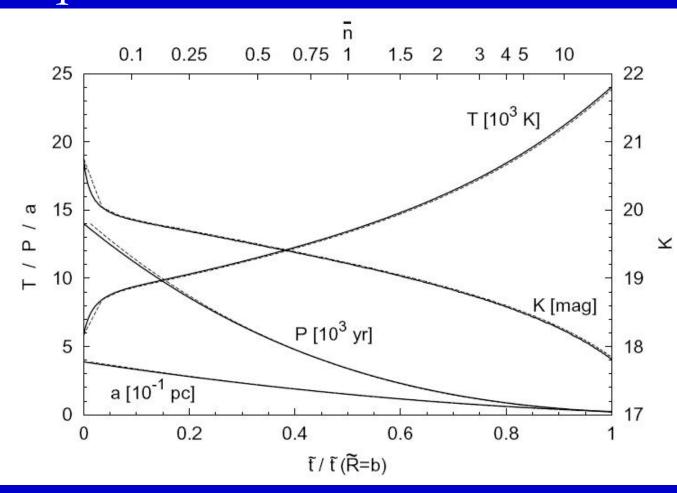


Mass determination from TDE

12 optically/UV selected TDE host galaxies



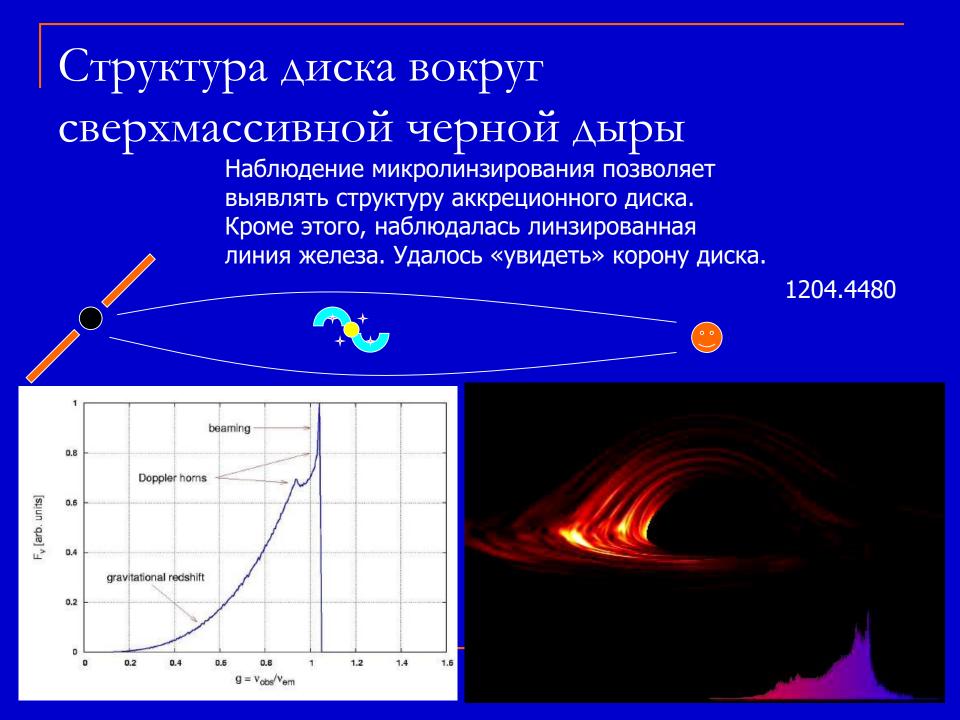
Squeezars



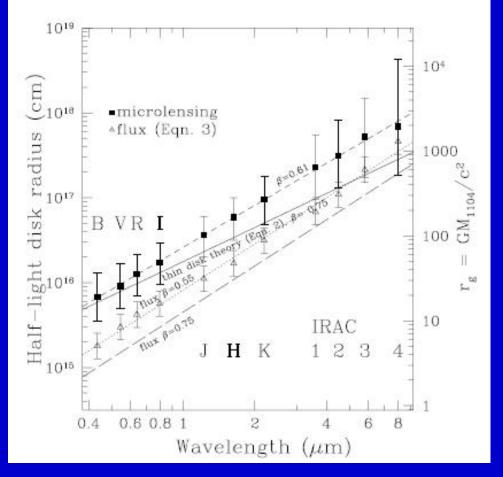
The rate of formation is lower than the rate of tidal disruption events, but the observable time is longer.

