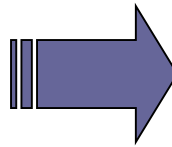
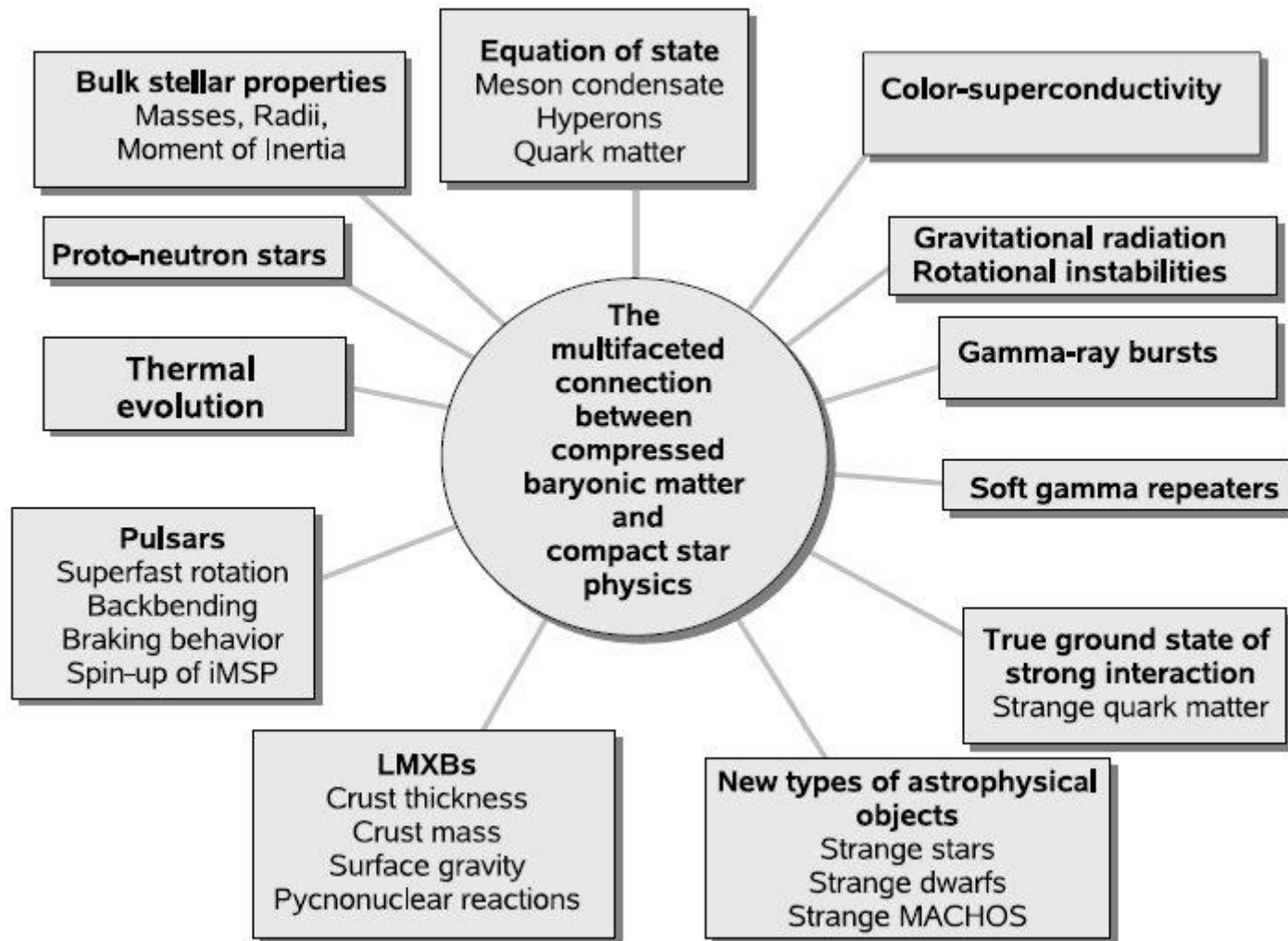


Internal structure of Neutron Stars

Artistic view



Astronomy meets QCD



Hydrostatic equilibrium for a star

$$\left\{ \begin{array}{ll} (1) & \frac{dP}{dr} = -\frac{Gm\rho}{r^2} \quad m = m(r) \\ (2) & \frac{dm}{dr} = 4\pi\rho r^2 \\ (3) & \cancel{\frac{dS}{dt} = Q} \\ (4) & P = P(\rho) \end{array} \right.$$

For NSs we can take $T=0$
and neglect the third equation

For a NS effects of GR are also important.

$$r_g = \frac{2GM}{c^2} \approx 2.95 \frac{M}{M_{\text{SUN}}} \text{ km}$$

$$M/R \sim 0.15 (M/M_{\odot})(R/10 \text{ km})^{-1}$$

$$J/M \sim 0.25 (1 \text{ ms}/P) (M/M_{\odot})(R/10 \text{ km})^2$$

Lane-Emden equation. Polytropes.

$$P = K\rho^\gamma, \quad K, \gamma = \text{const}, \quad \gamma = 1 + \frac{1}{n}$$

$$\frac{dP}{dr} = -\frac{Gm\rho}{r^2} = g\rho, \quad g = -\frac{Gm}{r^2} = -\frac{d\varphi}{dr}$$

$$\frac{dP}{dr} = -\rho \frac{d\varphi}{dr}, \quad \Delta\varphi = 4\pi G\rho$$

$$\rho = \rho_c \Theta^n, \quad \Theta = 1 \text{ при } r = 0$$

$$P = K\rho_c^{1+1/n} \Theta^{1+n}, \quad \frac{dP}{dr} = (n+1)K\rho_c^{1+1/n} \Theta^n \frac{d\Theta}{dr}$$

$$\frac{d\varphi}{dr} = -(n+1)K\rho_c^{1/n} \frac{d\Theta}{dr}$$

$$\Delta\Theta = -\frac{4\pi G\rho_c^{1-1/n}}{(n+1)K} \Theta^n$$

$$\xi = r/a, \quad a^2 = (n+1)K\rho_c^{1/n-1}/(4\pi G)$$

$$\frac{1}{\xi^2} \frac{d}{d\xi} \xi^2 \frac{d}{d\xi} \Theta = -\Theta^n$$

$$\Theta = \Theta(\xi)$$

$$0 \leq \xi \leq \xi_1$$

$$\Theta(0) = 1, \quad \Theta'(0) = 0$$

$$\Theta(\xi_1) = 0$$

Properties of polytropic stars

Analytic solutions:

$$\begin{array}{lll}
 n=0 & \Theta = 1 - \frac{\xi^2}{6} & \xi_1 = \sqrt{6} \\
 n=1 & \Theta = \frac{\sin \xi}{\xi} & \xi_1 = \pi \\
 n=5 & \Theta = \frac{1}{\sqrt{1 + \xi^2/3}} & \xi_1 = \infty
 \end{array}$$

$$M = 4\pi \int_0^R dr r^2 \rho = 4\pi \rho_c a^3 \xi_1^2 |\Theta'(\xi_1)|$$

$$\frac{\rho_c}{\rho} = \frac{4\pi R^3 \rho_c}{3M} = \frac{\xi_1}{3|\Theta'(\xi_1)|}$$

↓ $\gamma=5/3$

↓ $\gamma=4/3$

n	0	1	1.5	2	3
ξ_1	2.449	3.142	3.654	4.353	6.897
$ \Theta'_1 $	0.7789	0.3183	0.2033	0.1272	0.04243
$\rho_c / \bar{\rho}$	1	3.290	5.991	11.41	54.04

$$M \sim \rho_c^{(3-n)/(2n)}$$

$$R \sim \rho_c^{(1-n)/(2n)}$$

$$M \sim R^{(3-n)/(1-n)}$$

$$n=0 \quad M \sim R^3$$

$$n=1 \quad M \sim \rho_c \quad R = \text{const}$$

$$n=1.5 \quad M \sim \sqrt{\rho_c} \sim R^{-3}$$

$$n=3 \quad M = \text{const} \quad R \sim \rho_c^{-1/3}$$

Useful equations

White dwarfs

1. Non-relativistic electrons

$$\gamma=5/3, K=(3^{2/3} \pi^{4/3} /5) (\hbar^2/m_e m_u^{5/3} \mu_e^{5/3});$$

μ_e -mean molecular weight per one electron

$$K=1.0036 \cdot 10^{13} \mu_e^{-5/3} \text{ (CGS)}$$

2. Relativistic electrons

$$\gamma=4/3, K=(3^{1/3} \pi^{2/3} /4) (\hbar c/m_u^{4/3} \mu_e^{4/3});$$
$$K=1.2435 \cdot 10^{15} \mu_e^{-4/3} \text{ (CGS)}$$

Neutron stars

1. Non-relativistic neutrons

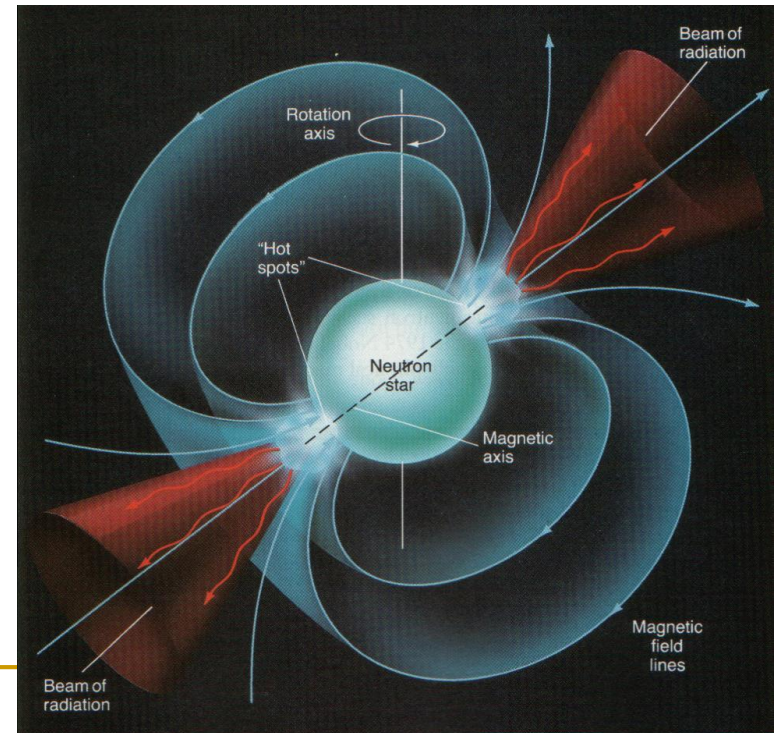
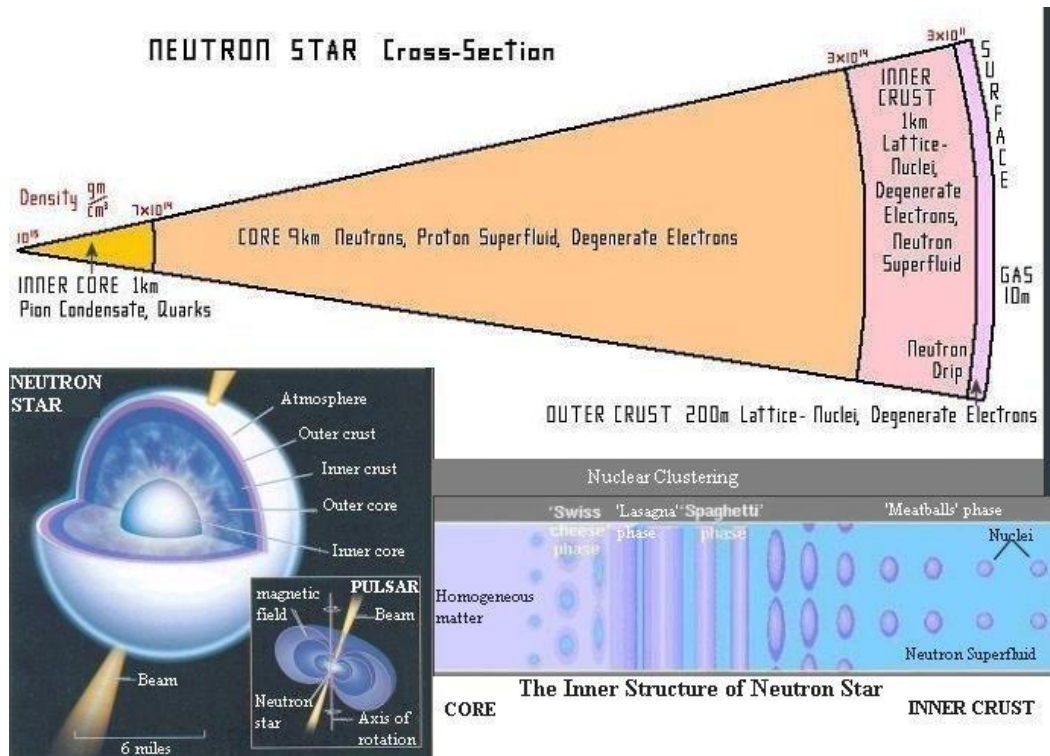
$$\gamma=5/3, K=(3^{2/3} \pi^{4/3} /5) (\hbar^2/m_n^{8/3});$$
$$K=5.3802 \cdot 10^9 \text{ (CGS)}$$

2. Relativistic neutrons

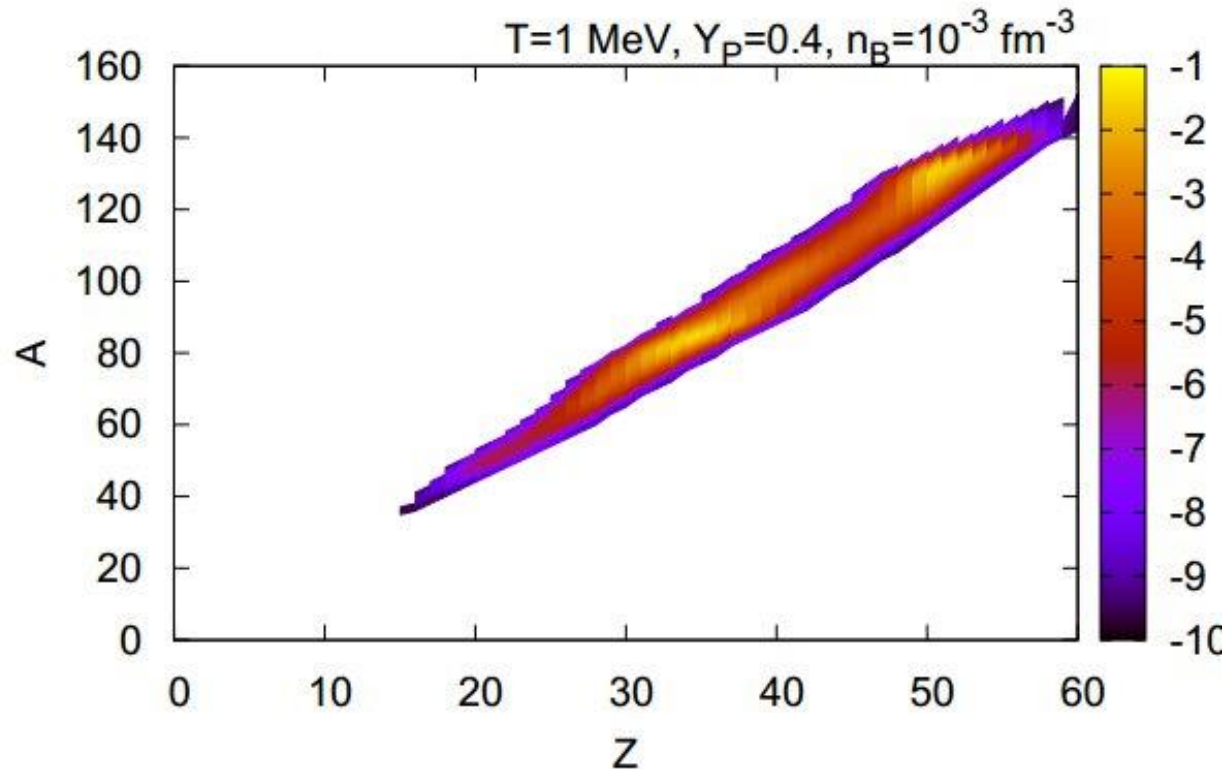
$$\gamma=4/3, K=(3^{1/3} \pi^{2/3} /4) (\hbar c/m_n^{4/3});$$
$$K=1.2293 \cdot 10^{15} \text{ (CGS)}$$

Neutron stars

Superdense matter and superstrong magnetic fields



Proto-neutron stars



Mass fraction of nuclei in the nuclear chart for matter at $T = 1 \text{ MeV}$, $n_B = 10^{-3} \text{ fm}^{-3}$, and $Y_p = 0.4$ (proton fraction). Different colors indicate mass fraction in Log_{10} scale.

1202.5791

NS EoS are also important for SN explosion calculation, see 1207.2184

EoS for core-collapse, proto-NS and NS-NS mergers

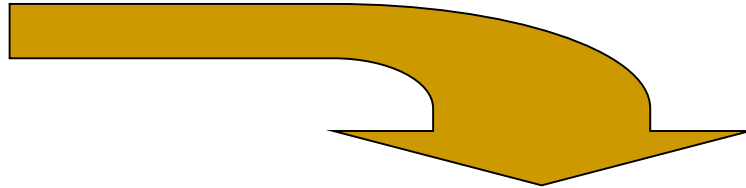
	Core-collapse supernovae	Proto-neutron stars	Mergers of compact binary stars
n/n_s	$10^{-8} - 10$	$10^{-8} - 10$	$10^{-8} - 10$
$T(\text{MeV})$	0 - 30	0 - 50	0 - 100
Y_e	0.35 - 0.45	0.01 - 0.3	0.01 - 0.6
$S(k_B)$	0.5 - 10	0 - 10	0 - 100

Wide ranges of parameters

Astrophysical point of view

**Astrophysical appearance of NSs
is mainly determined by:**

- Spin
- Magnetic field
- Temperature
- Velocity
- Environment






The first four are related to the NS structure!

Equator and radius

$$ds^2 = c^2 dt^2 e^{2\Phi} - e^{2\lambda} dr^2 - r^2 [d\theta^2 + \sin^2\theta d\varphi^2]$$

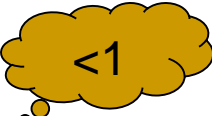
In flat space $\Phi(r)$ and $\lambda(r)$ are equal to zero.

