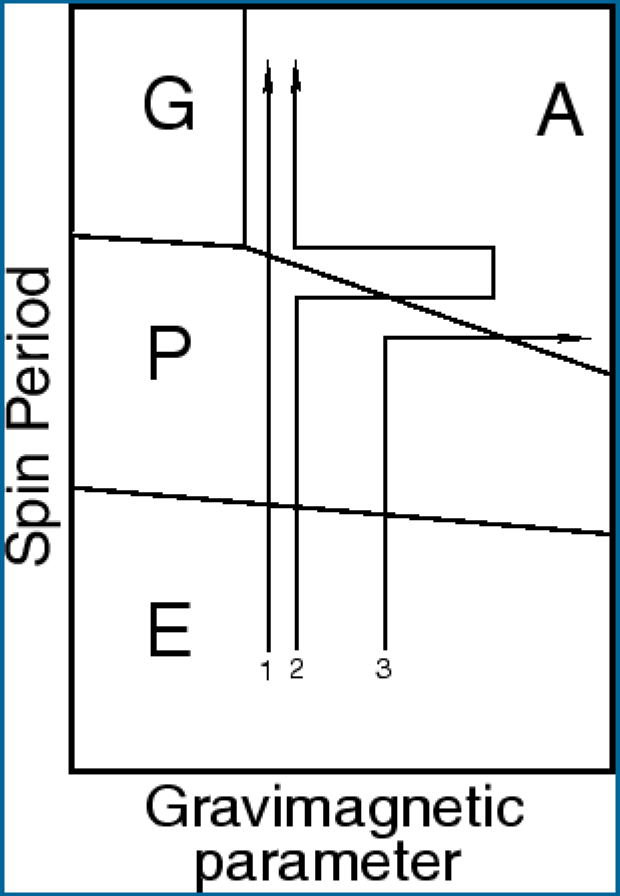
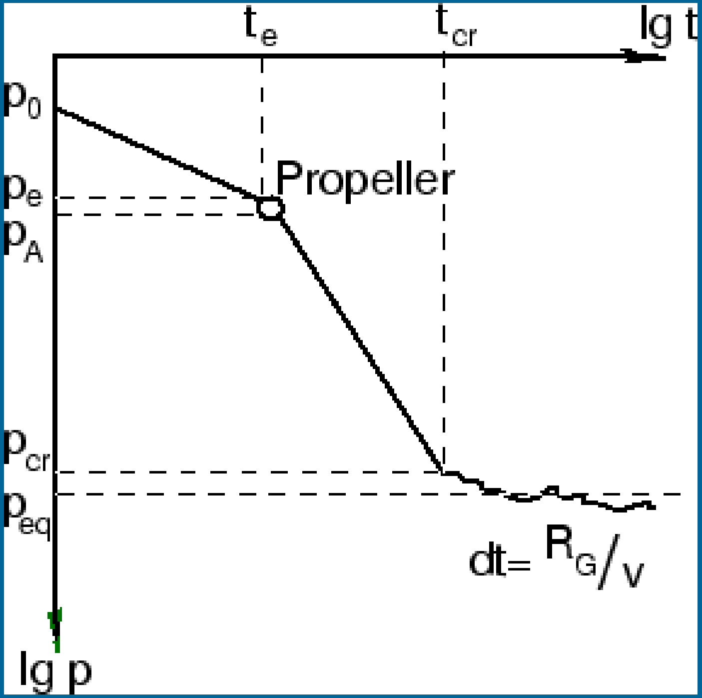
Thermal evolution of neutron stars

Evolution of neutron stars. I.: rotation + magnetic field

Ejector → Propeller → Accretor → Georotator



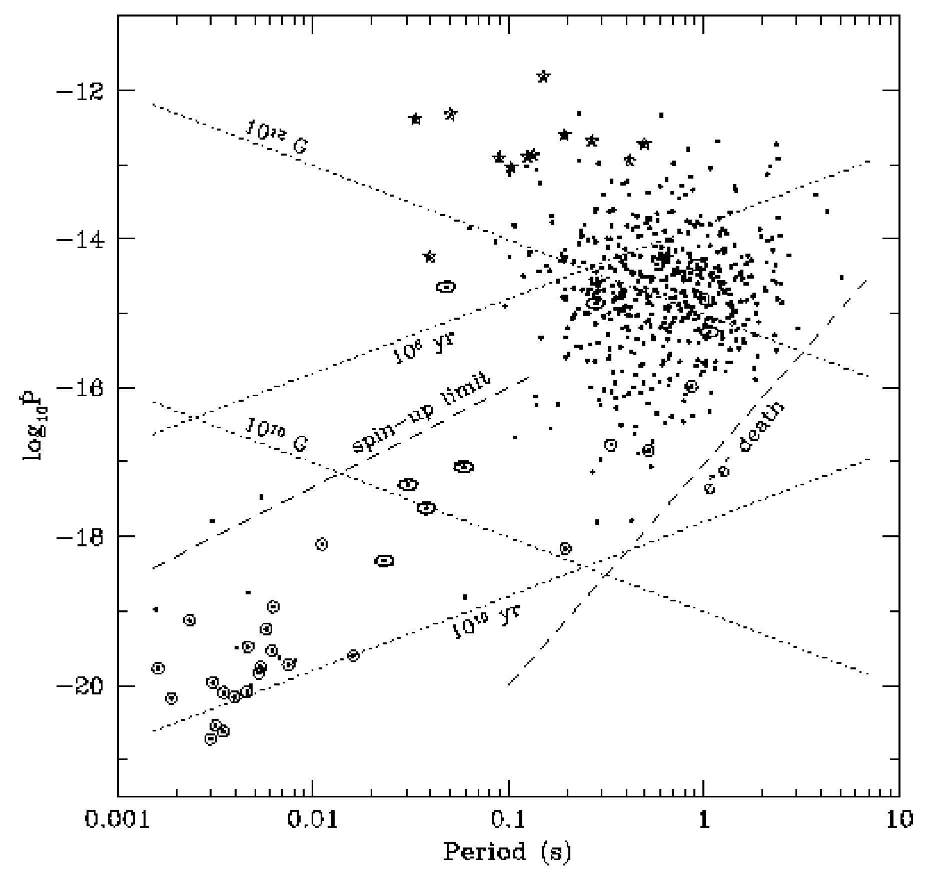
1 – spin down
2 – passage through a molecular cloud
3 – magnetic field decay



astro-ph/0101031

See the book by Lipunov (1987, 1992)

Magnetorotational evolution of radio pulsars



$$L_m = \frac{2}{3}\frac{\mu^2\omega^4}{c^3}\sin^2\beta = \kappa_t\frac{\mu^2}{R_l^3}\omega\,,$$