Graphs are plotted for a solar-type star orbiting the BH in the center of our Galaxy.

(astro-ph/0305061)



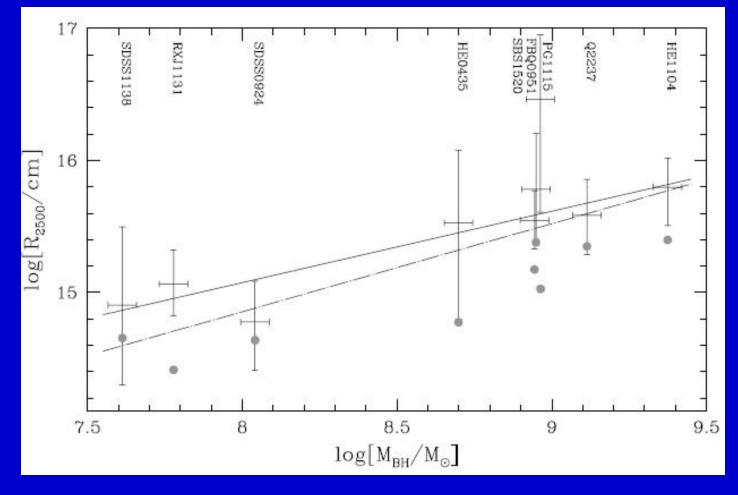
Disc structure from microlensing



Using the data on microlensing at wavelengths 0.4-8 microns it was possible to derive the size of the disc in the quasar HE1104-1805 at different wavelengths.

arXiv:0707.0003 Shawn Poindexter et al. «The Spatial Structure of An Accretion Disk»

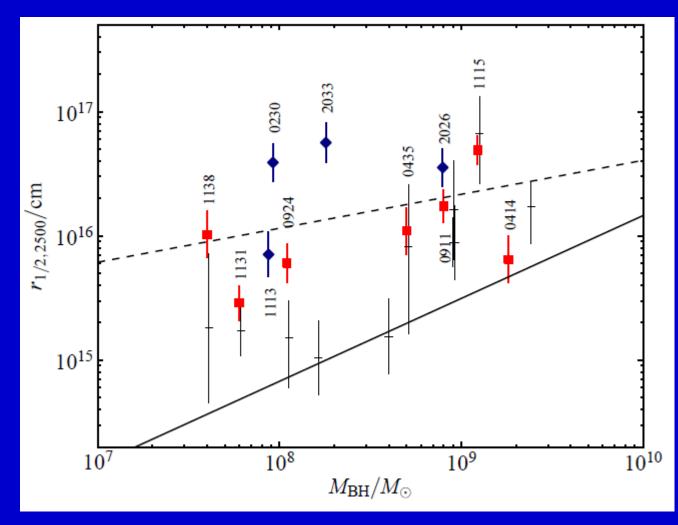
Disc size – BH mass



Disc size can be determined from microlensing.

arXiv:0707.0305 Christopher W. Morgan et al. «The Quasar Accretion Disk Size - Black Hole Mass Relation»

New data



IR and optics.

Chromatic microlensing: blue light from the inner regions is more strongly microlensed than red light from farther out

Solid line is the prediction of the thin disk model.