• $t = \text{const}, r = \text{const}, \theta = \pi/2, 0 < \varphi < 2\pi$  $l = 2\pi r$

• $t = \text{const}, \theta = \text{const}, \varphi = \text{const}, 0 < r < r_0$  $dl = e^\lambda dr$  $l = \int_0^{r_0} e^\lambda dr \neq r_0$

Gravitational redshift

$$d\tau = dt e^{\Phi},$$



$$\nu_r = \frac{dN}{d\tau} = e^{-\Phi} \frac{dN}{dt} \longrightarrow \text{Frequency emitted at } r$$

$$r \rightarrow \infty \quad \Phi \rightarrow 0 \quad \nu_{\infty} = \frac{dN}{dt} \longrightarrow \text{Frequency detected by an observer at infinity}$$

$$\nu_{\infty} = \nu_r e^{\Phi} \Rightarrow \Phi(r) \longrightarrow \text{This function determines gravitational redshift}$$

$$e^{2\lambda} \equiv \frac{1}{1 - \frac{2Gm}{c^2 r}}$$


It is useful to use $m(r)$ – gravitational mass inside r – instead of $\lambda(r)$

Outside of the star

$$r > R \Rightarrow m(r) = M = \text{const}$$

$$e^{2\Phi} = 1 - \frac{2GM}{c^2 r} = 1 - \frac{r_g}{r}, \quad r_g = \frac{2GM}{c^2}$$

$$ds^2 = \left(1 - \frac{r_g}{r}\right) c^2 dt^2 - \left(1 - \frac{r_g}{r}\right)^{-1} dr^2 - r^2 d\Omega^2$$

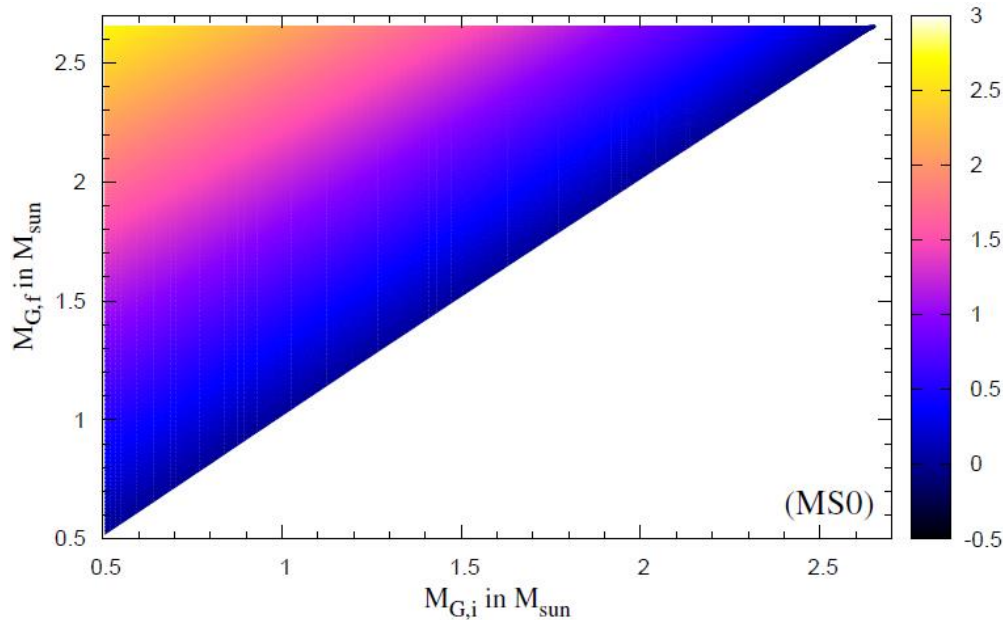
$$v_\infty = v_r \sqrt{1 - \frac{r_g}{r}}$$


redshift

Bounding energy  $\Delta M = M_b - M \sim 0.2 M_{\text{sun}}$

Apparent radius  $R_\infty = R / \sqrt{1 - r_g / R}$

Bounding energy



If you drop a kilo on a NS, then you increase its mass for < kilo

M_{acc} is shown with color

$M_{G,i}$ (M_{\odot})	ΔM_G (M_{\odot})	$M_{B,i}$ (M_{\odot})		$M_{\text{acc}} (\Delta M_B)$ (M_{\odot})	
		APR	MS0	APR	MS0
1.4	0.57	1.554	1.525	0.768	0.712
1.5	0.47	1.681	1.647	0.641	0.591
1.6	0.37	1.811	1.767	0.511	0.470
1.7	0.27	1.943	1.892	0.379	0.345
1.8	0.17	2.080	2.018	0.242	0.219
1.9	0.07	2.221	2.146	0.101	0.091


$$M_{\text{acc}} = \Delta M_G + \Delta BE / c^2 = \Delta M_B$$

BE- binding energy

$$BE = (M_B - M_G) c^2$$

TOV equation

$$R_{ik} - \frac{1}{2} g_{ik} R = \frac{8\pi G}{c^4} T_{ik}$$



(1) $\frac{dP}{dr} = -\frac{G\rho m}{r^2} \left(1 + \frac{P}{\rho c^2}\right) \left(1 + \frac{4\pi r^3 P}{mc^2}\right) \left(1 - \frac{2Gm}{rc^2}\right)^{-1}$

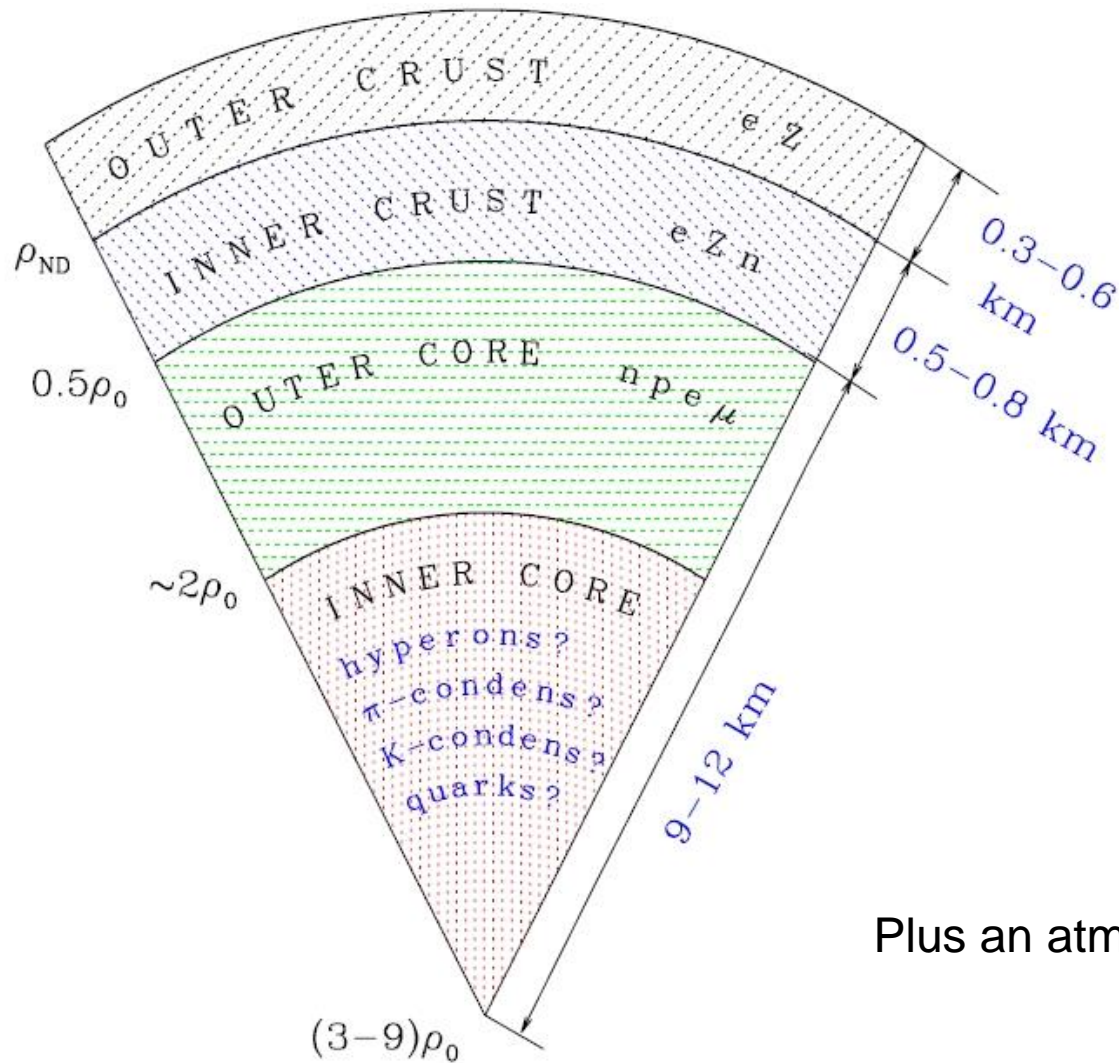
(2) $\frac{dm}{dr} = 4\pi r^2 \rho$

(3) $\frac{d\Phi}{dr} = -\frac{1}{\rho c^2} \frac{dP}{dr} \left(1 + \frac{P}{\rho c^2}\right)^{-1}$

(4) $P = P(\rho)$

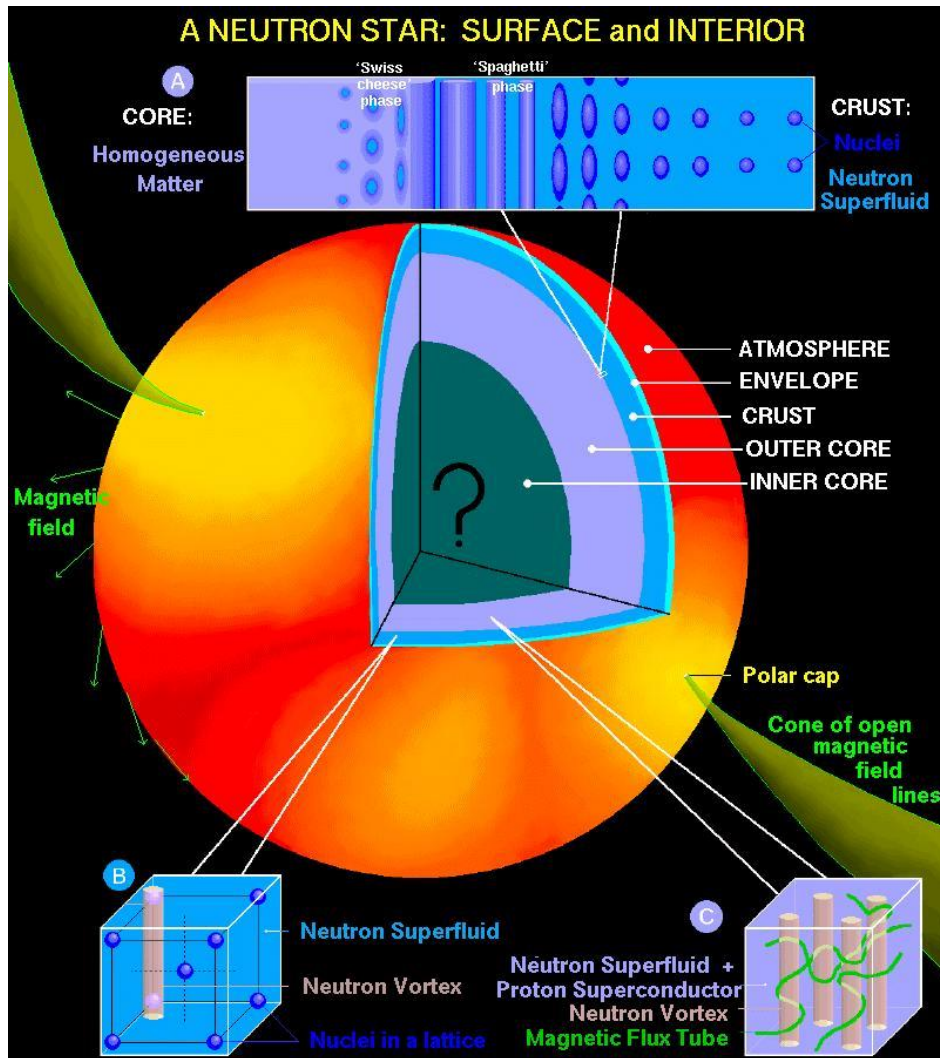
**Tolman (1939)
Oppenheimer-
Volkoff (1939)**

Structure and layers



Plus an atmosphere...

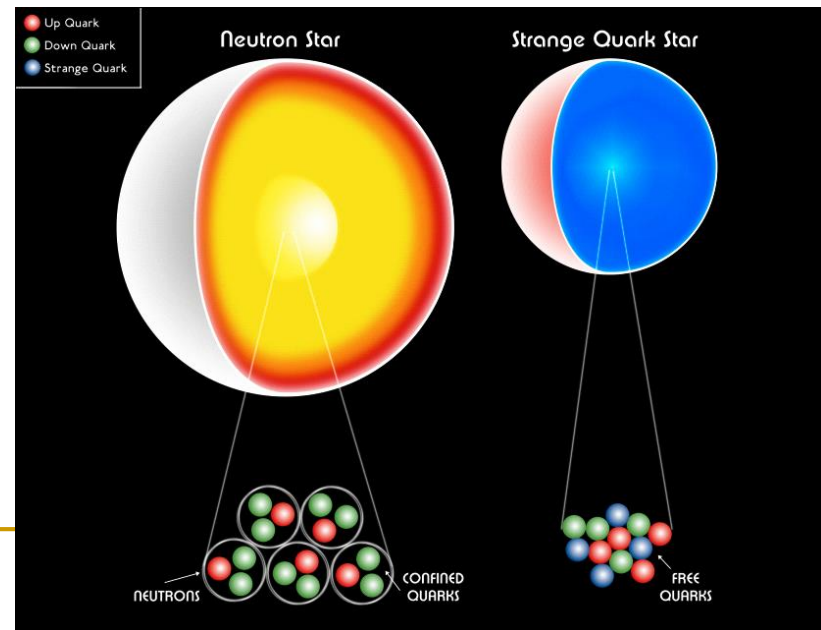
Neutron star interiors



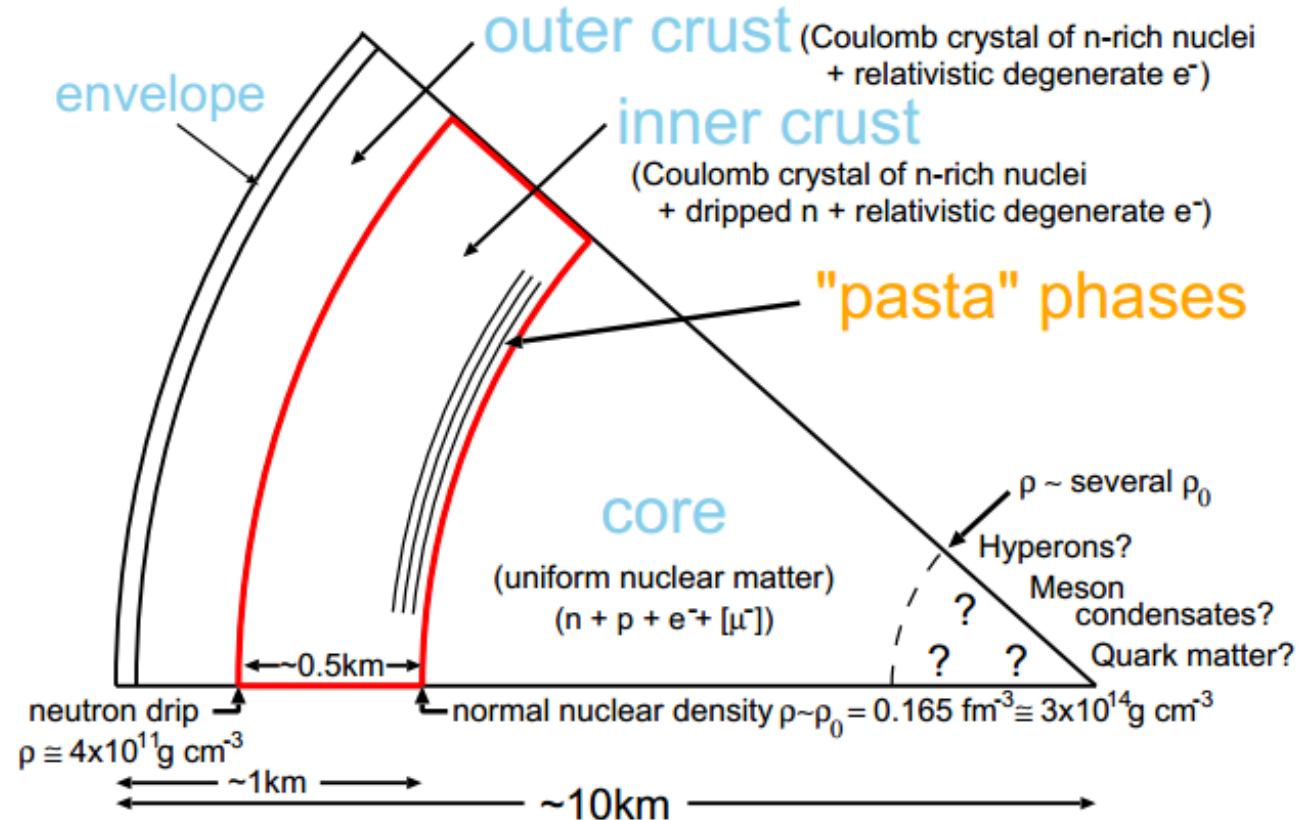
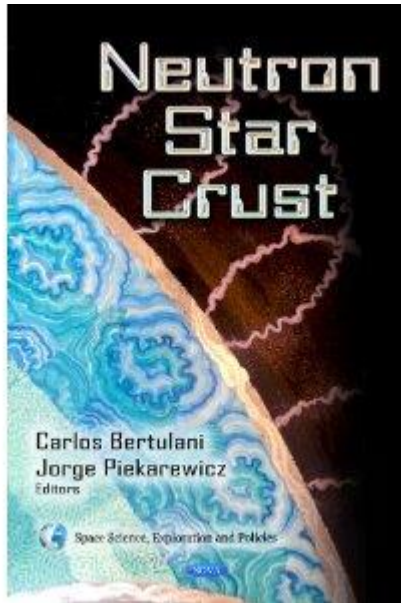
Radius: 10 km