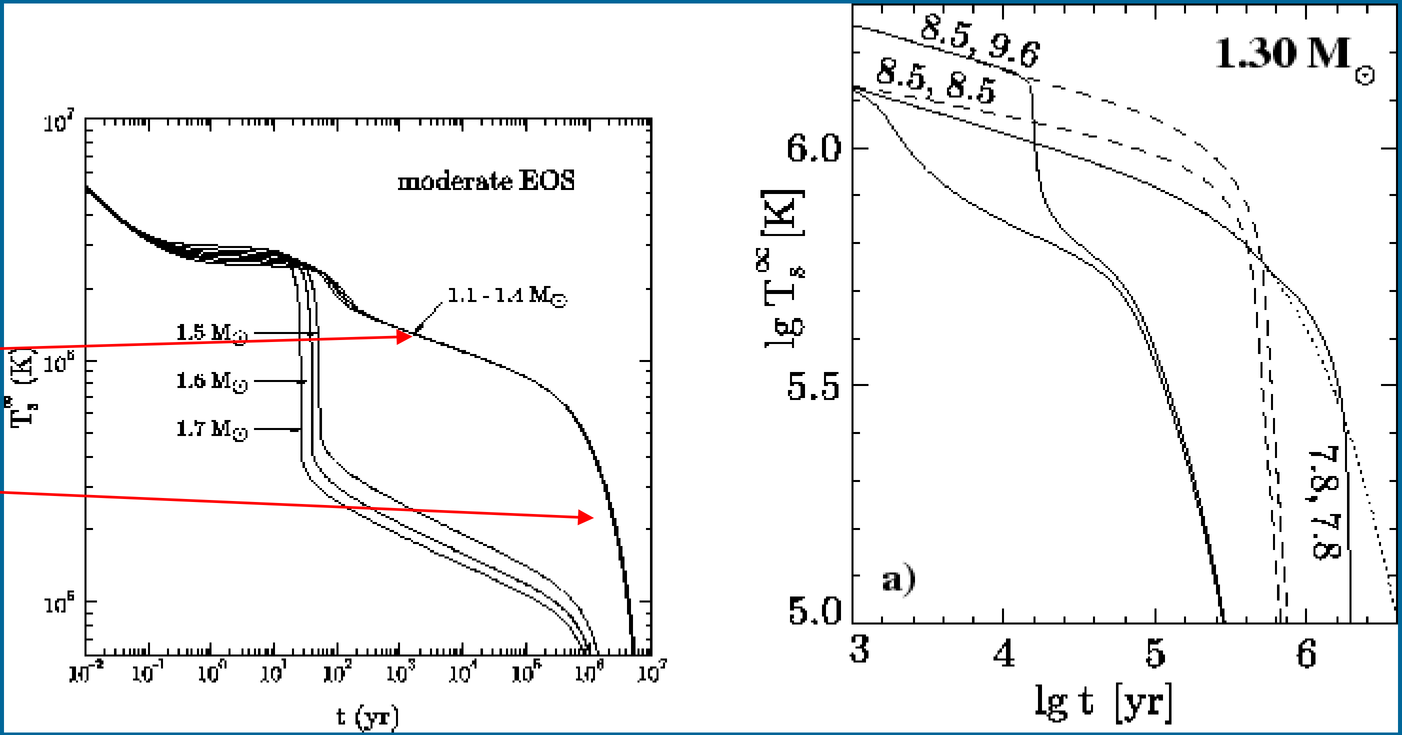
$$B \sim 3.2\times10^{19}\,(P\,dP/dt)^{1/2}\,\mathrm{G}.$$

Spin-down.
Rotational energy is released.
The exact mechanism is still unknown.

Evolution of NSs. II.: temperature

Neutrino
cooling stage

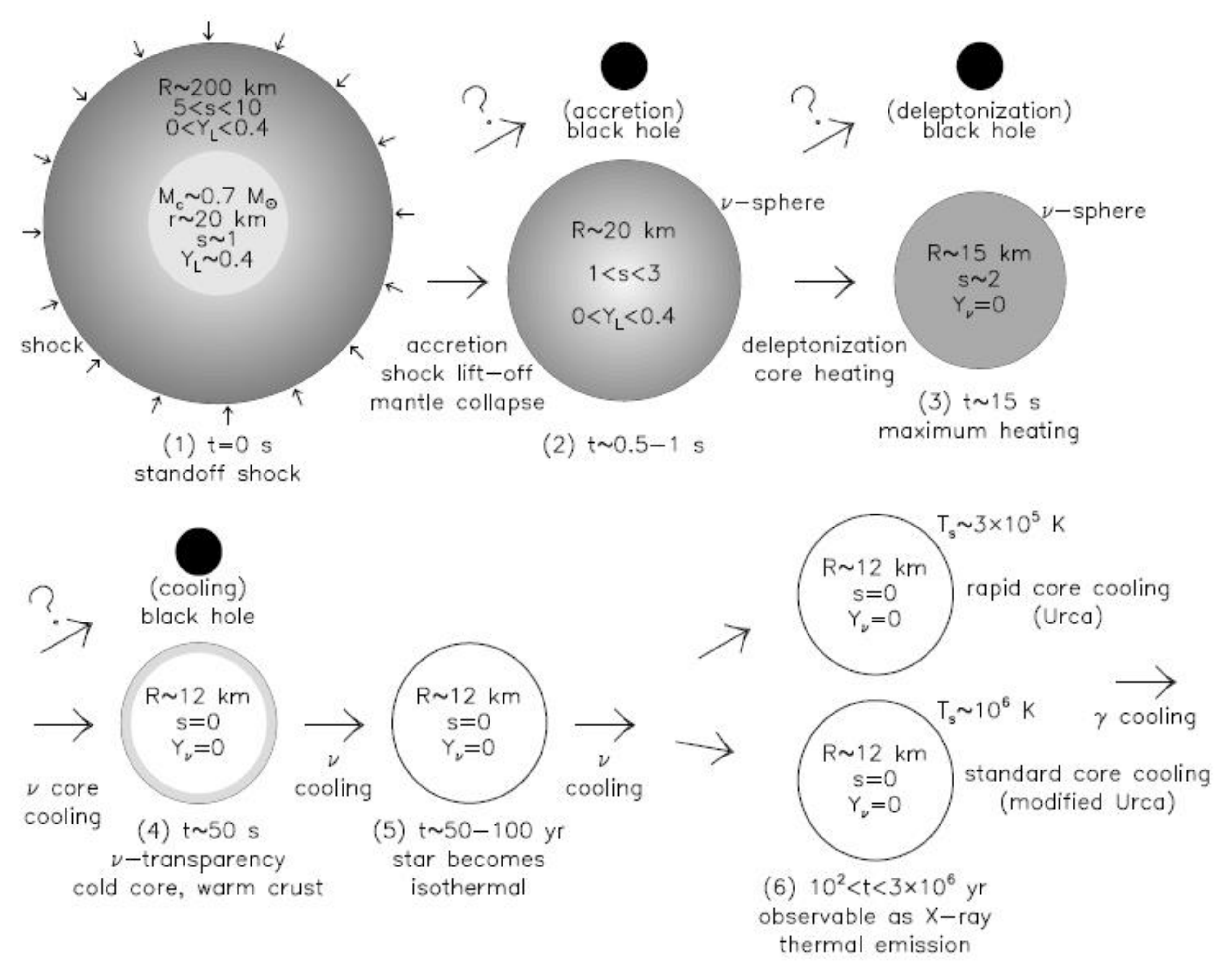
Photon
cooling stage



First papers on the thermal
evolution appeared already
in early 60s, i.e. before
the discovery of radio pulsars.