Standard disc properties

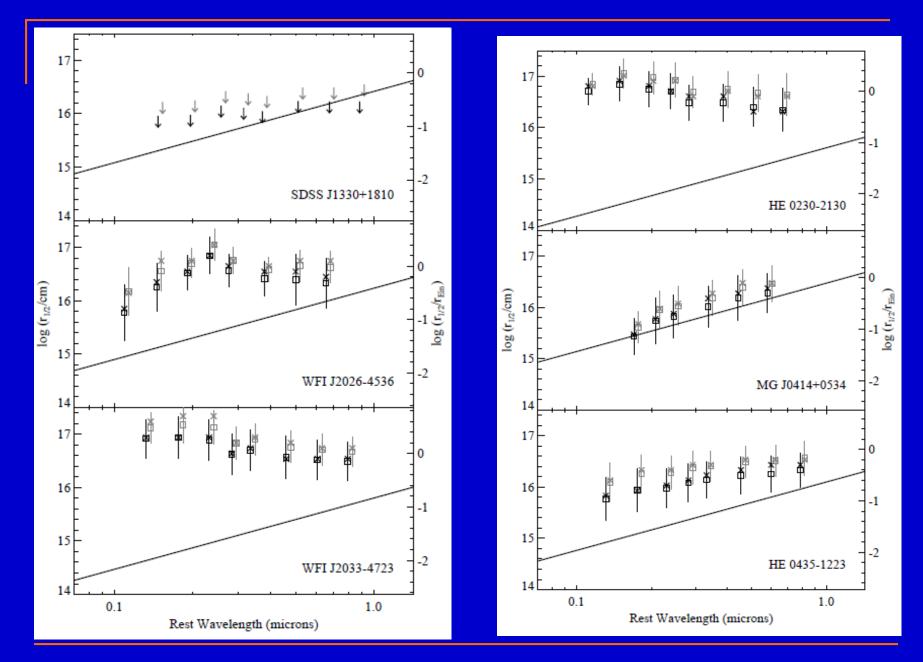
$$T_{\rm eff}(r) = \left(\frac{3G^2 M_{\rm BH}^2 m_p f_{\rm Edd}}{2c\sigma_B \sigma_T \eta r^3}\right)^{1/4} g(r_{\rm in}/r)^{1/4} \,.$$

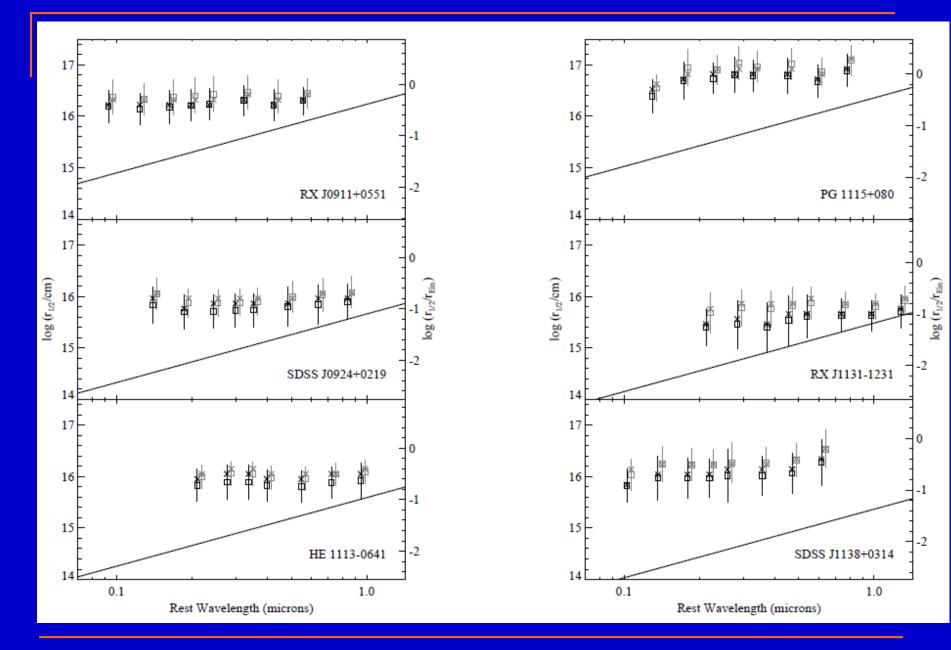
Standard disc model

$$\begin{aligned} r_{1/2} &= 2.44 \left[\frac{45G^2 M_{\rm BH}^2 m_p f_{\rm Edd} \lambda^4}{4\pi^5 h_P c^3 \sigma_T \eta} \right]^{1/3} \sqrt{\cos i} \\ &= 1.68 \times 10^{16} {\rm cm} \left(\frac{M_{\rm BH}}{10^9 M_{\odot}} \right)^{2/3} \left(\frac{f_{\rm Edd}}{\eta} \right)^{1/3} \left(\frac{\lambda}{\mu {\rm m}} \right)^{4/3} \end{aligned}$$