Mass: 1-2 solar

Density: above the nuclear
Strong magnetic fields



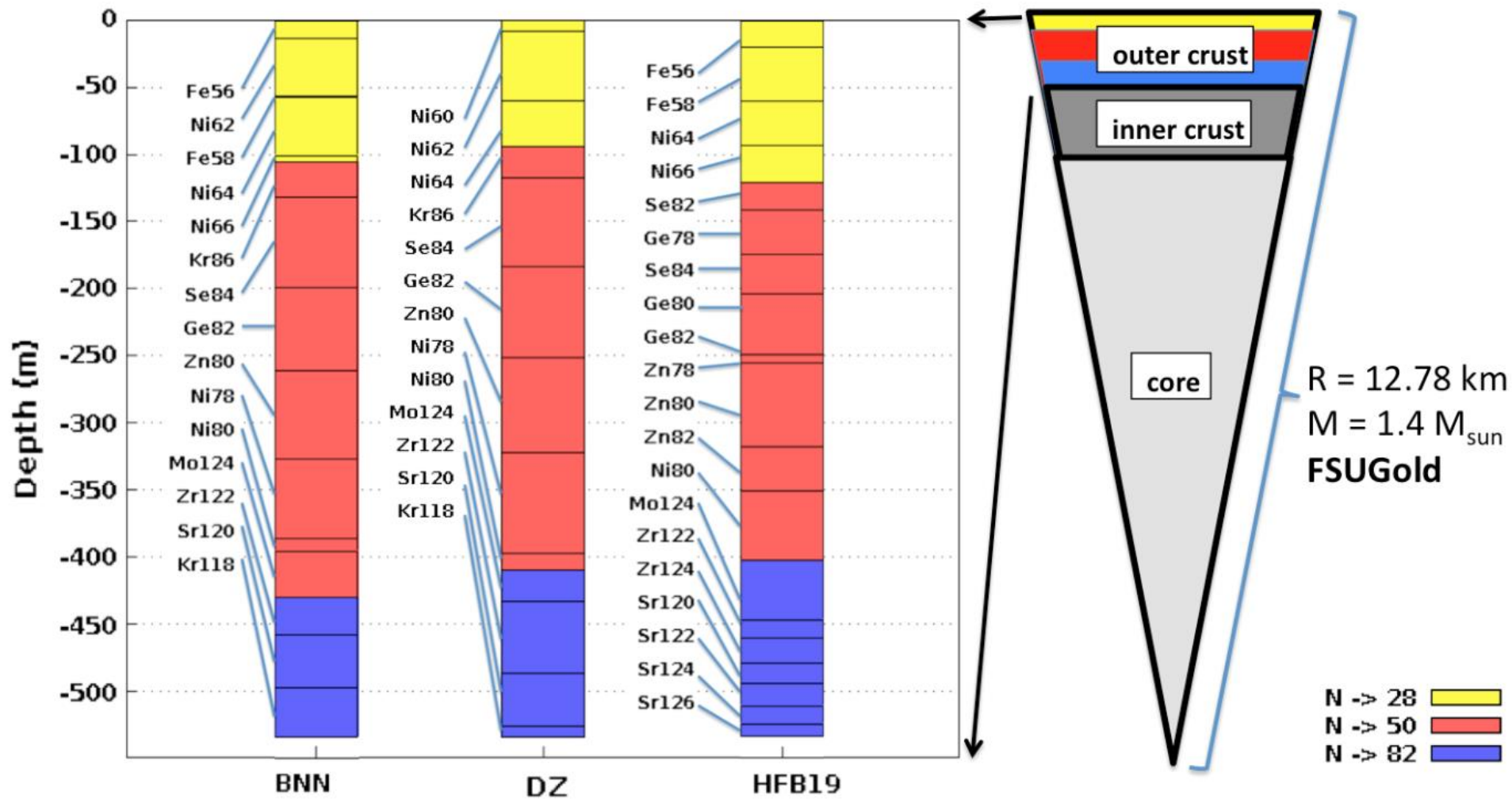
Neutron star crust



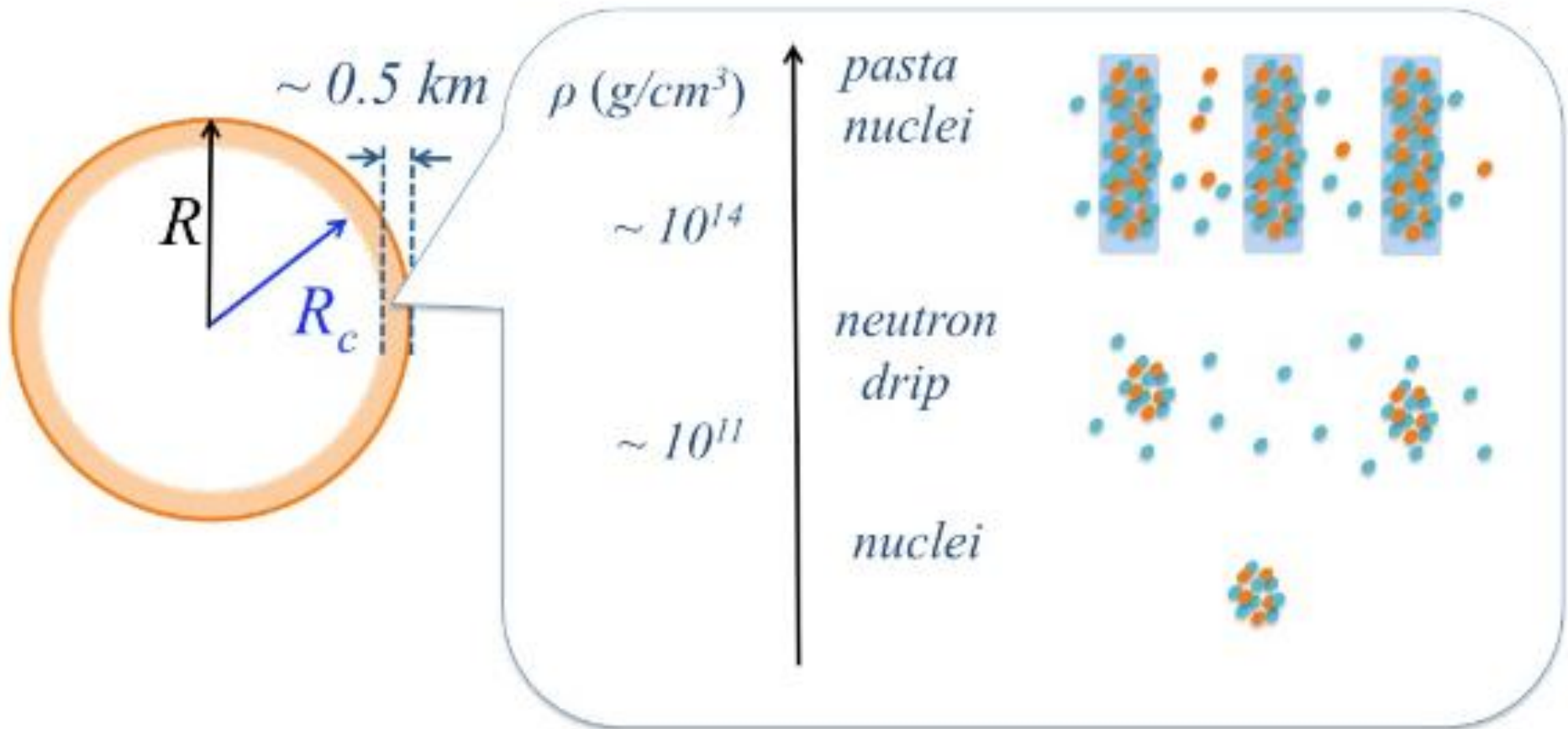
Many contributions to the book are available in the arXiv.

Mechanical properties of crusts are continuously discussed, see 1208.3258, 1808.06415

Element composition of the crust

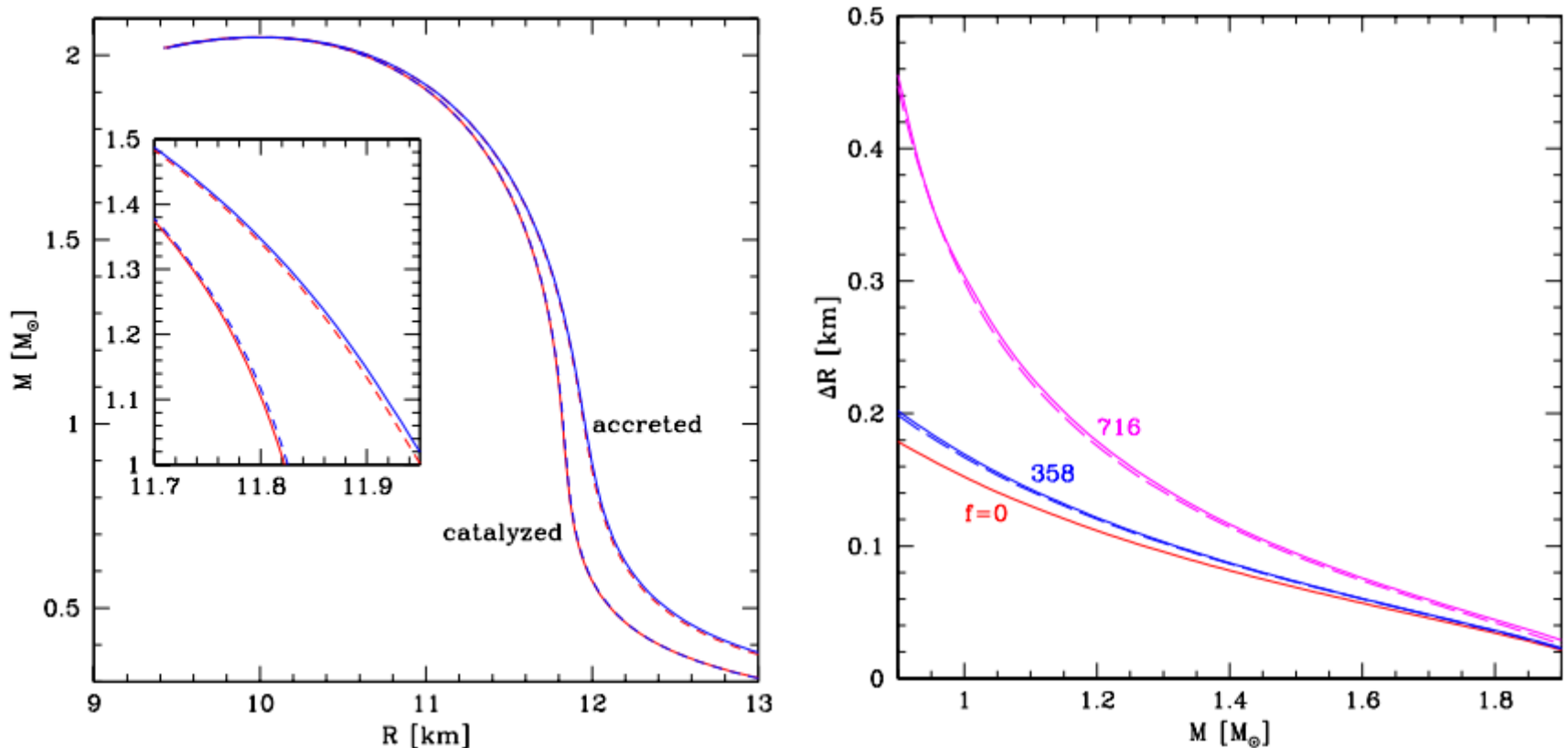


Inner crust properties



Accreted crust

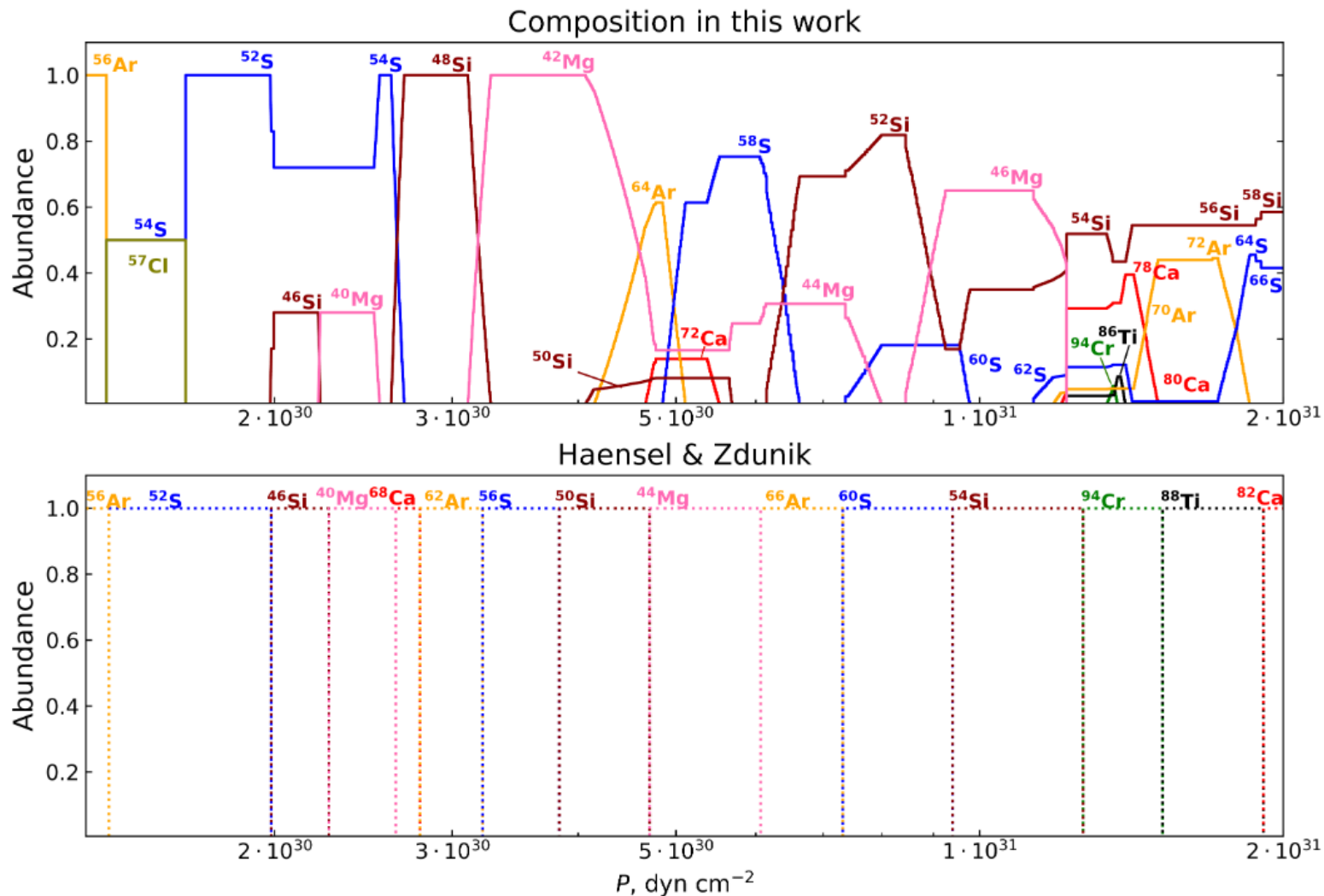
It is interesting that the crust formed by accreted matter differs from the crust formed from catalyzed matter. The former is thicker.



1104.0385

See new results in 1910.03932

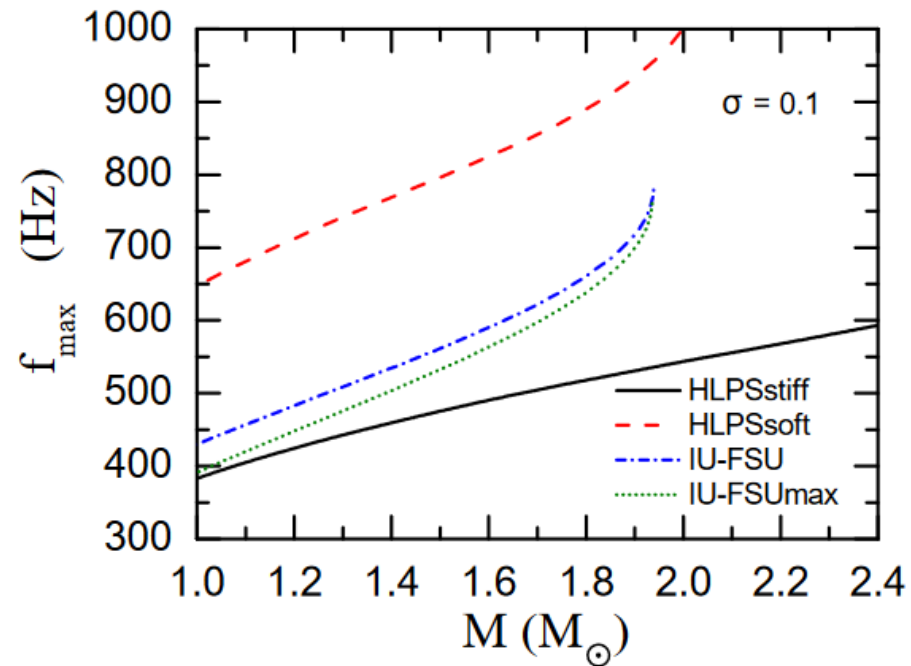
Composition of accreted crust



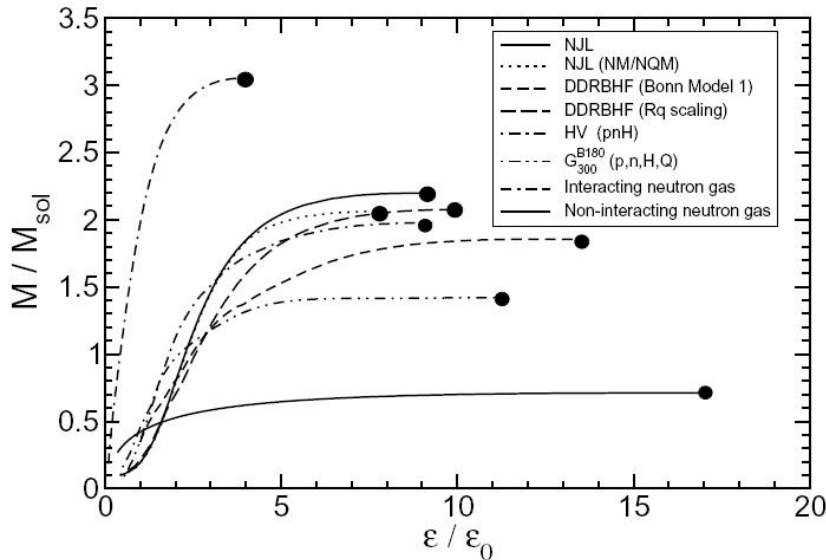
Crust and limiting rotation

Model	σ	$f_{\text{in}}^{1.4}$ (Hz)	$f_{\text{fin}}^{1.4}$ (Hz)	$f_{\text{in}}^{1.8}$ (Hz)	$f_{\text{fin}}^{1.8}$ (Hz)
HLPSStiff	0.05	0	326	35	368
	0.10	136	479	236	569
IU-FSU	0.05	349	515	909	1022
	0.10	781	947	1875	1988
IU-FSU _{max}	0.05	35	358	374	586
	0.10	232	555	854	1066

Failure of the crust can be the reason of the limiting frequency.
Spinning-up of a NS due to accretion can result in crust failure.
Then the shape of the star is deformed, it gains ellipticity.
So, GWs are emitted which slow down the compact object.