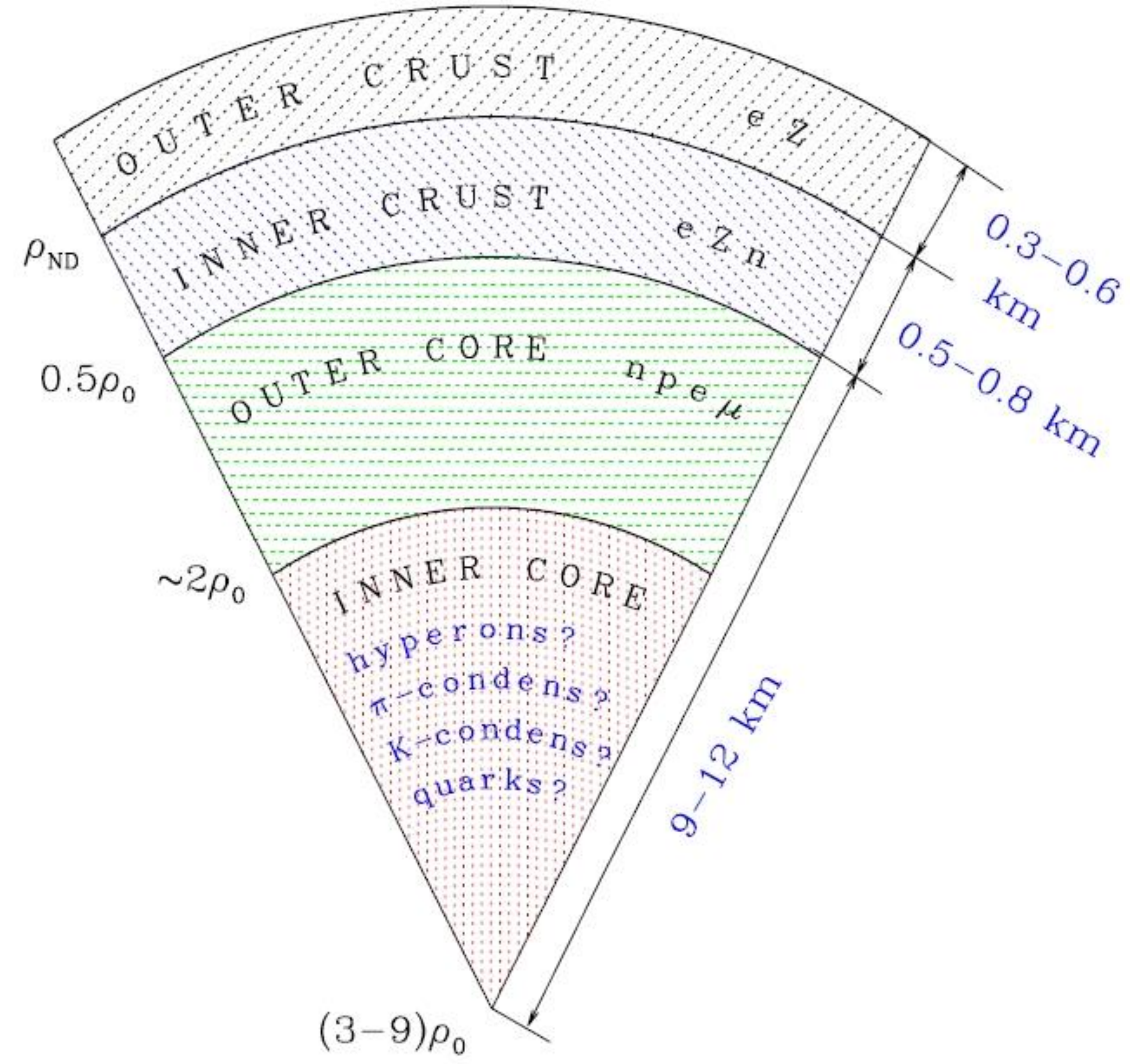
[Yakovlev et al. (1999) Physics Uspekhi]

Early evolution of a NS



(Prakash et al. astro-ph/0012136)

Structure and layers



Plus an atmosphere...

See Ch.6 in the book by Haensel, Potekhin, Yakovlev

$\rho_0 \sim 2.8\ 10^{14}$ g cm^{-3}

The total thermal energy of a nonsuperfluid neutron star is estimated as $U_T \sim 10^{48}\ T^2_9$ erg.

The heat capacity of an *npe* neutron star core with strongly superfluid neutrons and protons is determined by the electrons, which are not superfluid, and it is ~20 times lower than for a neutron star with a nonsuperfluid core.