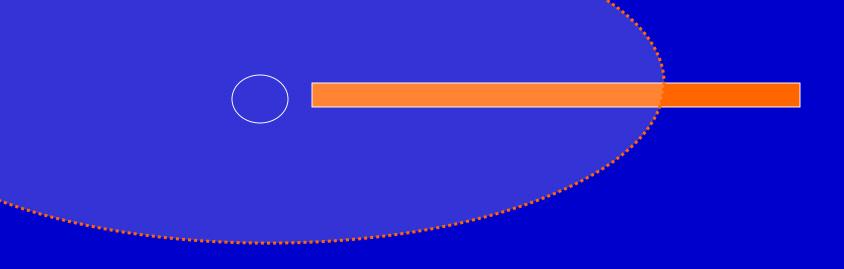
1007.1665

 $r_{1/2} \sim \lambda^{4/3}$



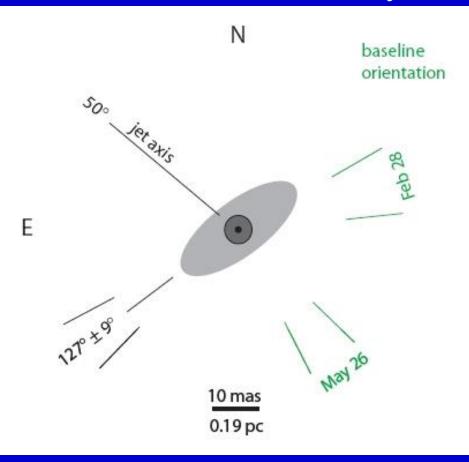


Super-Eddington discs



Super-Eddington accretion leads to formation of an optically-thick envelope scattering the radiation formed in the disc. This makes the apparent disc size larger and practically independent of wavelength

Discs observed by VLTI



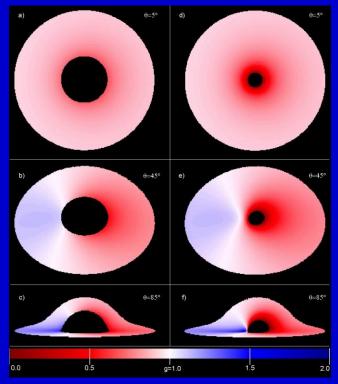
The structure of the disc in Cen A was studied in IR for scales <1 pc.

The data is consistent with a geometrically thin disc with diameter 0.6 pc.

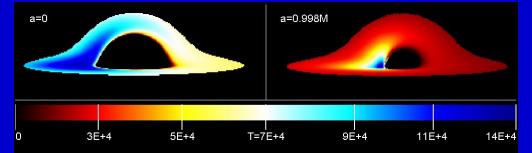
Observations on VLTI.

arXiv:0707.0177 K. Meisenheimer et al. «Resolving the innermost parsec of Centaurus A at mid-infrared wavelengths»

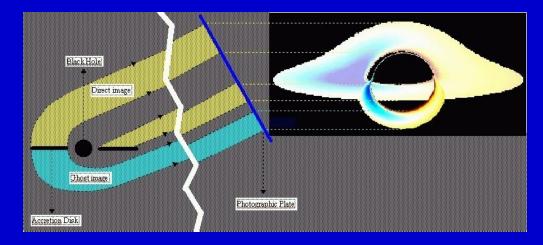
Discs around black holes: a look from aside



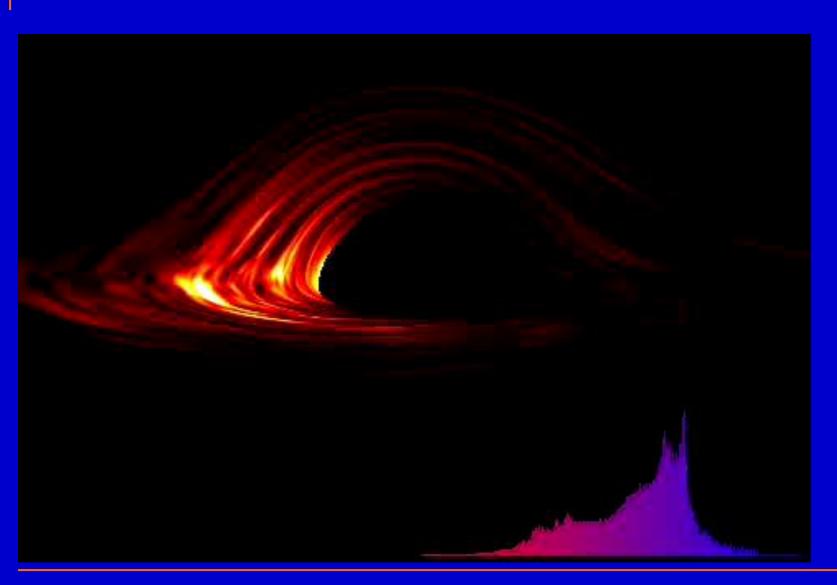
Discs observed from infinity. Left: non-rotating BH, Right: rotating.