Configurations

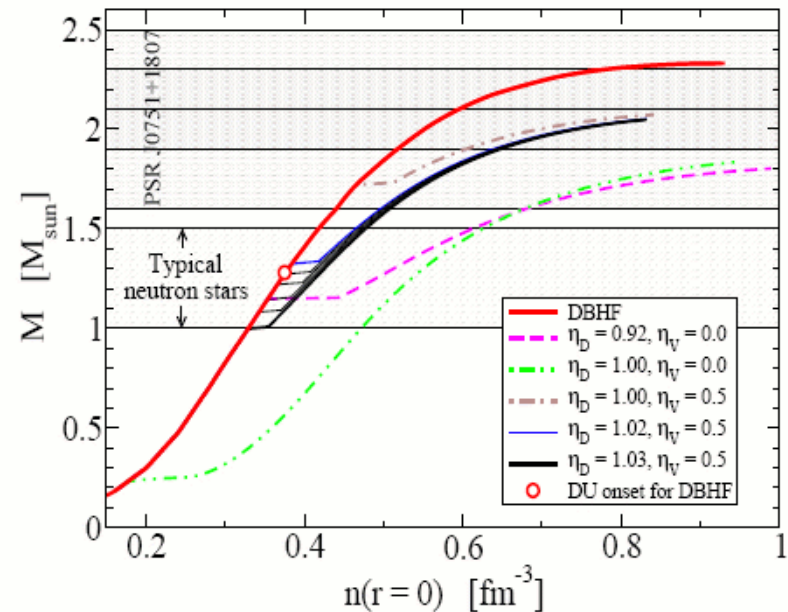


A RNS code is developed
and made available to the public
by Sterlgioulas and Friedman
ApJ 444, 306 (1995)
<http://www.gravity.phys.uwm.edu/rns/>

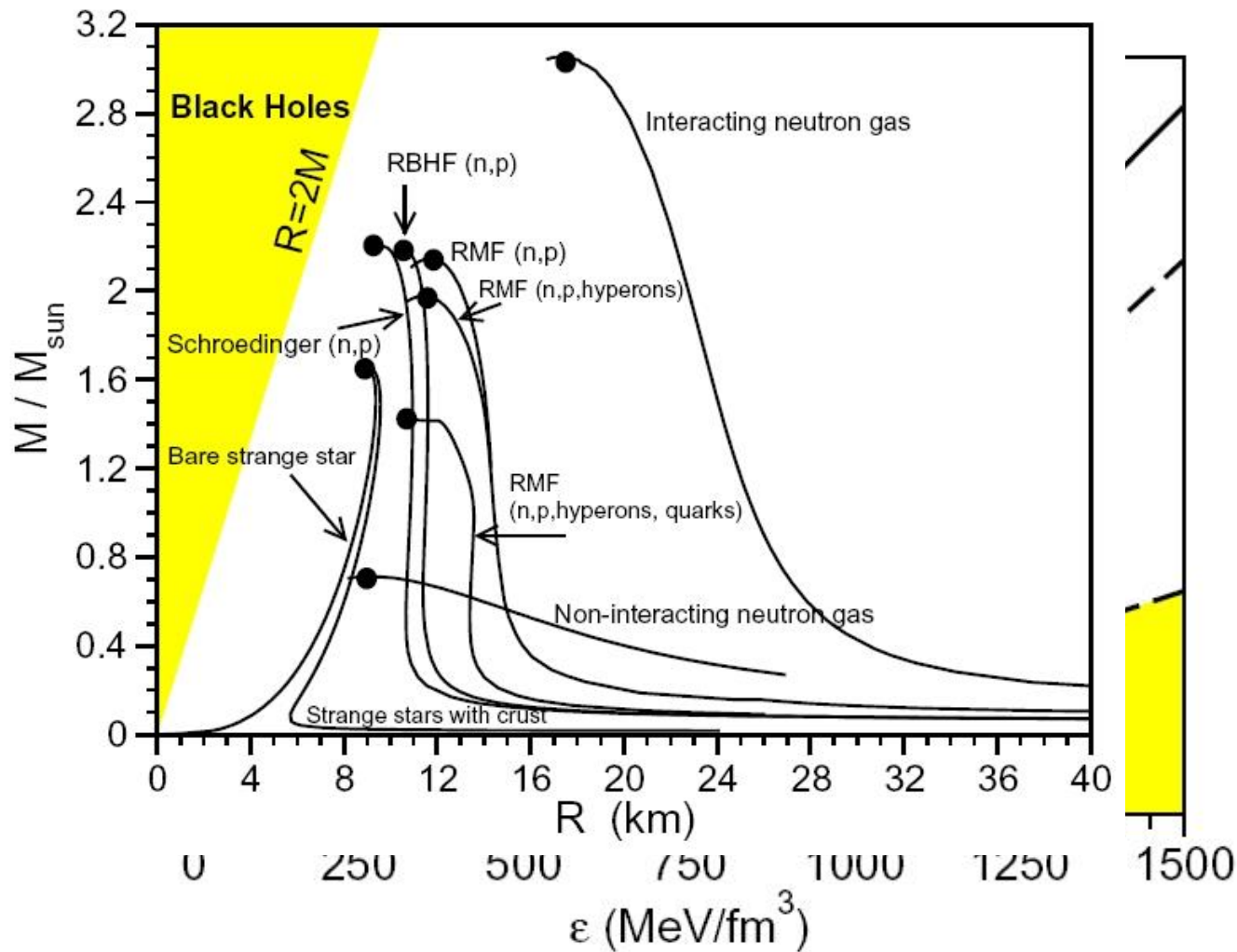
NS mass vs.
central density
(Weber et al.
arXiv: 0705.2708)



Stable configurations
for neutron stars and
hybrid stars
(astro-ph/0611595).



EoS



(Weber et al. ArXiv: 0705.2708)

Mass-radius

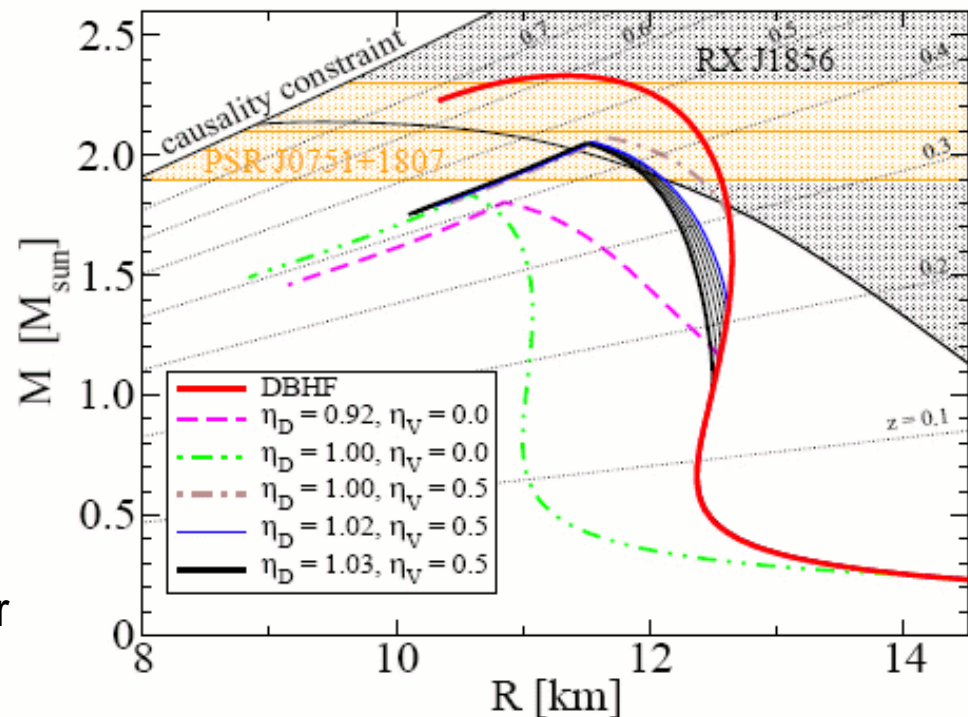
Mass and radius are macroscopical potentially measured parameters.

Thus, it is important to formulate EoS in terms of these two parameters.

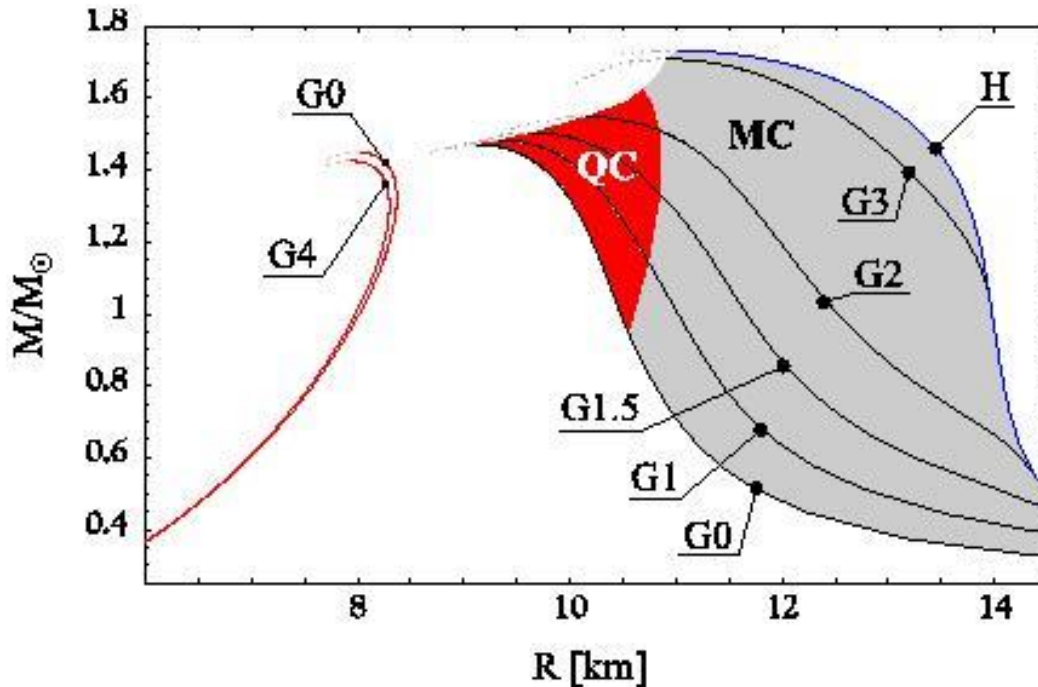
About hyperon stars see a review in 1002.1658.

About strange stars and some other exotic options – 1002.1793

Mass-radius relations for CSs with possible phase transition to deconfined quark matter.



Mass-radius relation



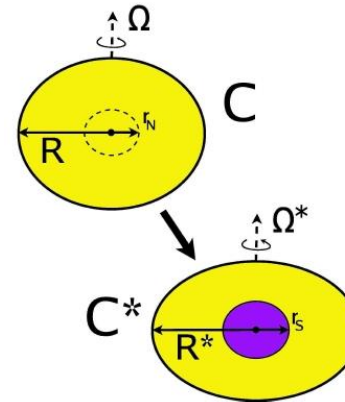
Main features

- Max. mass
- Diff. branches (quark and normal)
- Stiff and soft EoS
- Small differences for realistic parameters
- Softening of an EoS with growing mass

Rotation is neglected here.
Obviously, rotation results in:

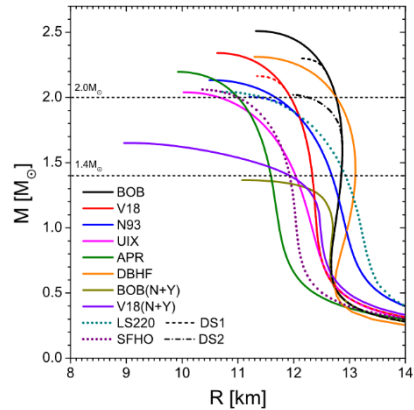
- larger max. mass
- larger equatorial radius

Spin-down can result in phase transition, as well as spin-up (due to accreted mass), see 1109.1179



Haensel, Zdunik
astro-ph/0610549

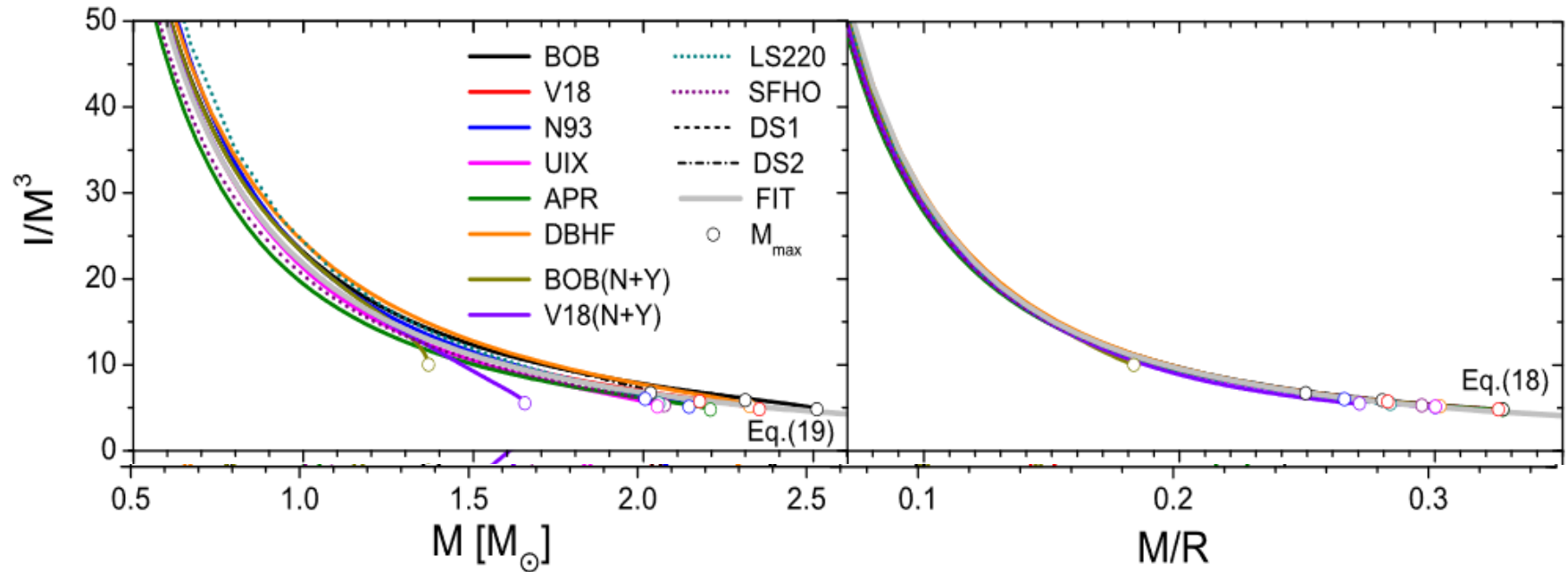
Fitting formulae for moment of inertia



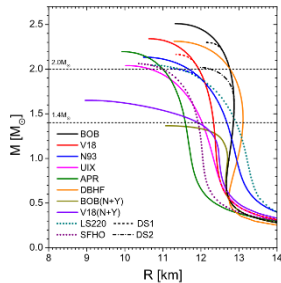
$$\frac{I}{M^3} \equiv 0.8134 \beta^{-1} + 0.2101 \beta^{-2} + 0.003175 \beta^{-3} - 0.0002717 \beta^{-4} \quad (18)$$

$$\frac{I}{M^3} \equiv 1.0334 M^{-1} + 30.7271 M^{-2} - 12.8839 M^{-3} + 2.8841 M^{-4} \quad (19)$$