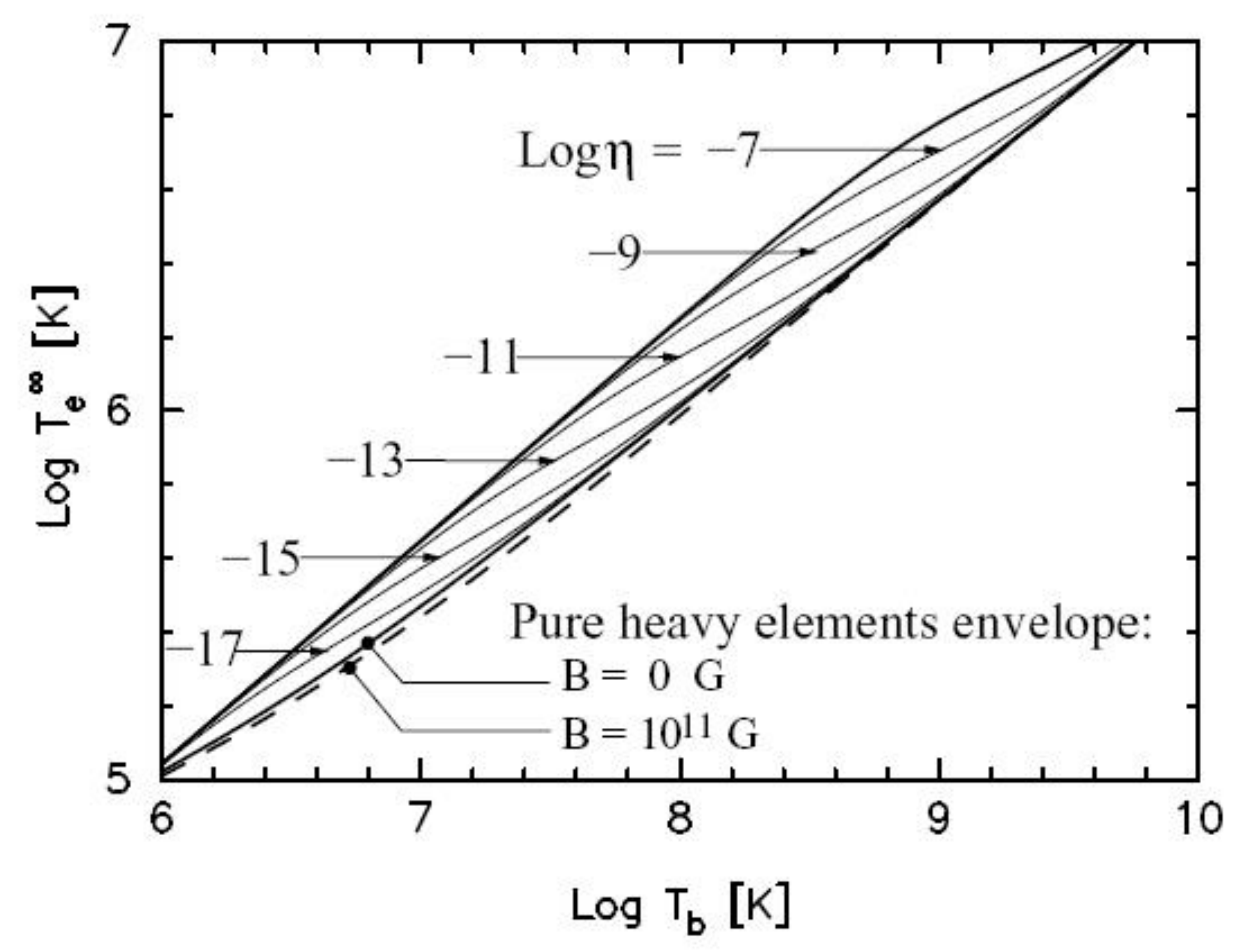
NS Cooling

- NSs are born very hot, T > 10^{10} K
- At early stages neutrino cooling dominates (exotic is possible – axions 1205.6940)
- The core is isothermal

$$\frac{dE_{th}}{dt} = C_V \frac{dT}{dt} = -\underbrace{L_\nu}_{\text{Neutrino luminosity}} - \overbrace{L_\gamma}^{\text{Photon luminosity}}$$

$$L_\gamma = 4\pi R^2 \sigma T_s^4, \quad T_s \propto T^{1/2+\alpha} \quad (|\alpha| << 1)$$

Core-crust temperature relation

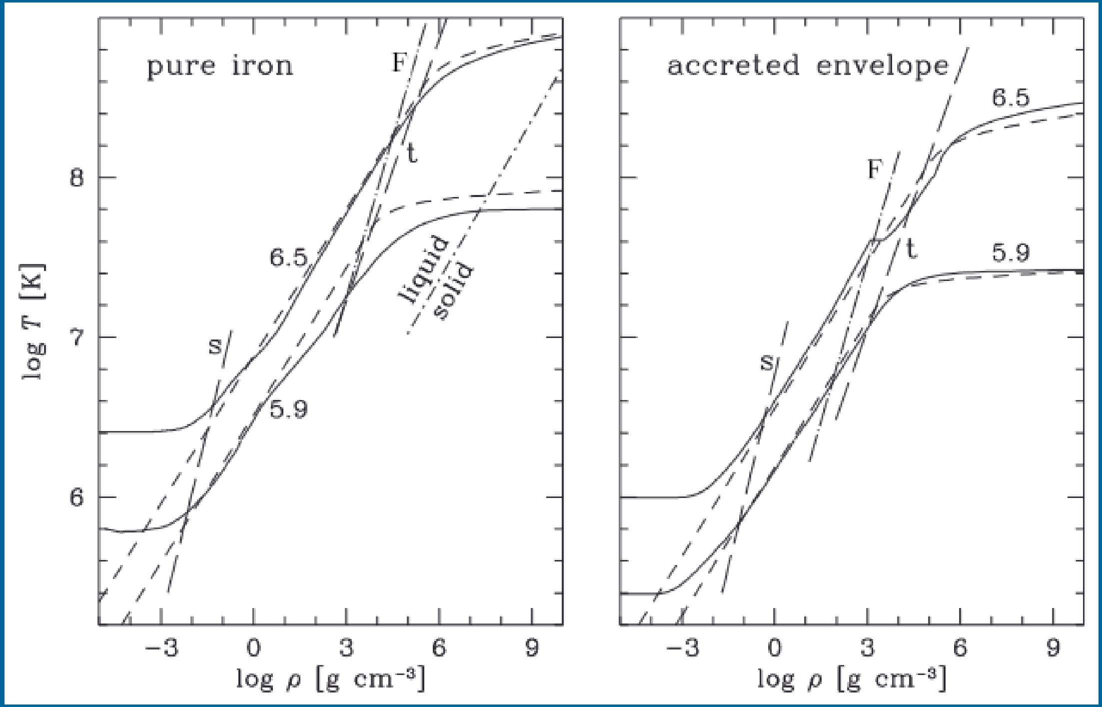


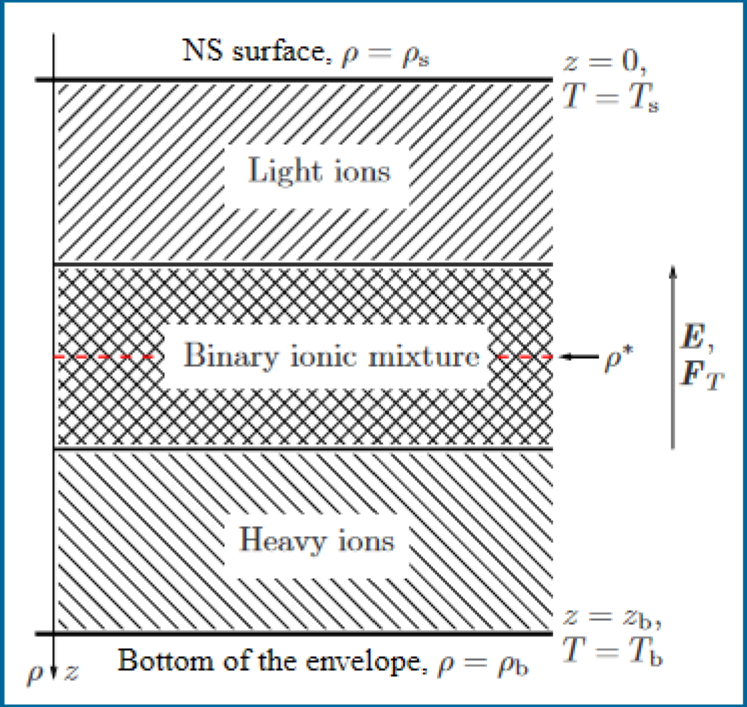
Heat blanketing envelope.
~100 meters
density ~10^{10} gcm^{-3}

See a review about crust properties related to thermal evolution in 1201.5602, 1507.06186, and in a recent large review – 2103.12422