Disc temperature

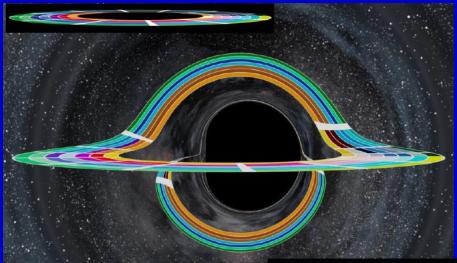


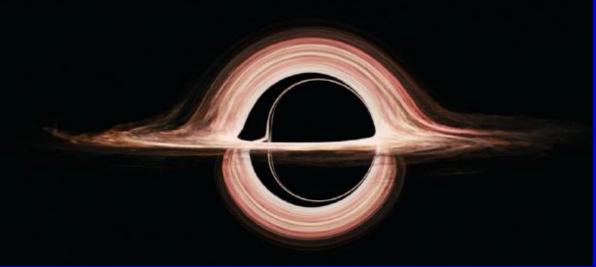
http://web.pd.astro.it/calvani/



from gr-qc/0506078

Discs from Interstellar movie

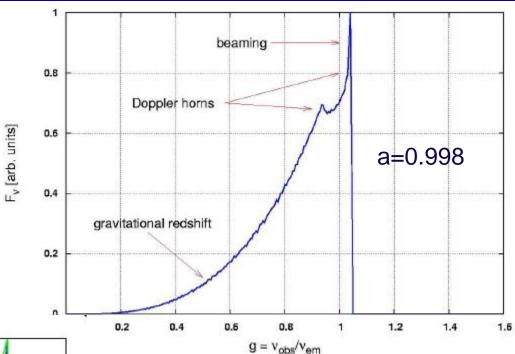


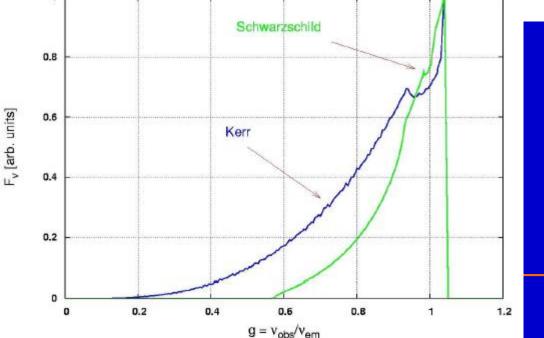


Different effects

$$r_{+} = \frac{GM}{c^2} + \left[\left(\frac{GM}{c^2} \right)^2 - \left(\frac{J}{Mc} \right)^2 \right]^{\frac{1}{2}}$$