$$\beta = Gm/(\bar{R}c^2)$$



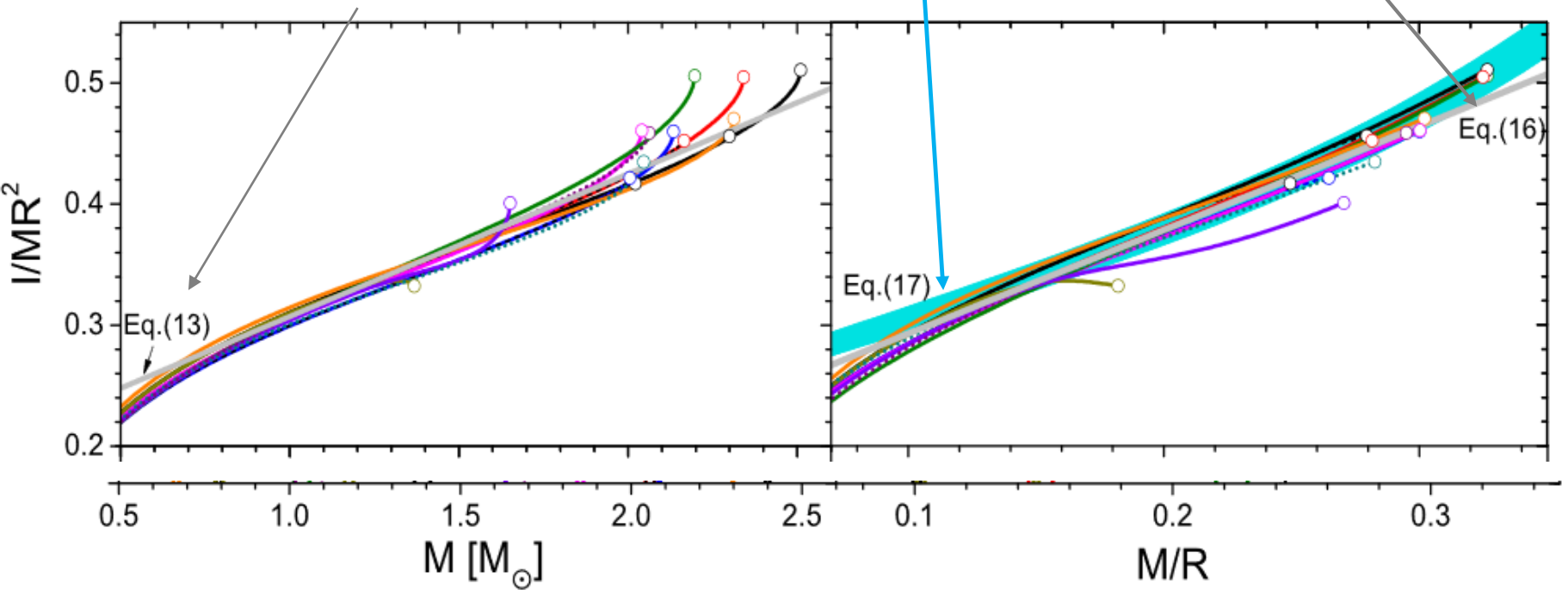
Fits for I/MR^2



$$\frac{I}{MR^2} \approx 0.189 + 0.118 \frac{M}{M_{\odot}} \pm 0.016.$$

$$\frac{I}{MR^2} \approx (0.237 \pm 0.008)(1 + 2.844\beta + 18.91\beta^4)$$

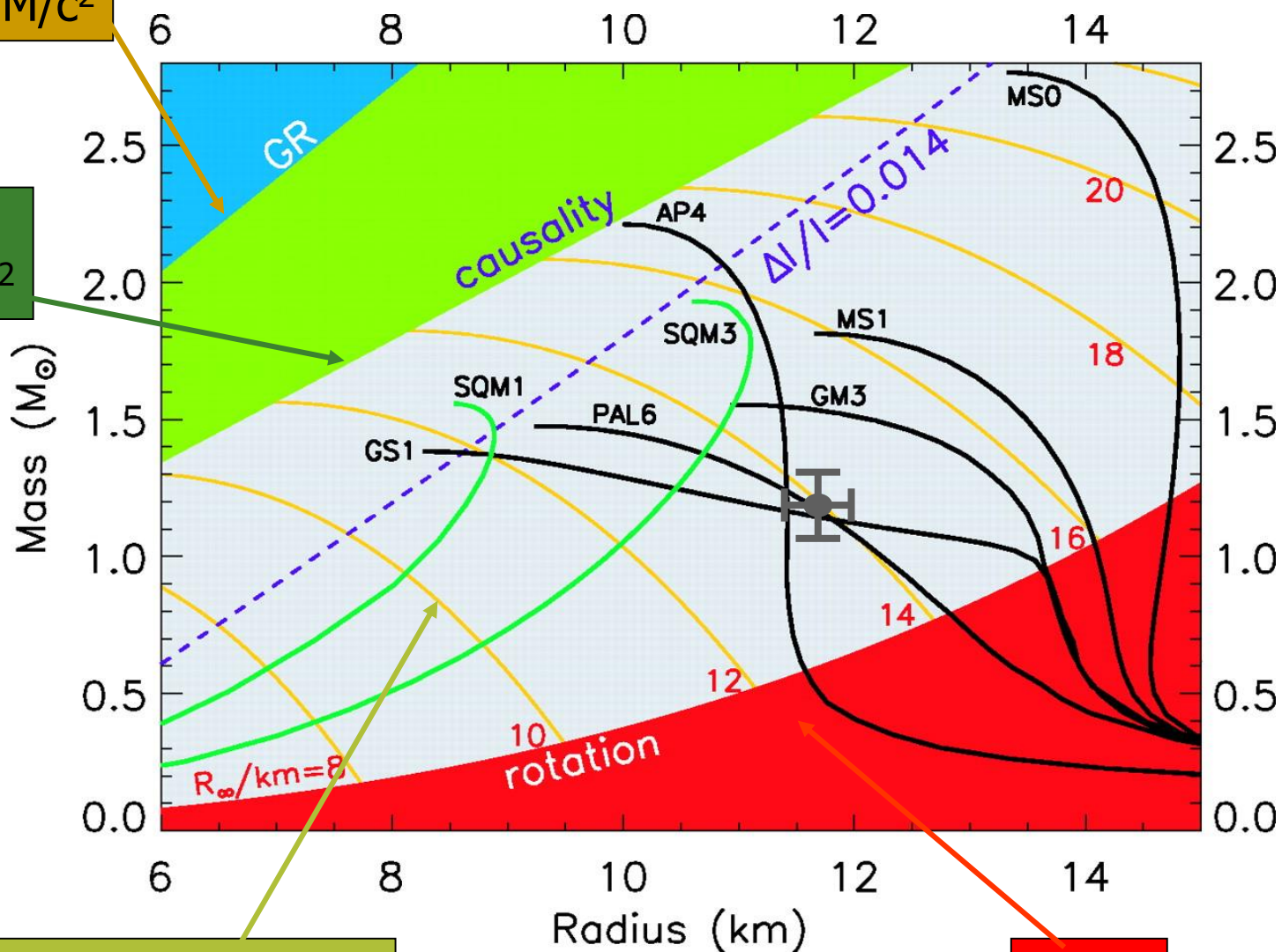
$$\frac{I}{MR^2} \equiv 0.207 + 0.857\beta \pm 0.011$$



$$R = 2GM/c^2$$

$$P = \rho$$

$$R \sim 3GM/c^2$$

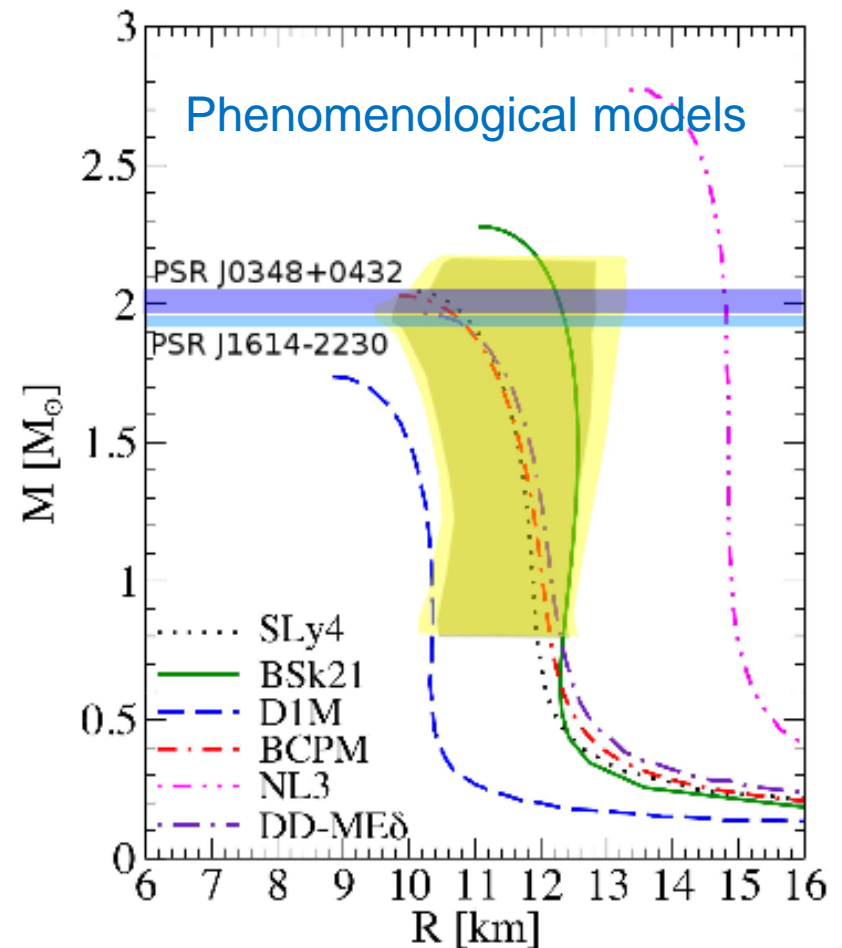
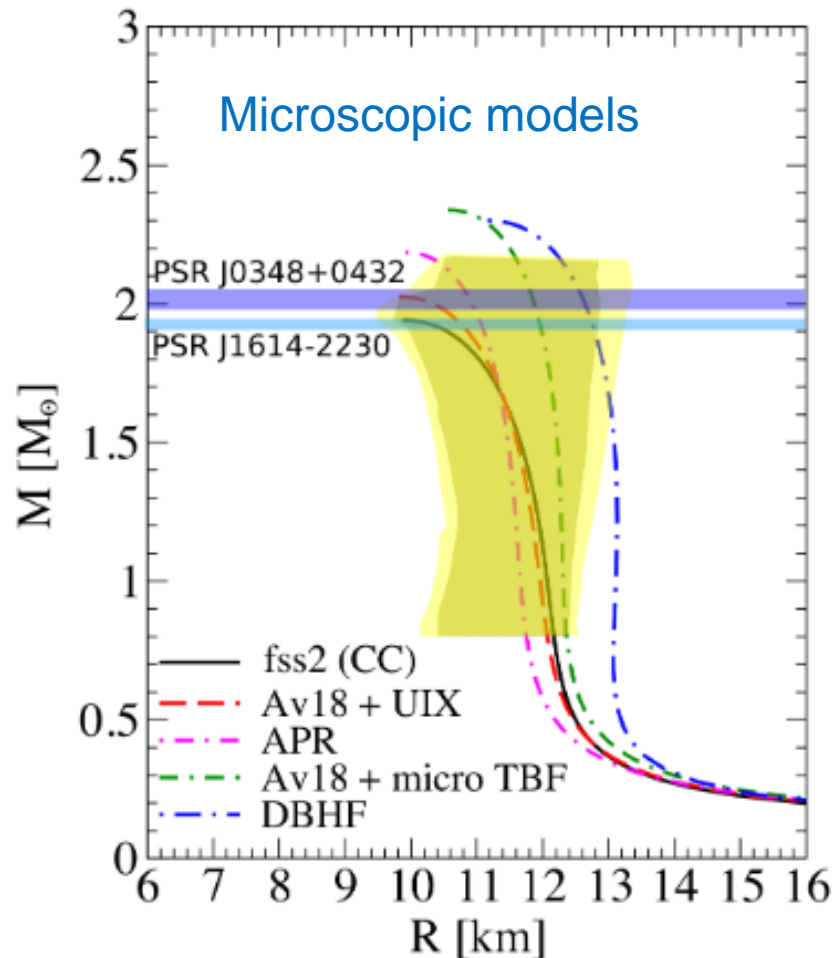


$$R_\infty = R(1 - 2GM/Rc^2)^{-1/2}$$

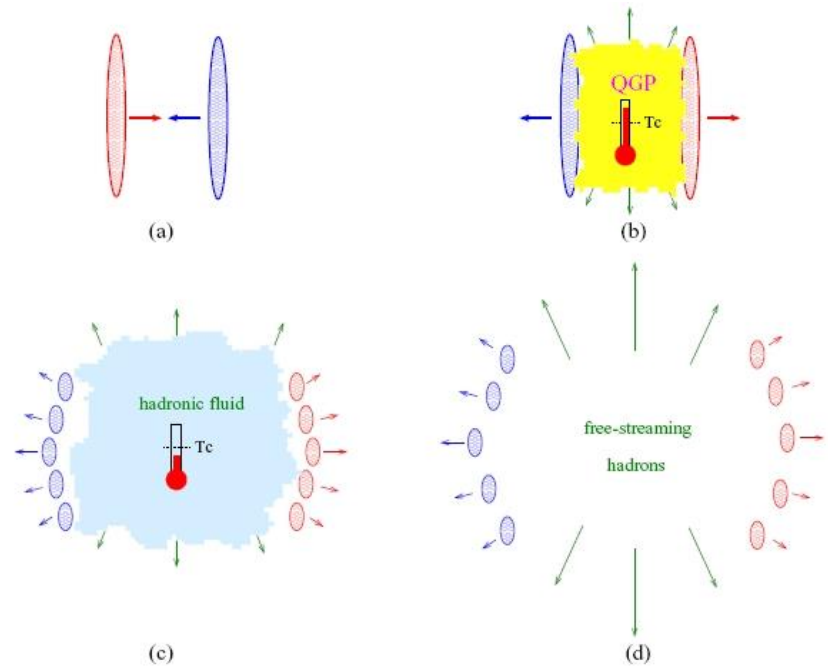
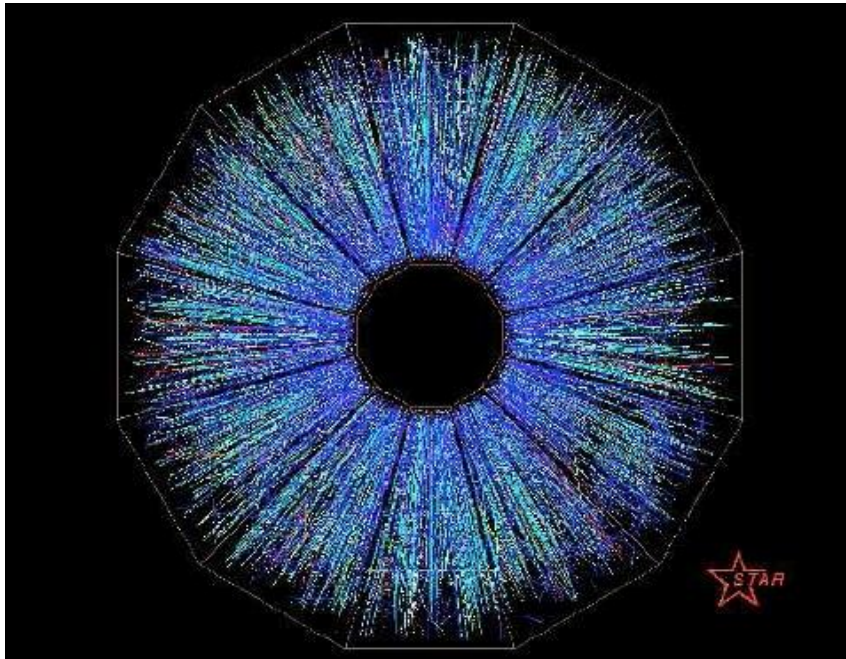
$$\omega = \omega_K$$

Lattimer & Prakash (2004)

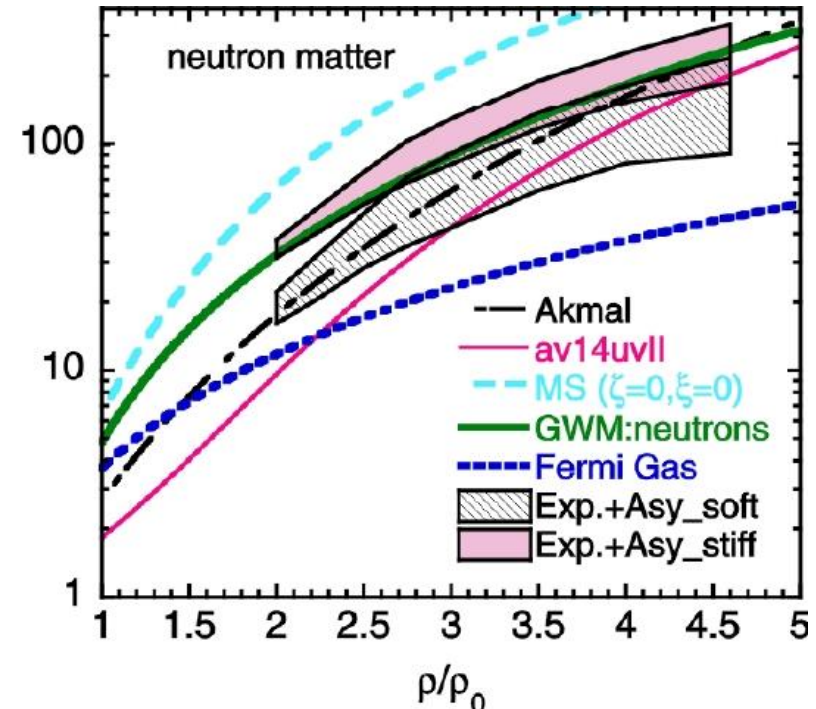
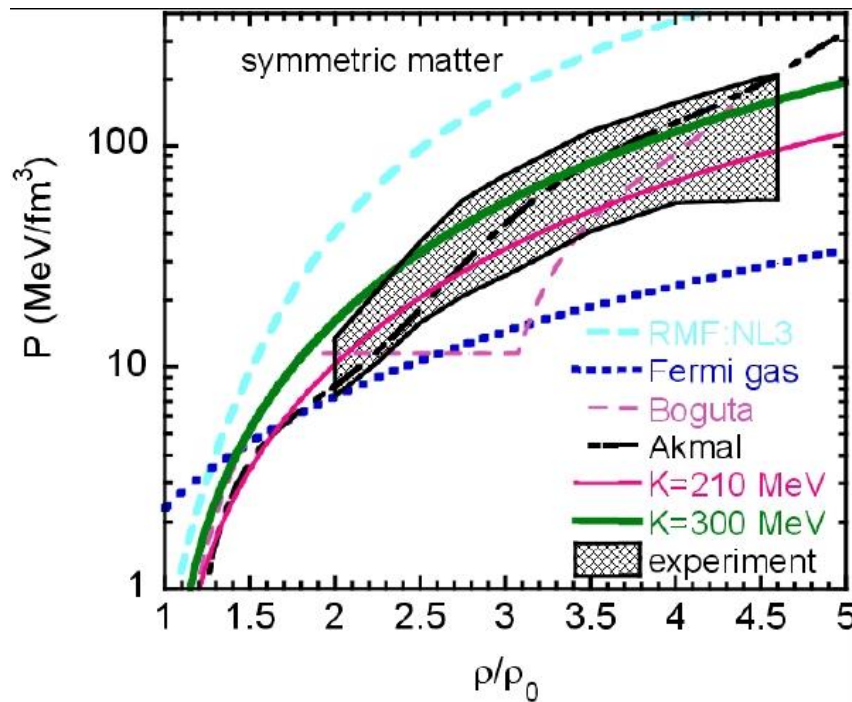
Theory vs. observations