Page et al. astro-ph/0508056

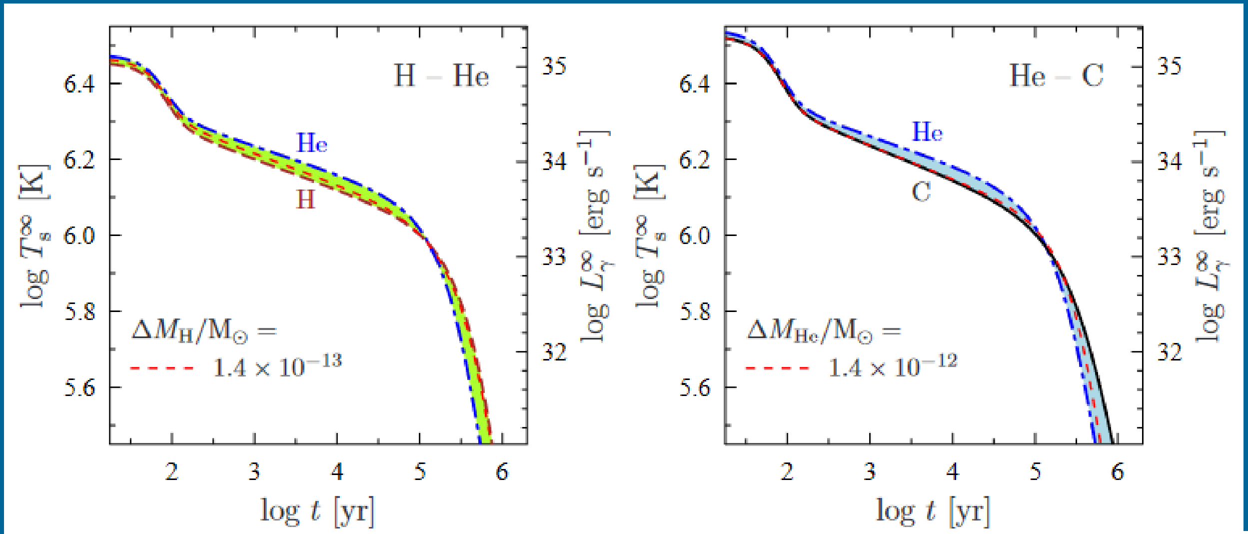
Temperature profile in an envelope

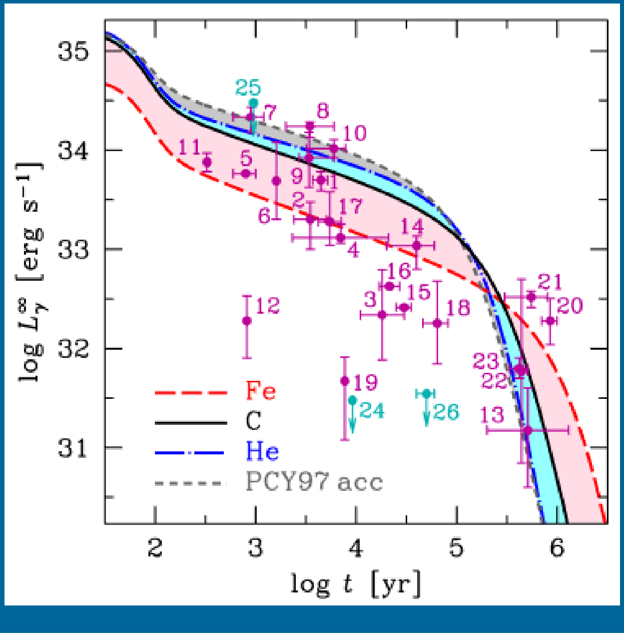


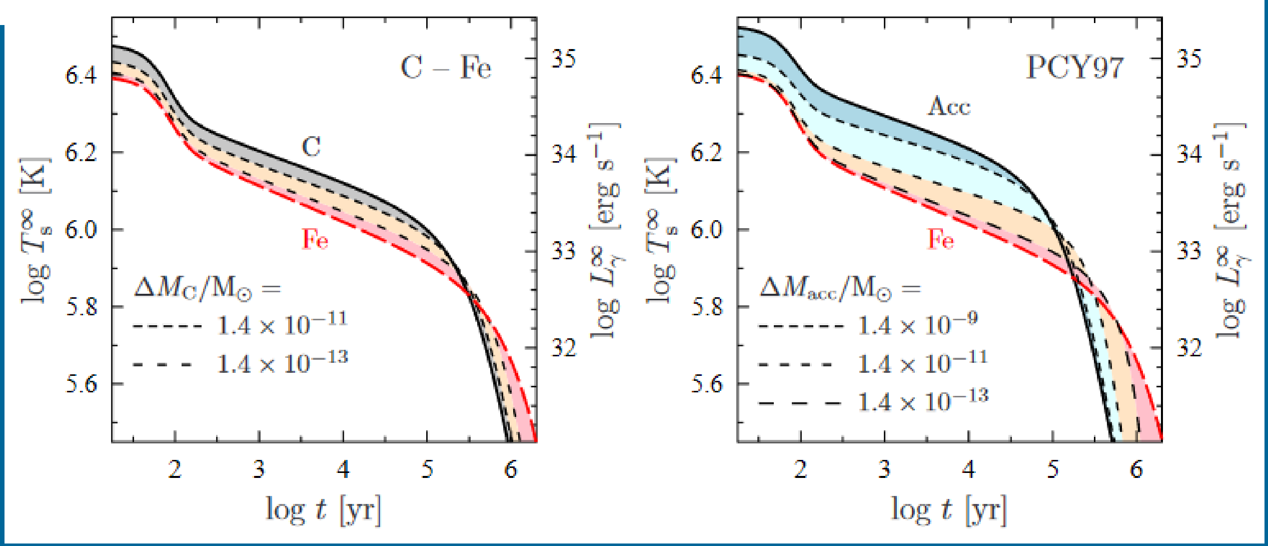


2103.12422

Composition dependence







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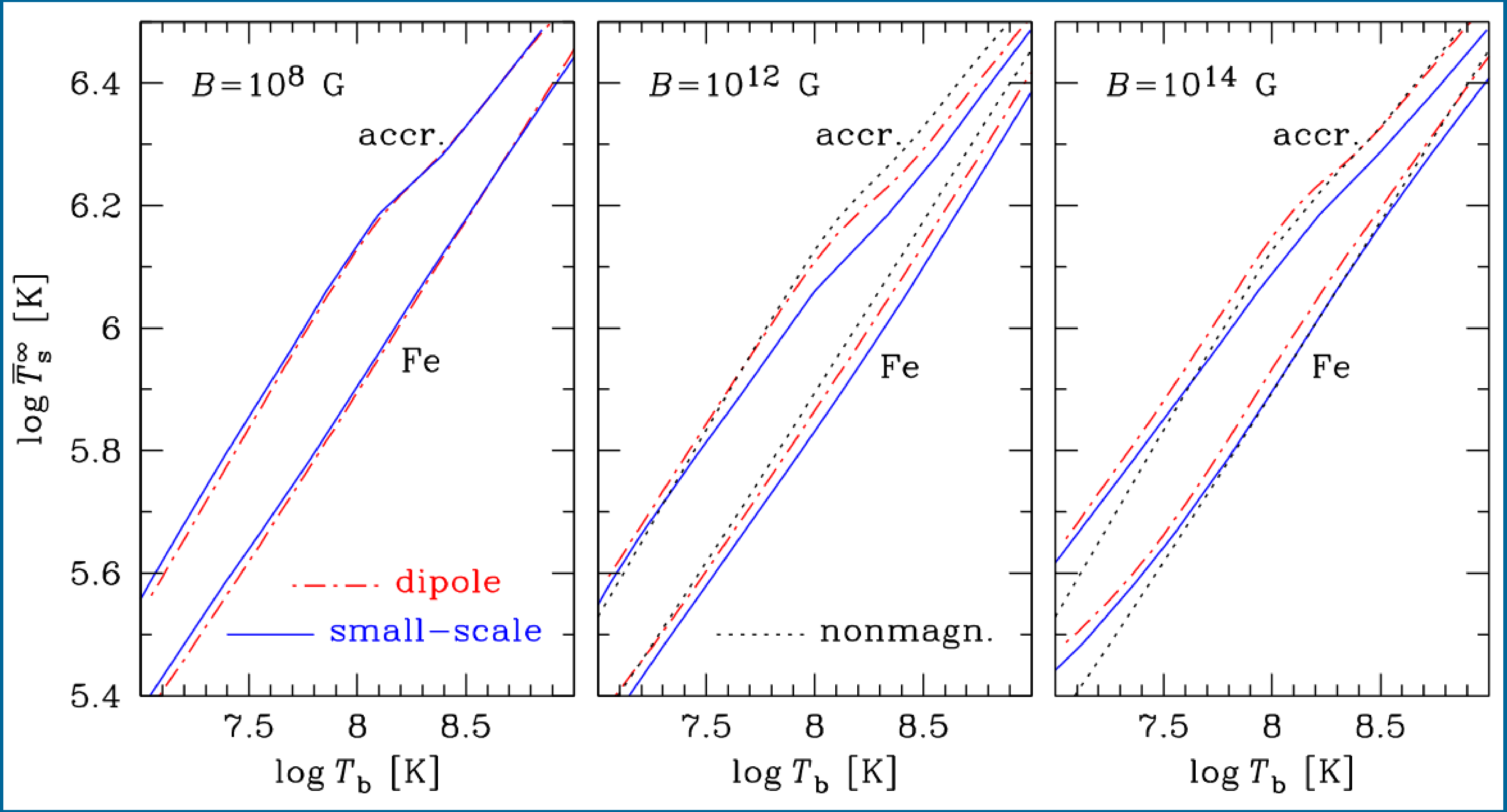
Composition and field dependences