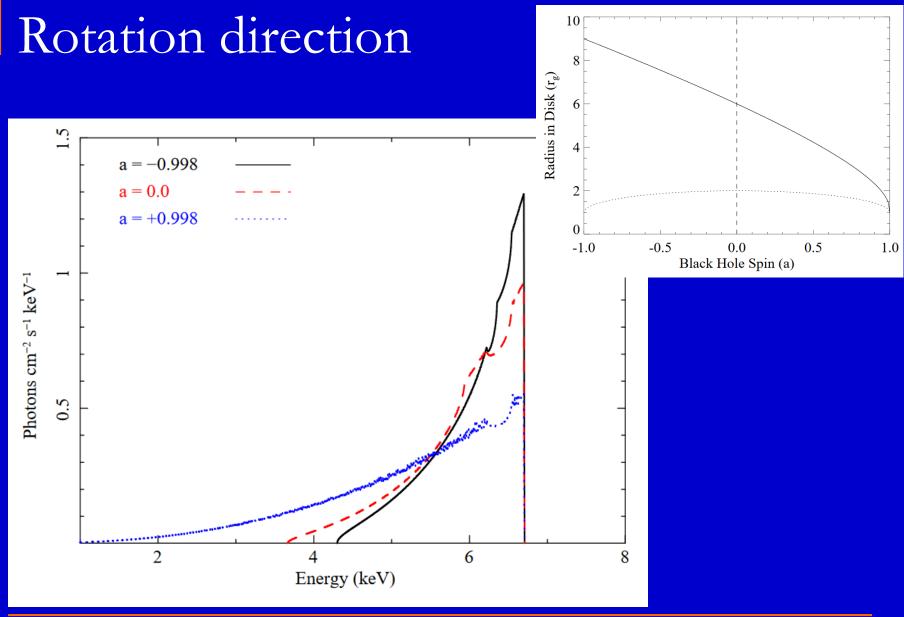
For maximal rotation r_{ISCO}=r₊



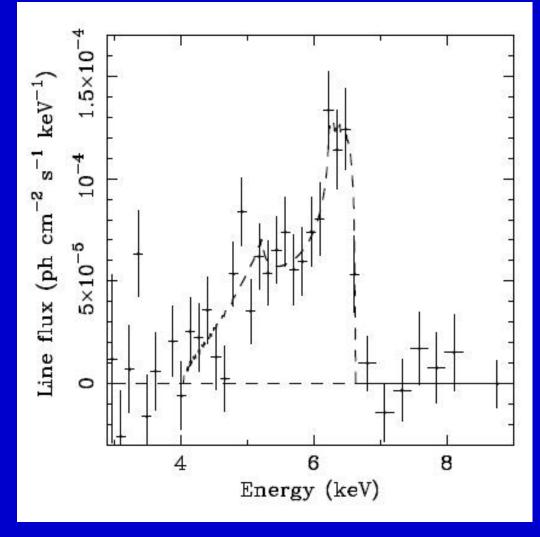


- Doppler effect
- Relativistic beaming
- GR light bending
- GR grav. redshift

arXiv: 0907.3602



Fluorescent lines

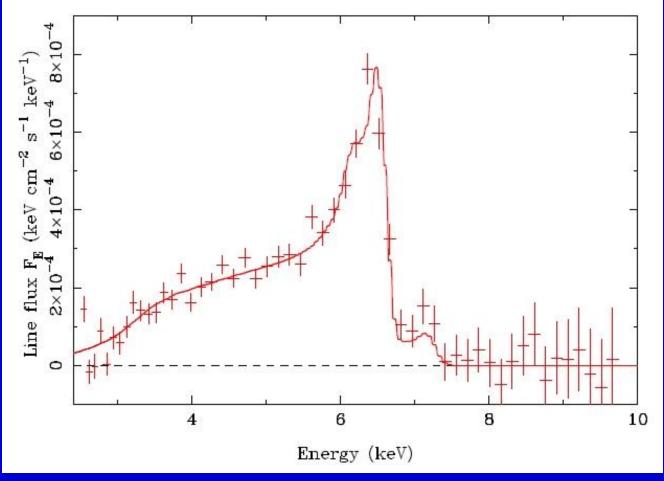


The Kα iron line observed by ASCA (1994 г.). Seyfert galaxy MCG-6-30-15

Dashed line: the model with non-rotating BH, disc inclination 30 degrees.

astro-ph/0212065

Lines and rotation of BHs



XMM-Newton data (astro-ph/0206095)

The fact that the line extends to the red side below 4 keV is interpreted as the sign of rapid rotation (the disc extends inside 3R_g).

see astro-ph/0212065

Suzaku spin measurements program

NGC 3783 z = 0.00973

A very complicated model. a > 0.93 (90% confidence)

AGN	a	A little in 1307
MCG–6-30-15 <u></u> ^{<i>a</i>}	≥ 0.98	and a l
Fairall 9^{b}	$0.65_{-0.05}^{+0.05}$	
SWIFT J2127.4+5654 c	$0.6^{+0.2}_{-0.2}$	
$1H0707-495^{d}$	≥ 0.98	
Mrk 79 <u>e</u>	$0.7^{+0.1}_{-0.1}$	
Mrk 335 <u></u>	$0.70_{-0.01}^{+0.12}$	
NGC 7469^{f}	$0.69^{+0.09}_{-0.09}$	
NGC $3783\underline{^g}$	≥ 0.98	