Au-Au collisions



Experimental results and comparison



$$1 \text{ MeV/fm}^3 = 1.6 \cdot 10^{32} \text{ Pa}$$

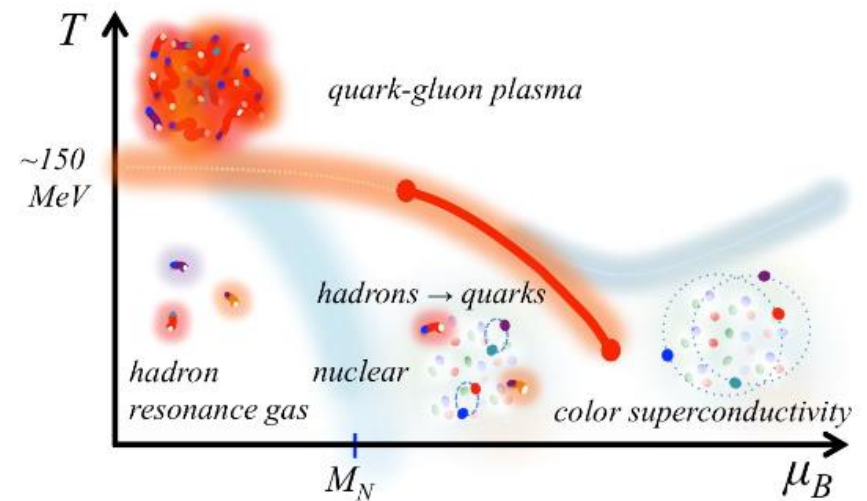
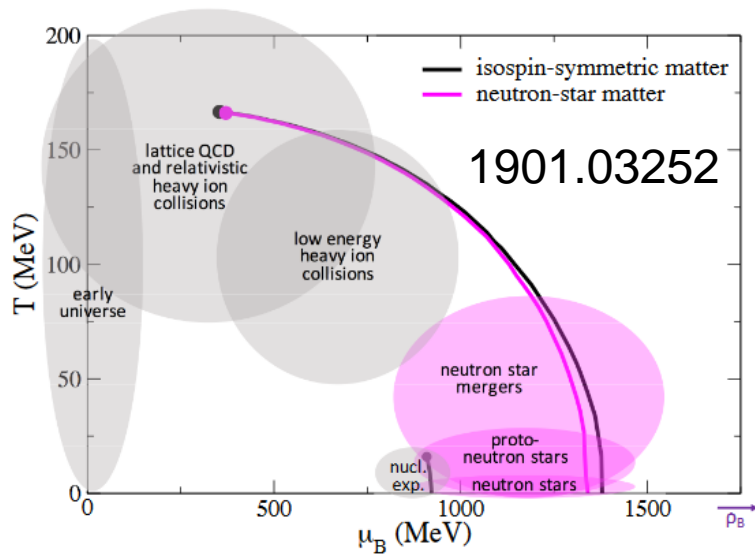
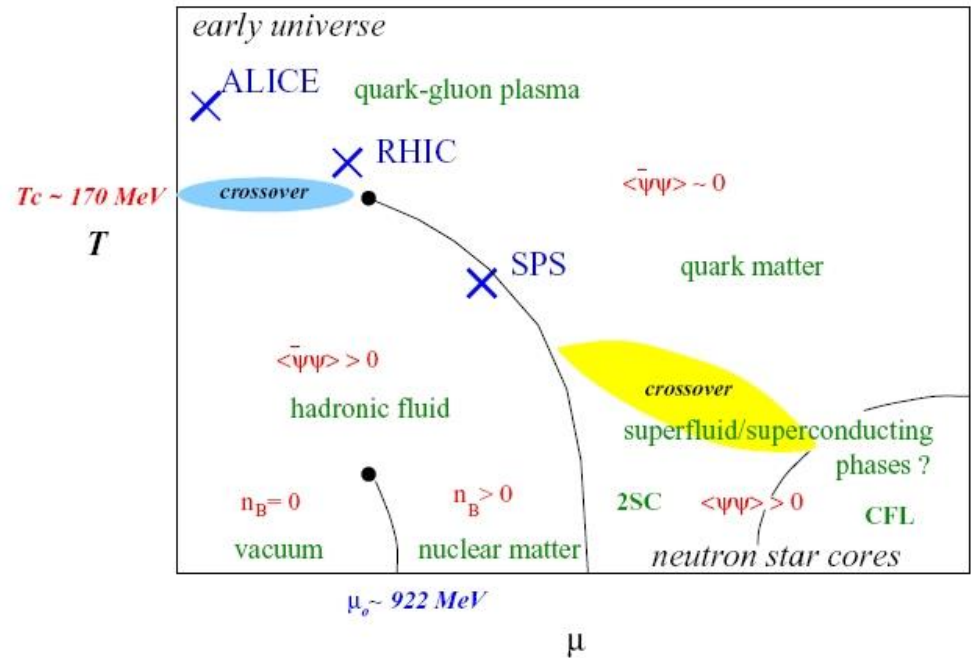
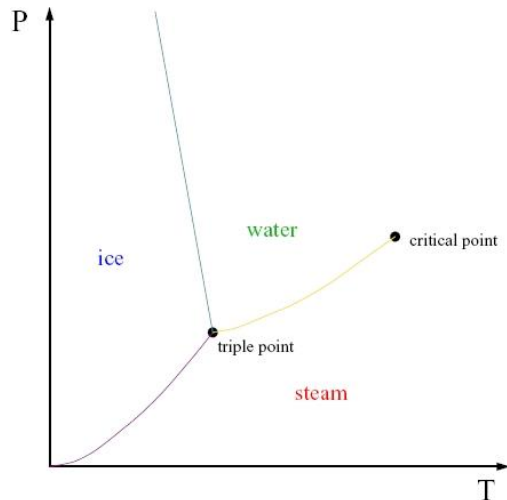
Danielewicz et al. nucl-th/0208016

GSI-SIS and AGS data

New heavy-ion data and discussion: 1211.0427

Also laboratory measurements of lead nuclei radius can be important, see 1202.5701

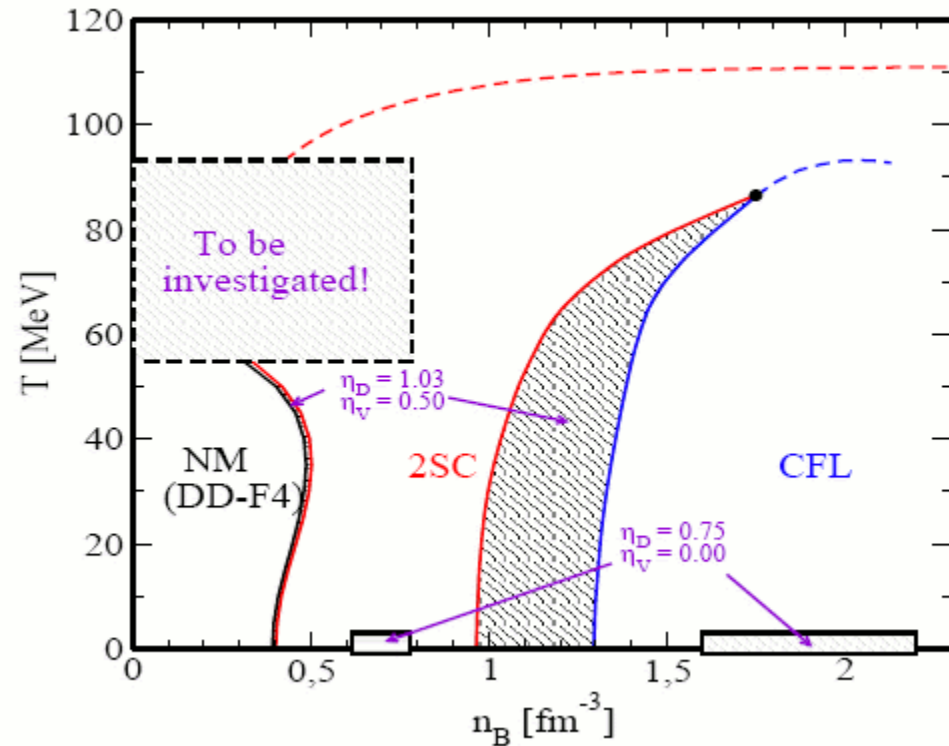
Phase diagram



See 1803.01836

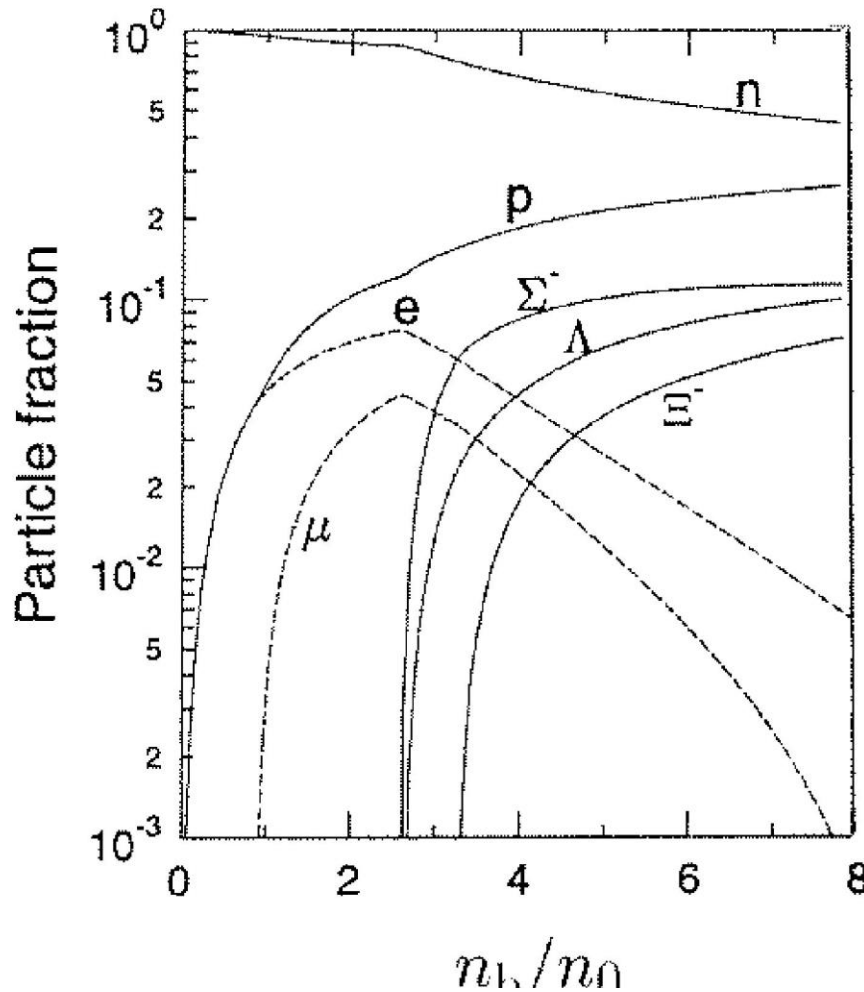
Phase diagram

Phase diagram for isospin symmetry using the most favorable hybrid EoS studied in astro-ph/0611595.

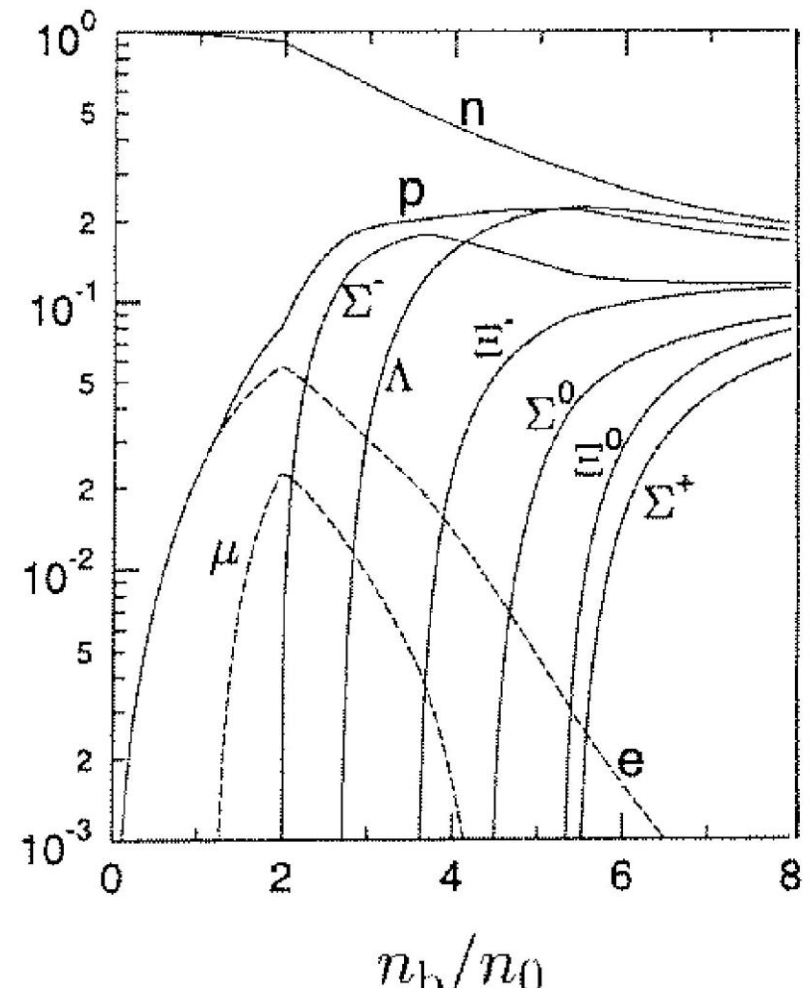


(astro-ph/0611595)

Particle fractions

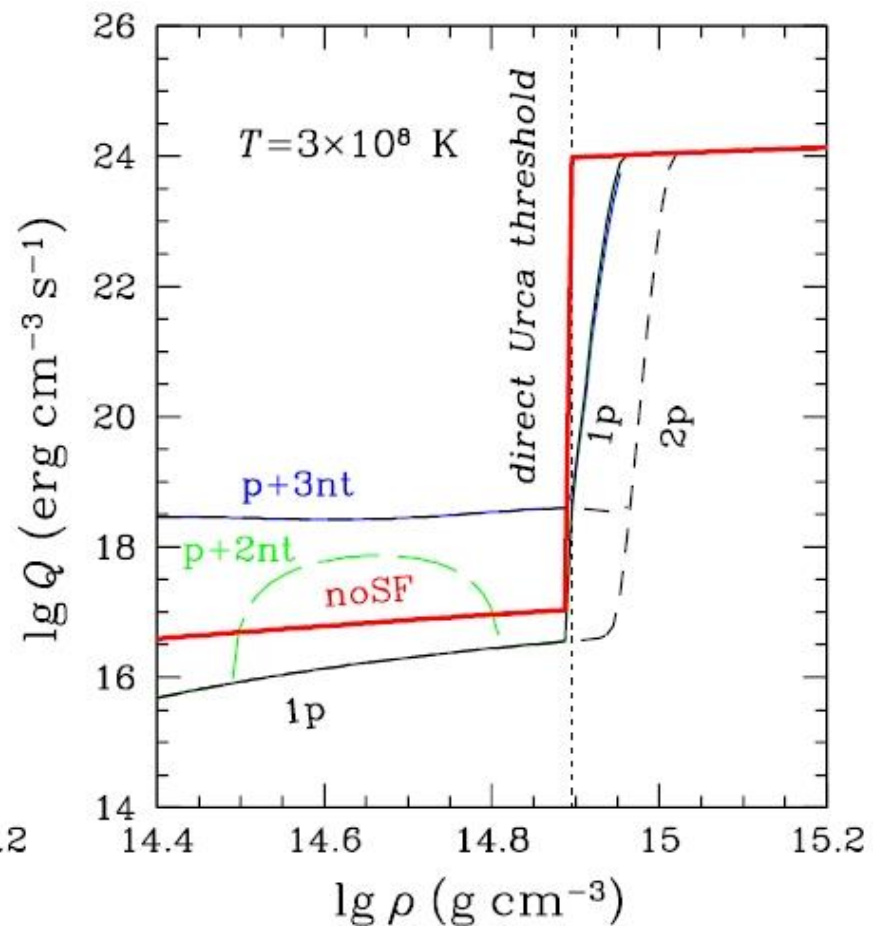
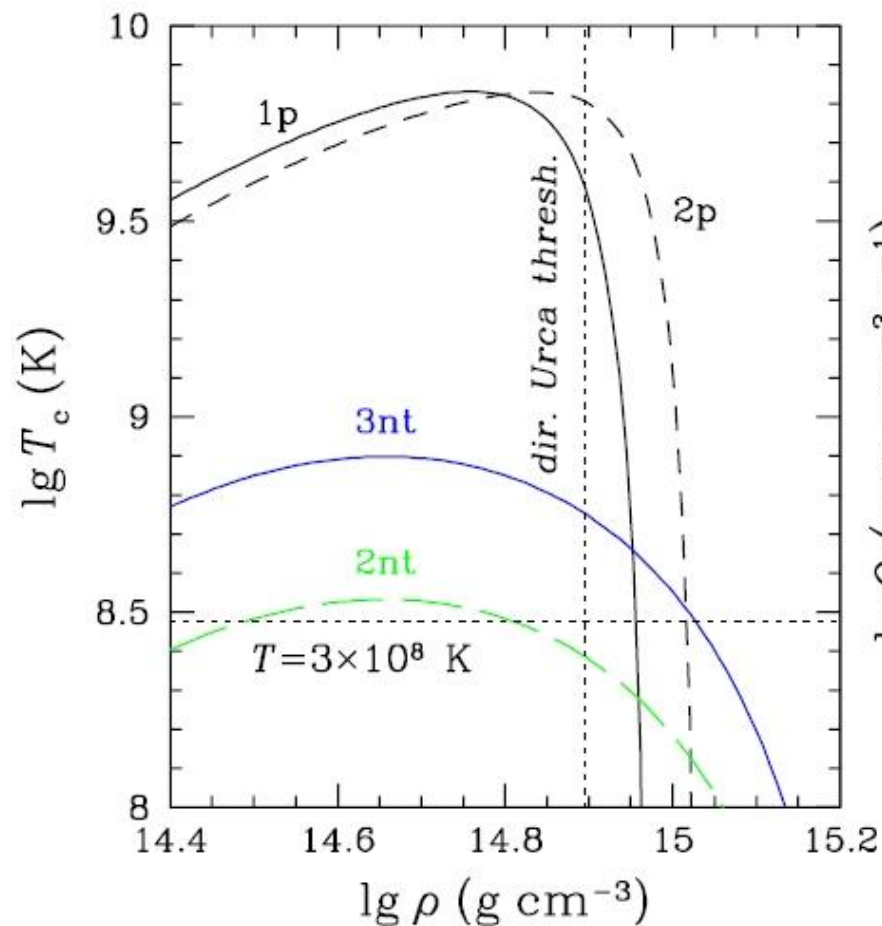