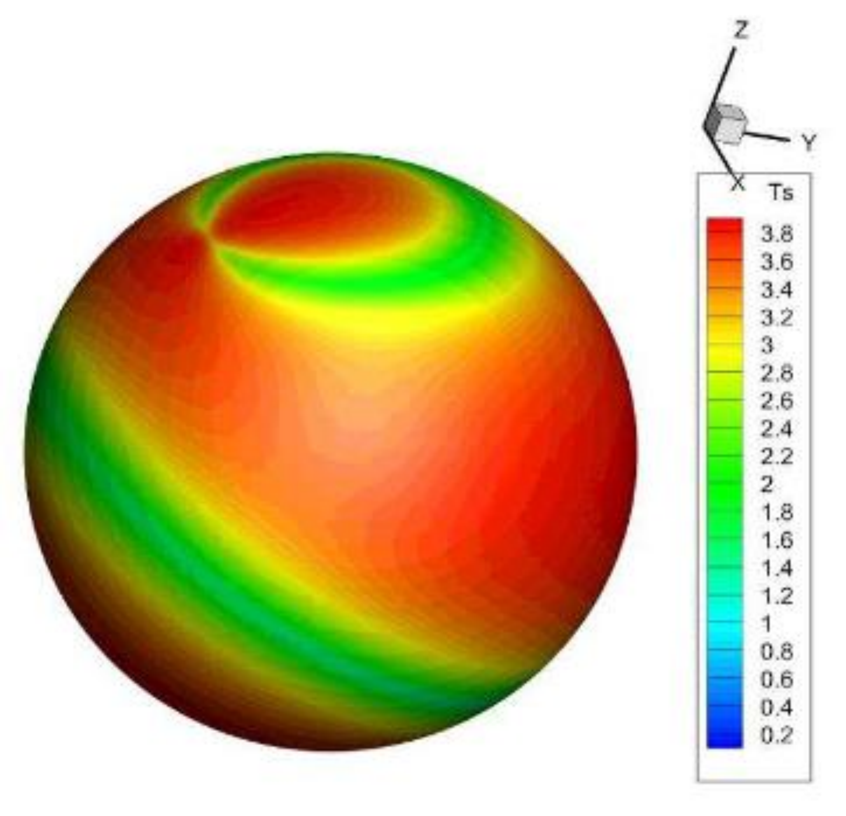
Envelope boundary: neutron drip density 4×10^{11} g cm^{-3}

astro-ph/0508415

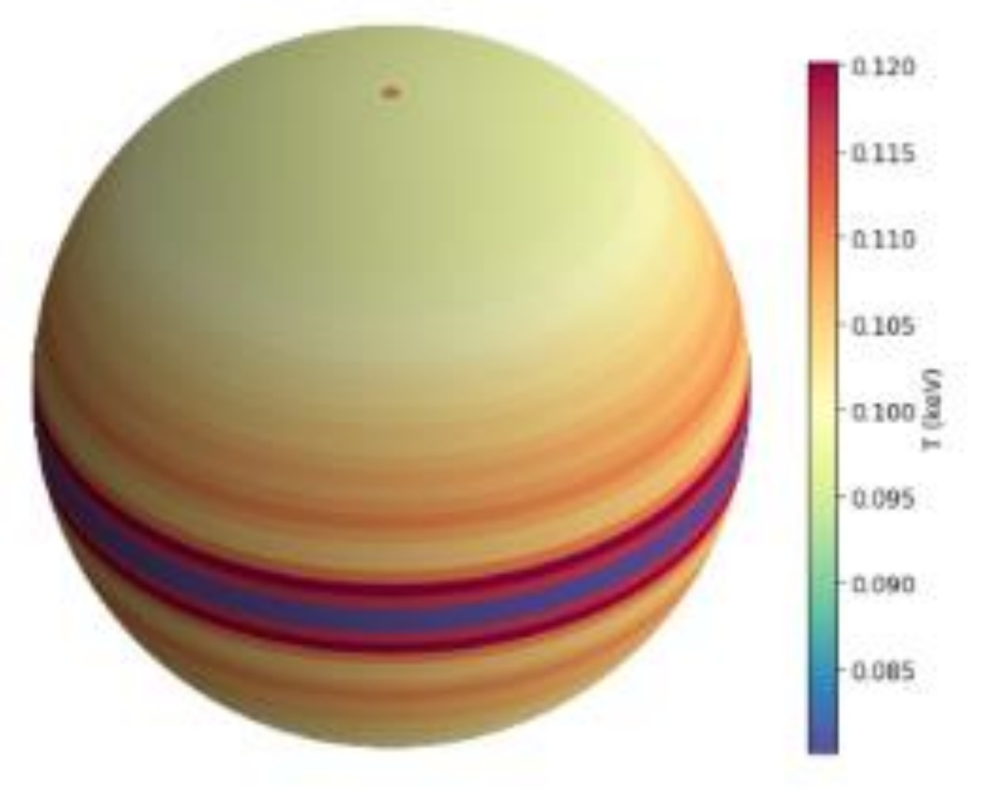
$L = 4\pi R^2 \sigma \overline{T}_s^4$ $L^\infty = (1 - r_g/R)\,L$ $T_s^\infty = T_s\,\sqrt{1 - r_g/R}$

Field and surface temperature

Complicated field topology in the crust can result
in complicated surface temperature distribution.