A little bit more data in 1307.3246 and a big review in 1309.6334

Relativistic reflection fitting of spectra

16

14

12

10

8

6

4

2

0

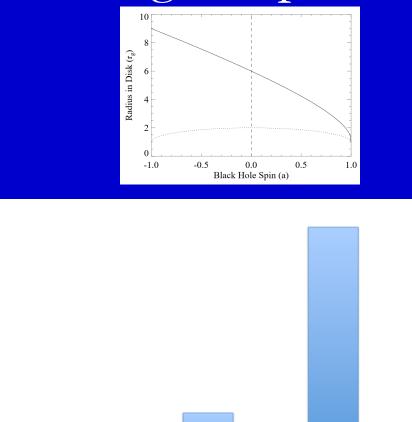
0

0.2

0.4

Number of SMBHs

AGN	2	les M	T /T	Host
	a	log M	$L_{\rm bol}/L_{\rm Edd}$	
MCG-6-30-15 ^a	$\geq +0.98$	$6.65\substack{+0.17\\-0.17}$	$0.40\substack{+0.13\\-0.13}$	E/S0
Fairall 9 ^b	$+0.52^{+0.19}_{-0.15}$	$8.41_{-0.11}^{+0.11}$	$0.05\substack{+0.01\\-0.01}$	\mathbf{Sc}
SWIFT J2127.4+5654 ^{c}	$+0.6^{+0.2}_{-0.2}$	$7.18^{+0.07}_{-0.07}$	$0.18^{+0.03}_{-0.03}$	
$1 \text{ H0707}-495^d$	$\geq +0.98$	$6.70\substack{+0.40\\-0.40}$	$\sim 1.0_{-0.6}$	
Mrk 79 ^e	$+0.7^{+0.1}_{-0.1}$	$7.72_{-0.14}^{+0.14}$	$0.05^{+0.01}_{-0.01}$	SBb
Mrk 335 ^f	$+0.70^{+0.12}_{-0.01}$	$7.15_{-0.13}^{+0.13}$	$0.25_{-0.07}^{+0.07}$	S0a
NGC 3783 ^g	$\geq +0.98$	$7.47^{+0.08}_{-0.08}$	$0.06^{+0.01}_{-0.01}$	SB(r)ab
Ark 120 ^h	$+0.94^{+0.1}_{-0.1}$	$8.18^{+0.05}_{-0.05}$	$0.04^{+0.01}_{-0.01}$	Sb/pec
$3C \ 120^{i}$	≥ 0.95	$7.74_{-0.22}^{+0.20}$	$0.31^{+0.20}_{-0.19}$	S0
$1 \text{ H0419}{-}577^{j}$	$\geq +0.88$	$8.18^{+0.12}_{-0.12}$	$1.27^{+0.42}_{-0.42}$	
Ark 564 ^j	$+0.96^{+0.01}_{-0.06}$	≤ 6.90	≥ 0.11	\mathbf{SB}
Mrk 110 ^j	$\geq +0.99$	$7.40^{+0.09}_{-0.09}$	$0.16\substack{+0.04\\-0.04}$	
SWIFT J0501.9-3239 ^j	$\geq +0.96$			SB0/a(s) pec
Ton S180 ^{j}	$+0.91^{+0.02}_{-0.09}$	$7.30^{+0.60}_{-0.40}$	$2.15^{+3.21}_{-1.61}$	
RBS 1124^{j}	$\geq +0.98$	8.26	0.15	
Mrk 359 ^j	$+0.66\substack{+0.30\\-0.54}$	6.04	0.25	\mathbf{pec}
Mrk 841 ^j	$\geq +0.52$	7.90	0.44	Е
IRAS 13224-3809 ^j	$\geq +0.995$	7.00	0.71	
Mrk 1018 ^j	$+0.58^{+0.36}_{-0.74}$	8.15	0.01	S0
IRAS 00521-7054 ^l	$\geq +0.84$			
NGC 4051 ^m	$\geq +0.99$	6.28	0.03	SAB(rs)bc
NGC 1365^k	$+0.97\substack{+0.01\\-0.04}$	$6.60^{+1.40}_{-0.30}$	$0.06\substack{+0.06\\-0.04}$	SB(s)b