Effective chiral model of
Hanauske et al. (2000)



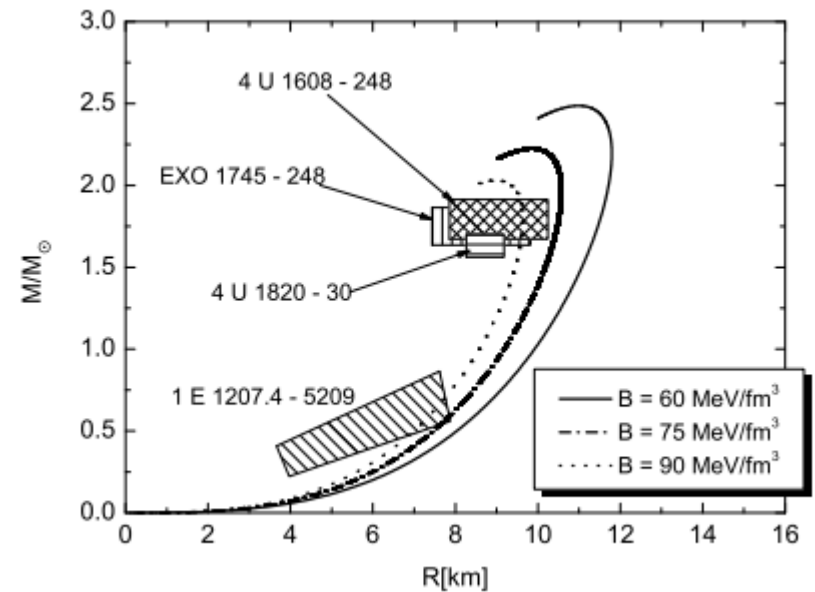
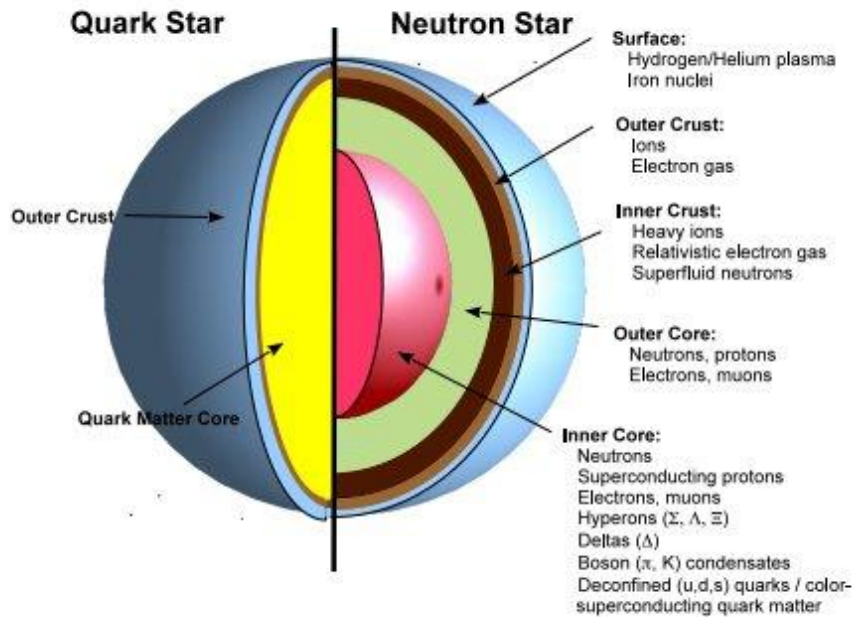
Relativistic mean-field model
TM1 of Sugahara & Toki (1971)

Superfluidity in NSs



(Yakovlev)

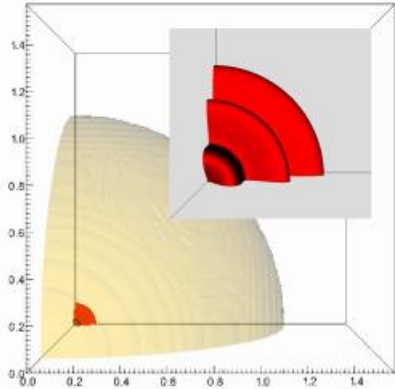
Quark stars



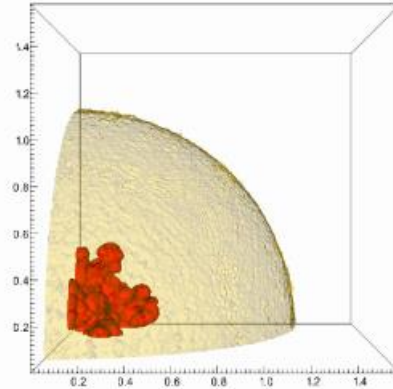
1210.1910

See also 1112.6430

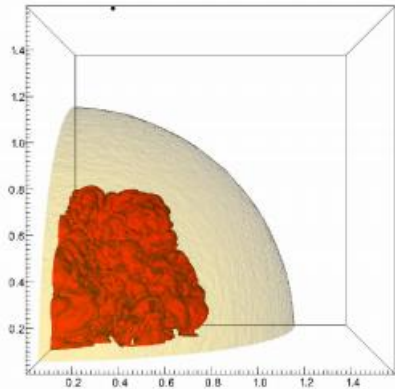
Formation of quark stars



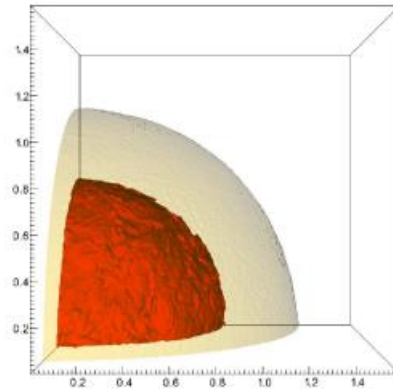
(a) $t = 0$



(b) $t = 0.7 \text{ ms}$



(c) $t = 1.2 \text{ ms}$

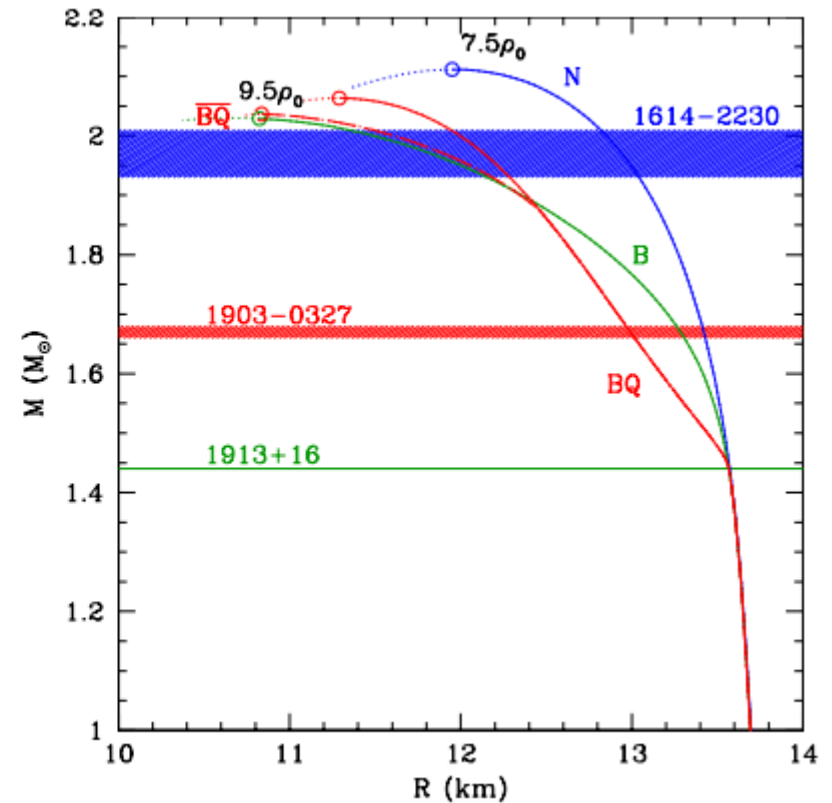
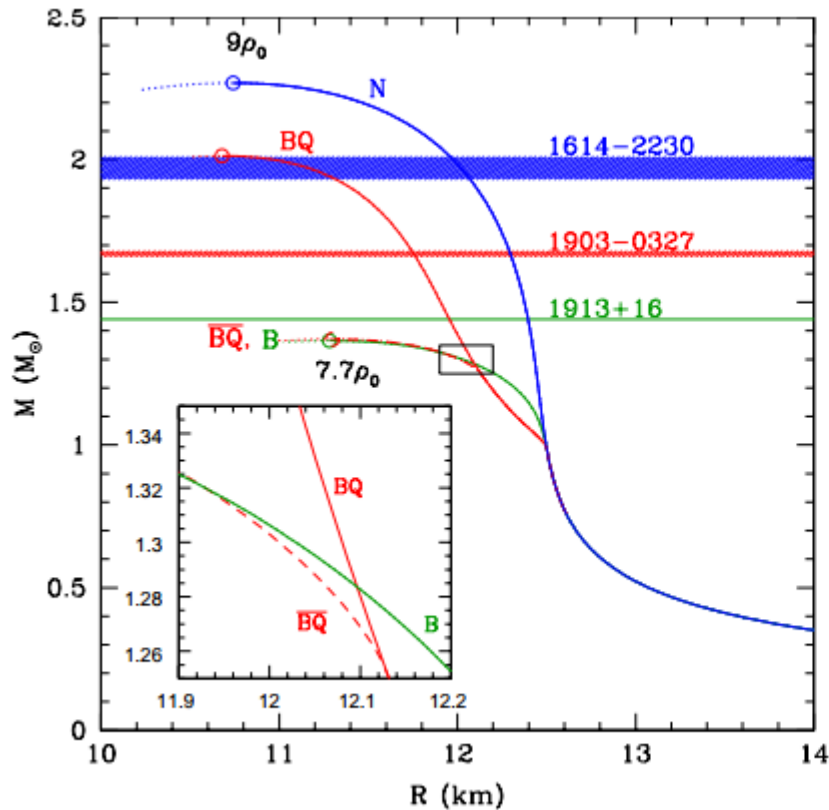


(d) $t = 4.0 \text{ ms}$

Turbulent deflagration,
as in SNIa.

Neutrino signal due to
conversion of a NS into
a quark star was calculated
in 1304.6884

Hybrid stars

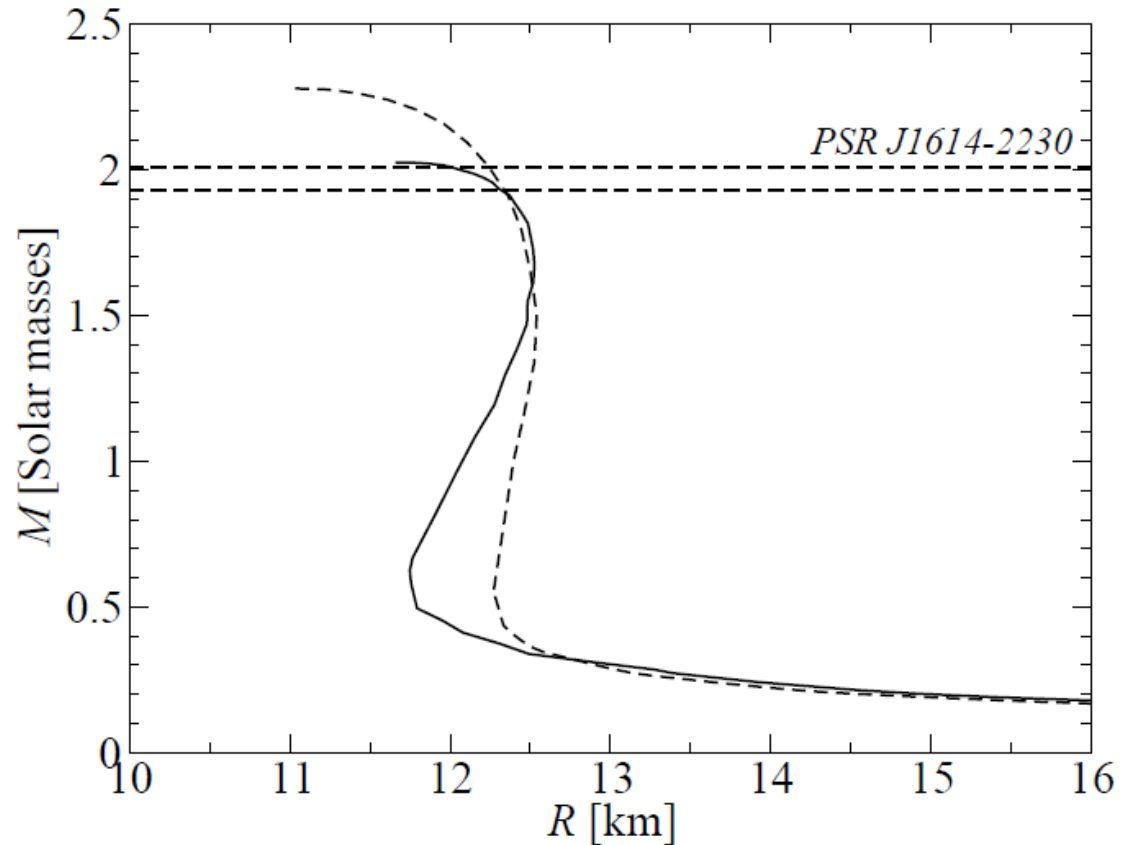


Stars with quark cores are reviewed in 1904.05471

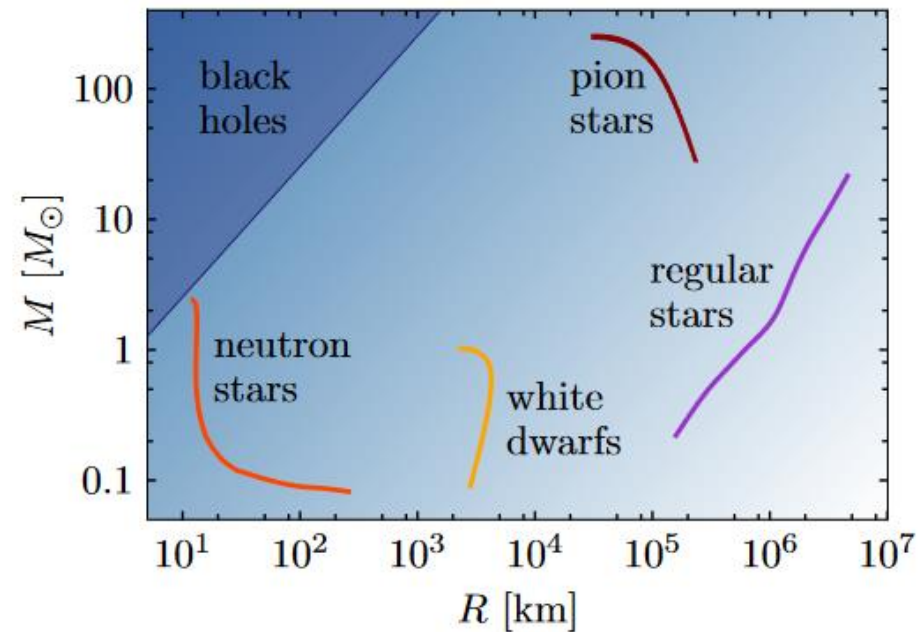
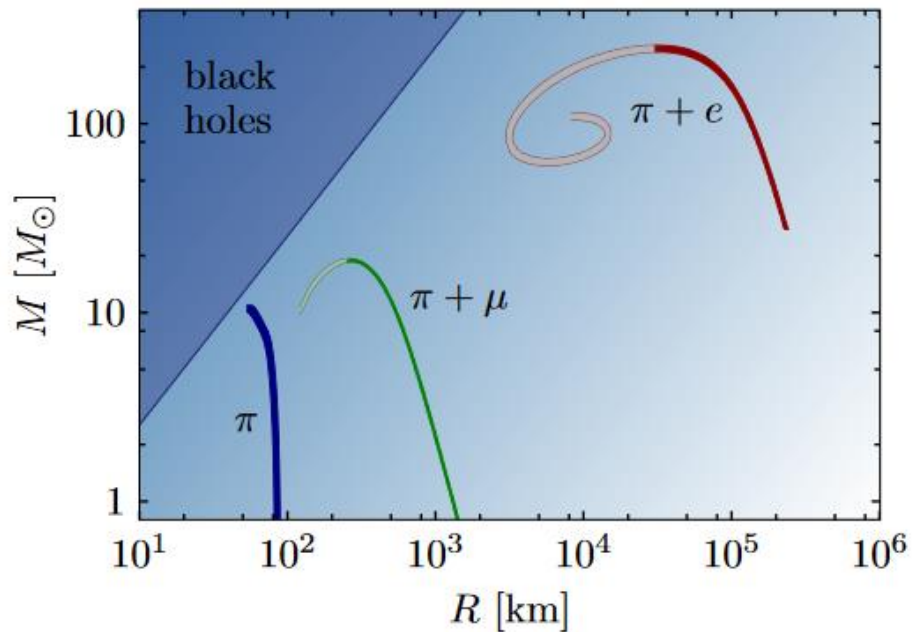
See also [1302.4732](#), [1903.08963](#), [1903.09121](#)

Massive hybrid stars

Stars with quark cores can be massive, and so this hypothesis is compatible with existence of pulsars with $M > 2$ Msolar



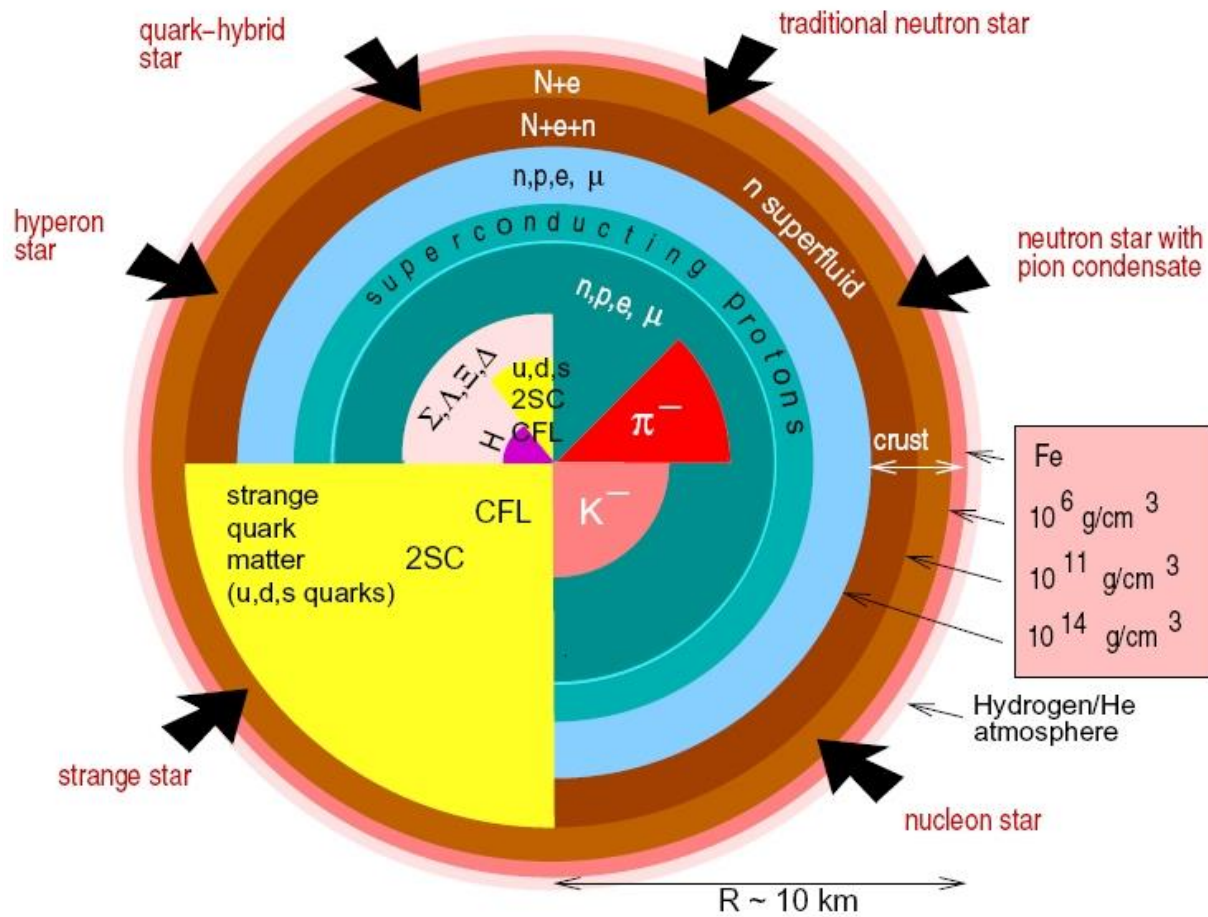
Pion stars



New exotic solution.