Dipole + quadrupole



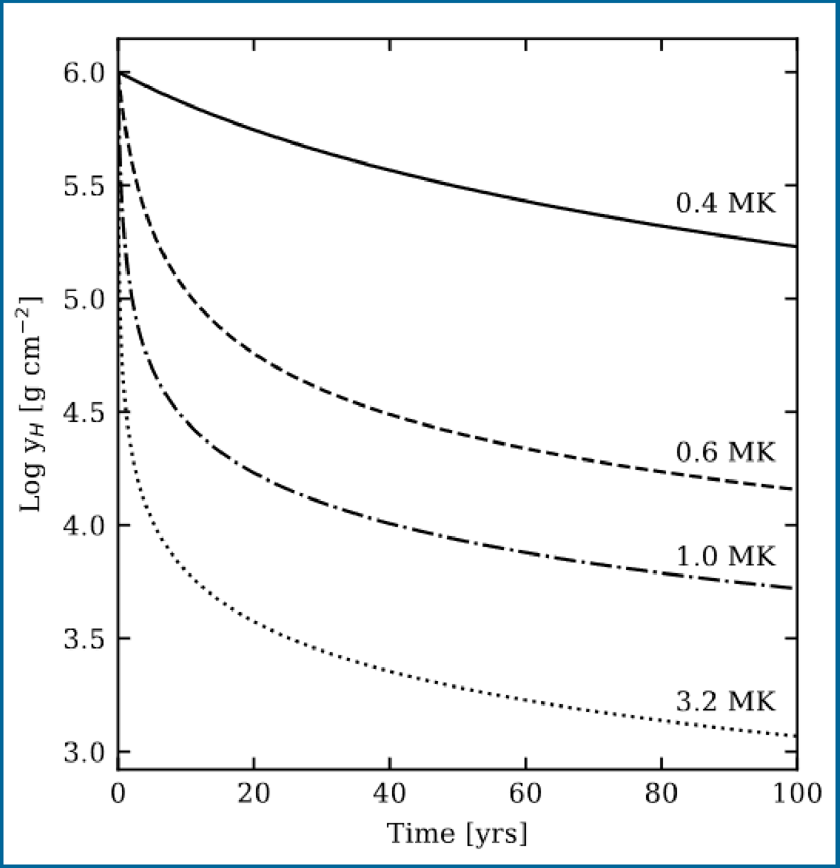
Poloidal + toroidal

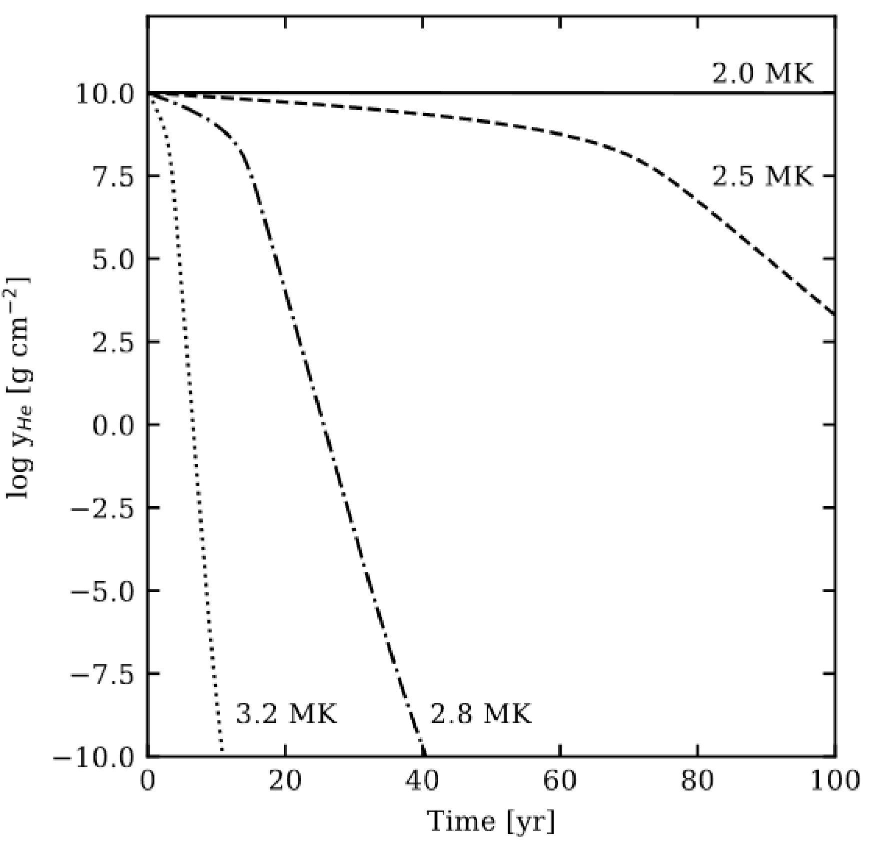
2007.11035

2009.04331

Diffusive nuclear burning

Time dependent envelope composition in early years.