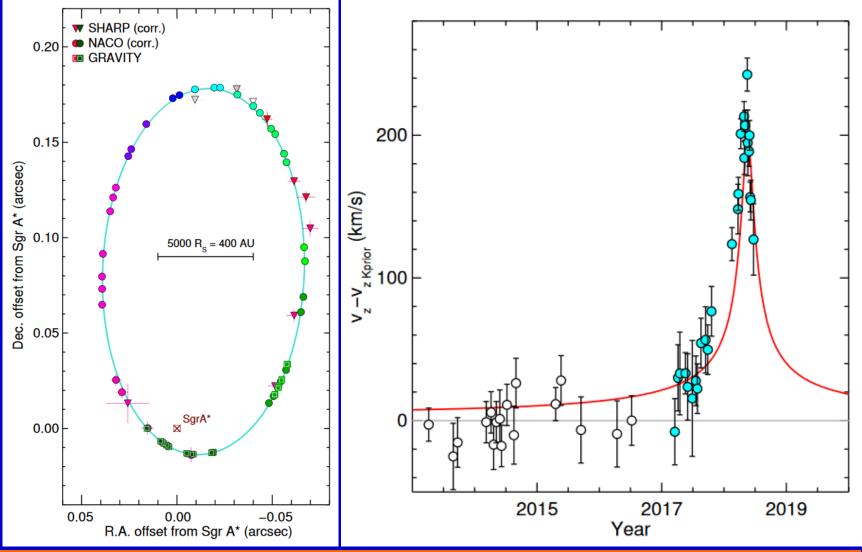
0.6

Spin Value (a)

0.8

1

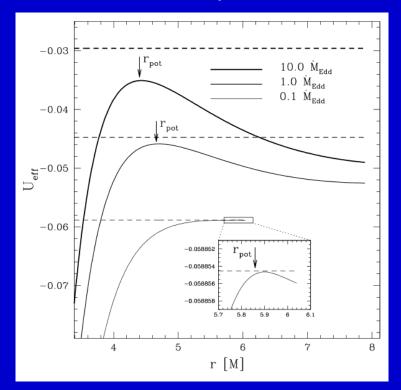
Gravitational redshift of the S2 star

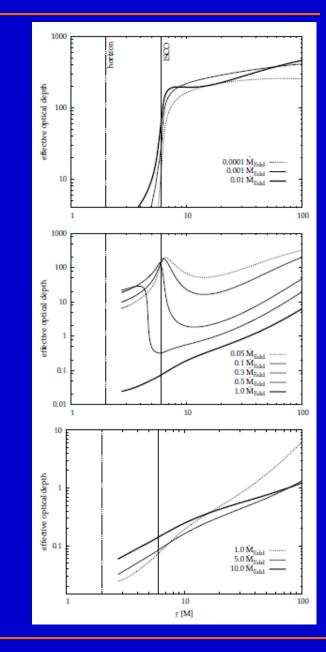


The inner edge

The place where the fluorescent line is formed is not necessarily the standard ISCO.

Especially for large accretion rates the situation is complicated.





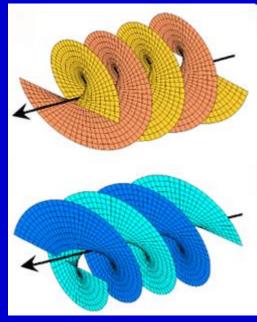
Measuring spins of stellar-mass BHs

Source	Spin a_*
GRS 1915+105	> 0.98
LMC X–1	$0.92^{+0.05}_{-0.07}$
M33 X-7	0.84 ± 0.05
$4U \ 1543 – 47$	0.80 ± 0.05
GRO J1655 -40	0.70 ± 0.05
XTE J1550–564	$0.34_{-0.28}^{+0.20}$
LMC X–3	$< 0.3^{\rm b}$
A0620-00	0.12 ± 0.18

Different methods used 1101.0811

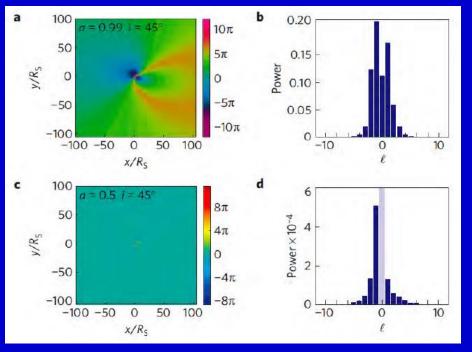
See a review on BH spin (both XRB and AGN) in 1507.06153

Twisted light



This effect can be used to learn about spin of accreting BHs.

If the source of the gravitational field also rotates, it drags space-time with it. Because of the rotation of the central mass, each photon of a light beam propagating along a null geodesic will experience a well-defined phase variation.



See <u>http://www.physics.gla.ac.uk/Optics/play/photonOAM/</u> and astro-ph/0307430 about orbital angular momentum of photons