It is not clear if it can be applied to any known type of sources.

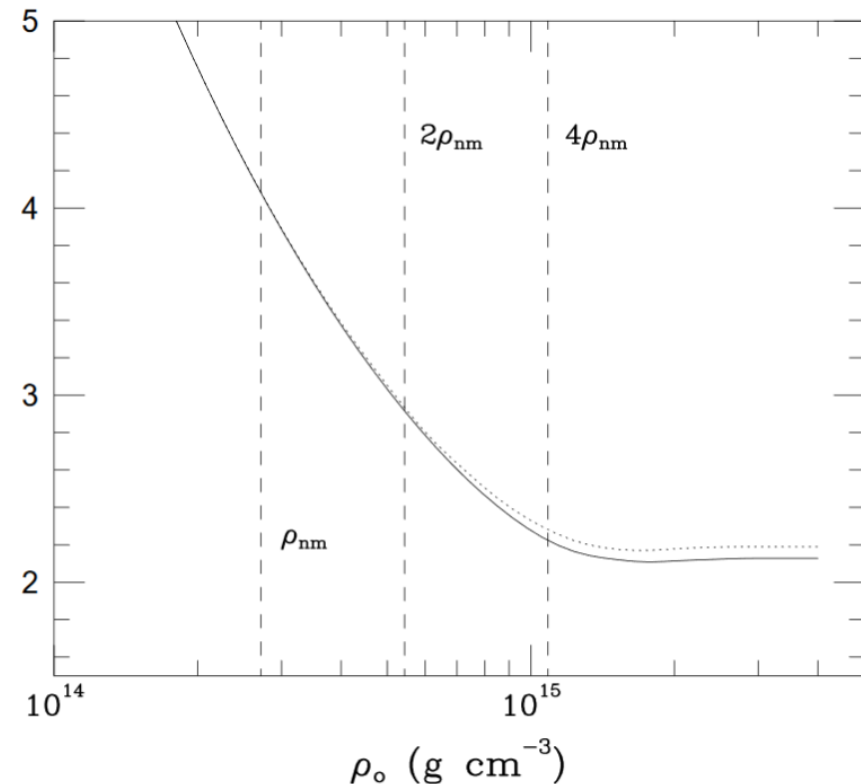
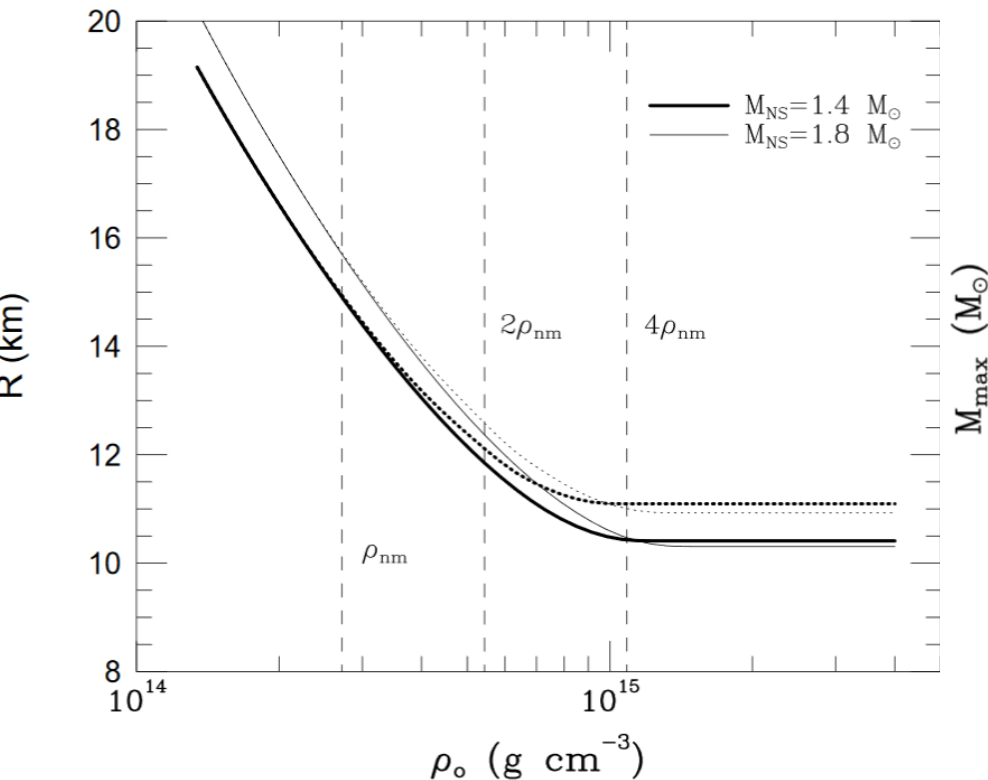
NS interiors: resume



(Weber et al. ArXiv: 0705.2708)

Maximum mass

Maximum mass of NSs depends on the EoS, however, it is possible to make calculations on the base of some fundamental assumptions.



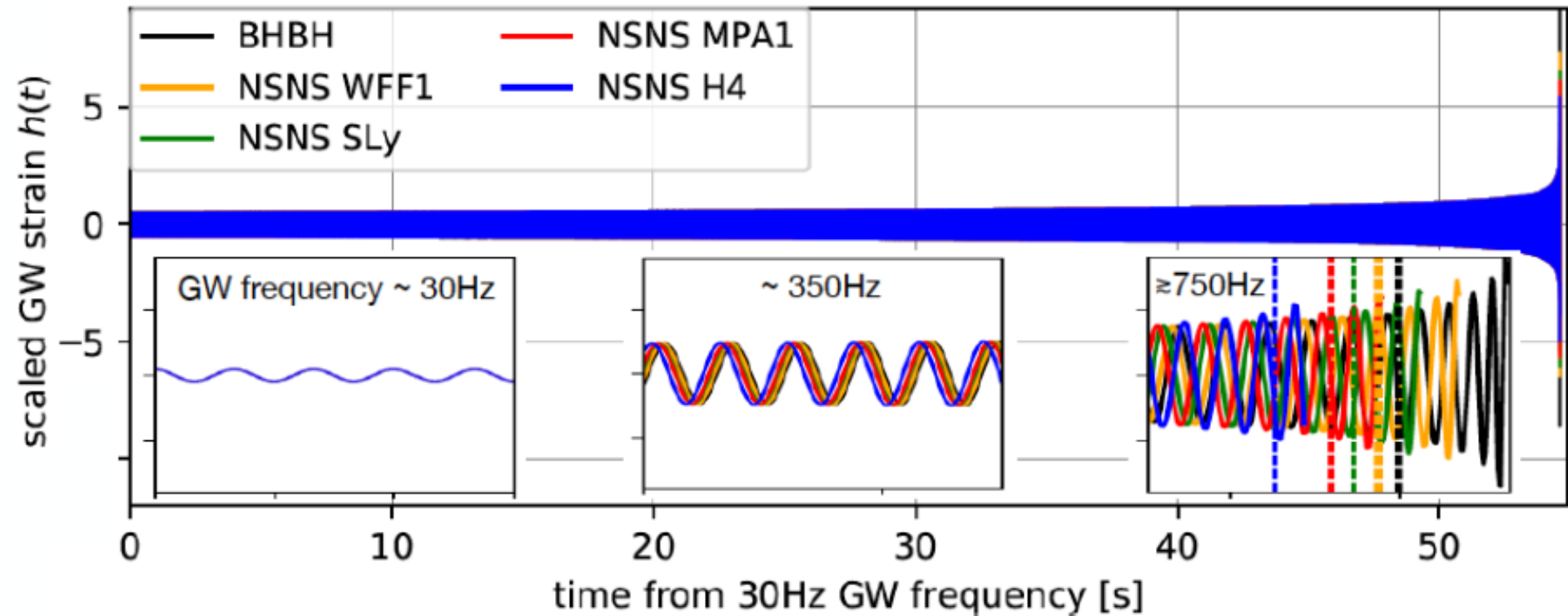
astro-ph/9608059

Seminal paper: Rhoades, Ruffini 1974

http://prl.aps.org/abstract/PRL/v32/i6/p324_1

$$c_s^2 = \frac{dP}{d\rho} = c^2.$$

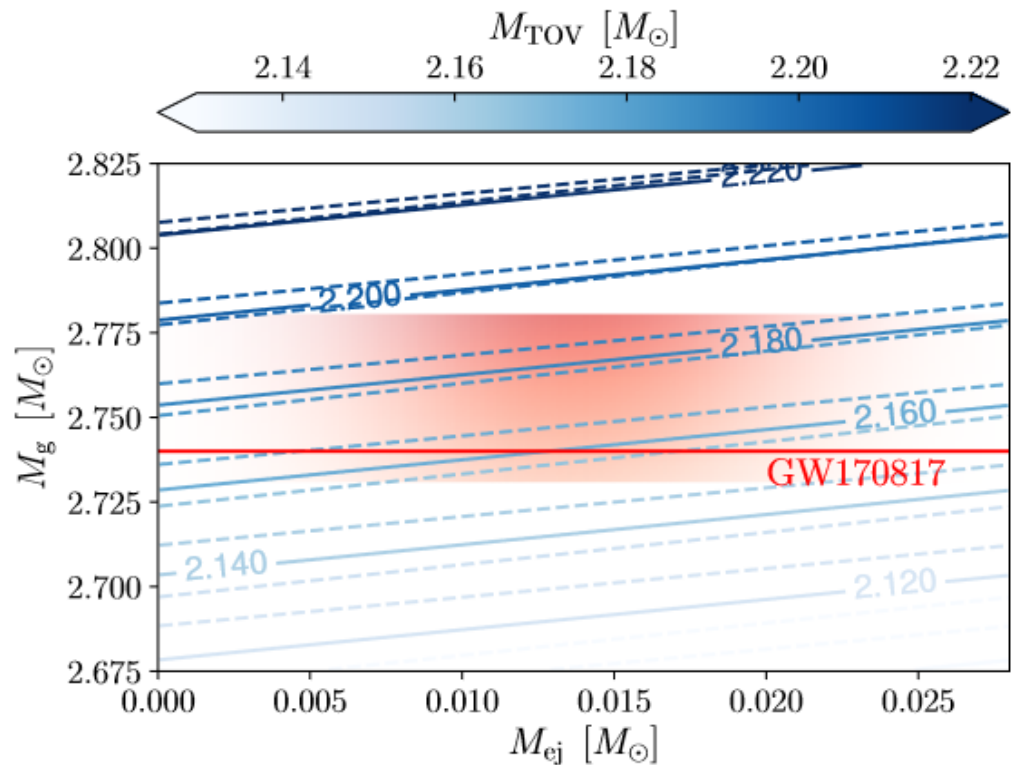
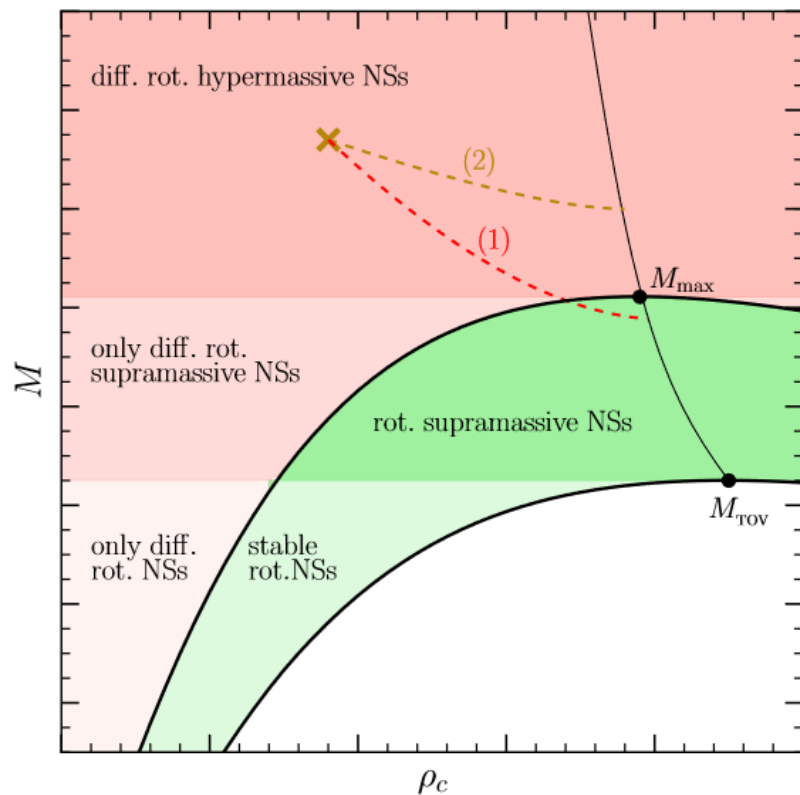
NS-NS coalescence and EoS



GW signals for different EoS

Calculations based on recent data on NS-NS coalescence

What uniform rotation can give: $M_{\max} = (1.20^{+0.02}_{-0.02}) M_{\text{TOV}}$ independently of the EOS



Another constraint from GW170817

$$M_{\text{NSNS}} \approx 2.74 \lesssim M_{\text{thresh}} \approx \alpha M_{\text{max}}^{\text{sph}}. \quad \longleftarrow \text{As there was no prompt collapse}$$

Here $\alpha \approx 1.3 - 1.7$ is the ratio of the HMNS threshold mass limit to the NS spherical maximum mass as gleaned from multiple numerical experiments of merging NSNSs

$$M_{\text{NSNS}} \approx 2.74 \gtrsim M_{\text{max}}^{\text{sup}} \approx \beta M_{\text{max}}^{\text{sph}},$$

where $\beta \approx 1.2$ is the ratio of the uniformly rotating supra-massive NS limit to the nonrotating spherical maximum

$$M_{\text{max}}^{\text{sph}} = 4.8 \left(\frac{2 \times 10^{14} \text{ gr/cm}^3}{\rho_m/c^2} \right)^{1/2} M_{\odot},$$

$$M_{\text{max}}^{\text{sup}} = 6.1 \left(\frac{2 \times 10^{14} \text{ gr/cm}^3}{\rho_m/c^2} \right)^{1/2} M_{\odot},$$

$$\beta \approx 1.27.$$

$$2.74/\alpha \lesssim M_{\text{max}}^{\text{sph}} \lesssim 2.74/\beta$$

$$M_{\text{max}}^{\text{sph}} \lesssim 2.16. \quad \beta \approx 1.27.$$

$$M_{\text{max}}^{\text{sph}} \lesssim 2.28. \quad \beta = 1.2$$

Papers to read

1. astro-ph/0405262 Lattimer, Prakash "Physics of neutron stars"
2. 0705.2708 Weber et al. "Neutron stars interiors and equation of state ..."
3. physics/0503245 Baym, Lamb "Neutron stars"
4. 0901.4475 Piekarewicz "Nuclear physics of neutron stars" (first part)
5. 0904.0435 Paerels et al. "The Behavior of Matter Under Extreme Conditions"
6. 1512.07820 Lattimer, Prakash "The EoS of hot dense matter"
7. 1001.3294 Schmitt "Dense matter in compact stars - A pedagogical introduction "
8. 1303.4662 Hebeler et al. "EoS and NS properties vs. nuclear phys. and observation"
9. 1210.1910 Weber et al. Structure of quark star
10. 1302.1928 Stone "High density matter "
11. 1707.04966 Baym et al. "From hadrons to quarks in neutron stars: a review"
12. 1804.03020. Burgio, Fantina "Nuclear EoS for Compact Stars and Supernovae"
13. 1803.01836 Blaschke, Chamel. "Phases of dense matter in compact stars"
14. 1904.05471 Alford et al. "Signatures of quark matter..."
15. 1904.08907 Coleman Miller et al. "Constraining the EoS"
16. 1912.11876 Pethick "Dense matter and neutron stars"

+ the book by Haensel, Yakovlev, Potekhin

Lectures on the Web

Lectures can be found at my homepage:

<http://xray.sai.msu.ru/~polar/html/presentations.html>