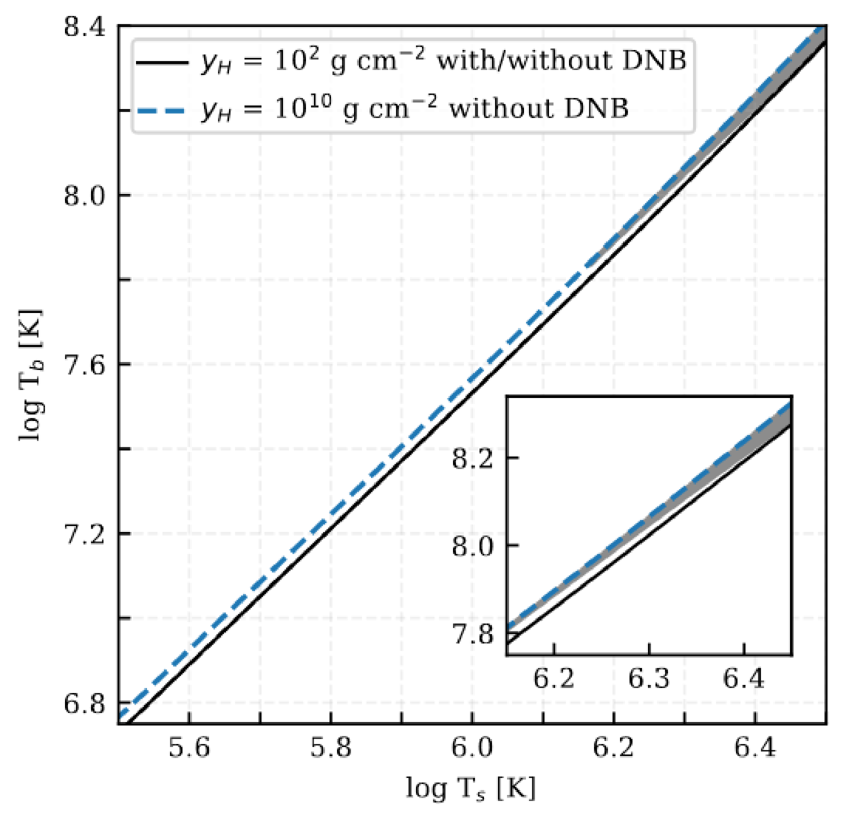


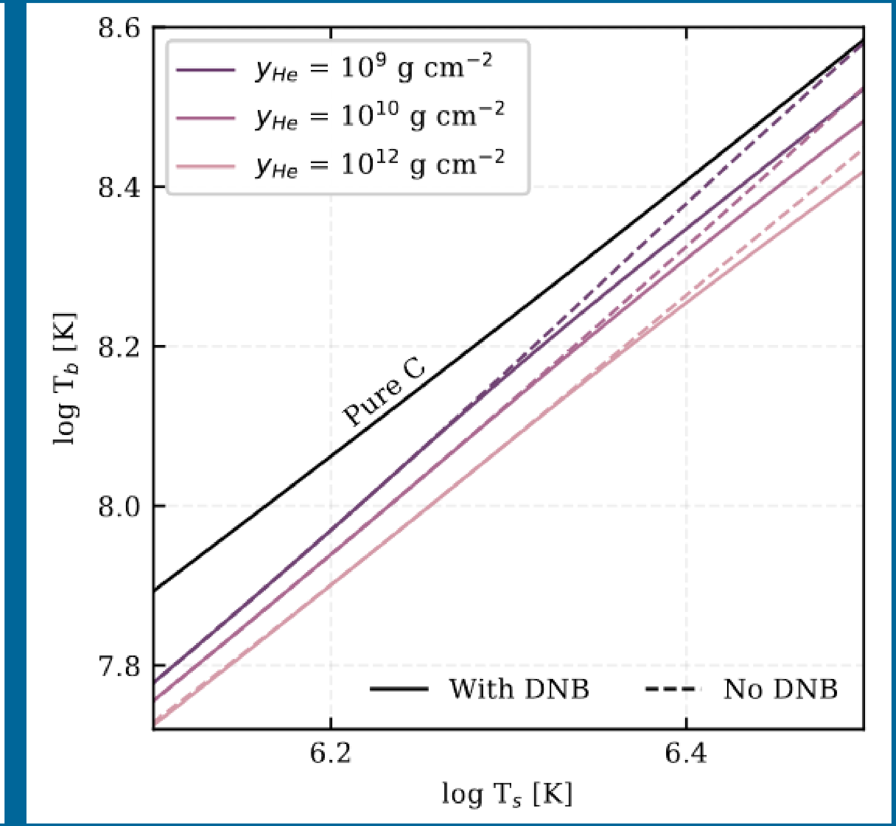


1901.01012

Cooling with DNB

Evolution of composition modifies cooling curves via Ts-Tb relation





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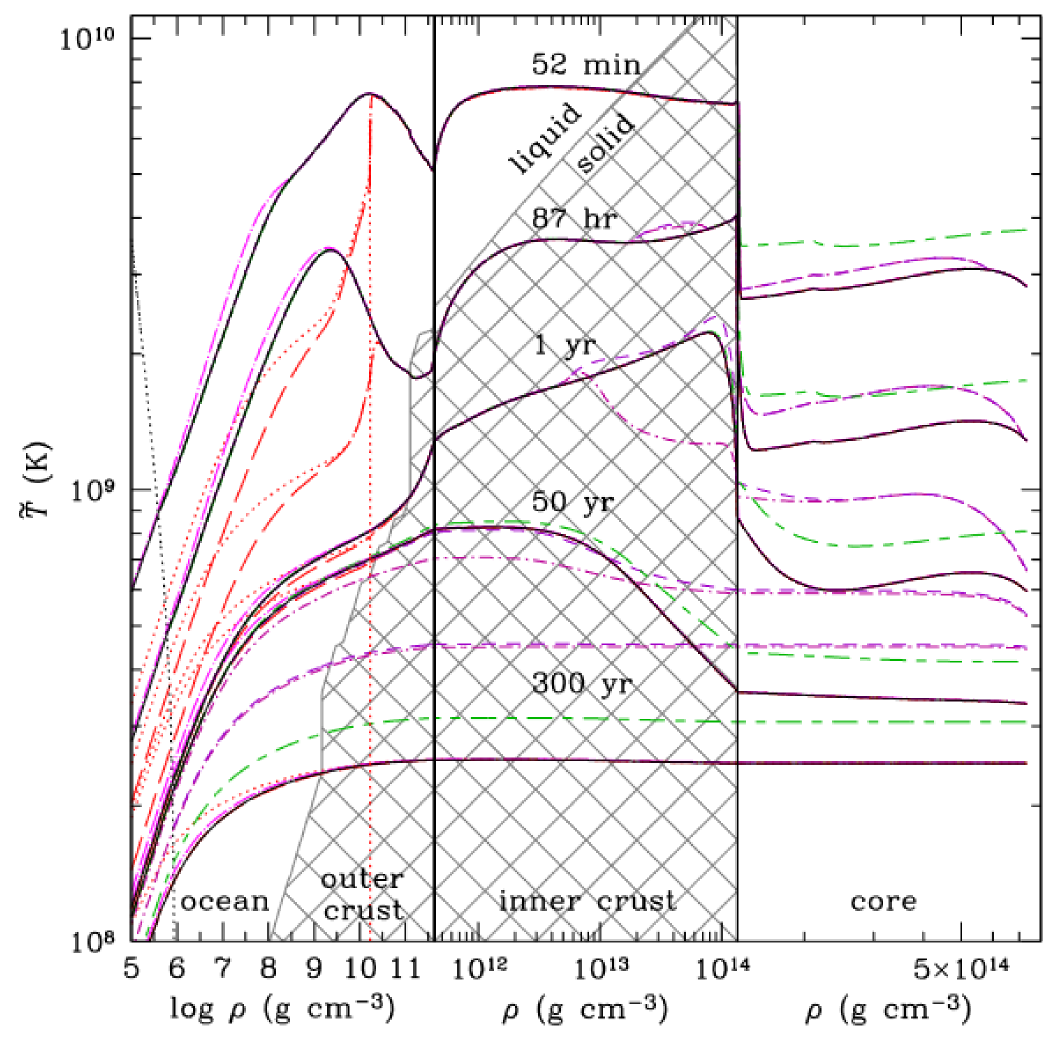
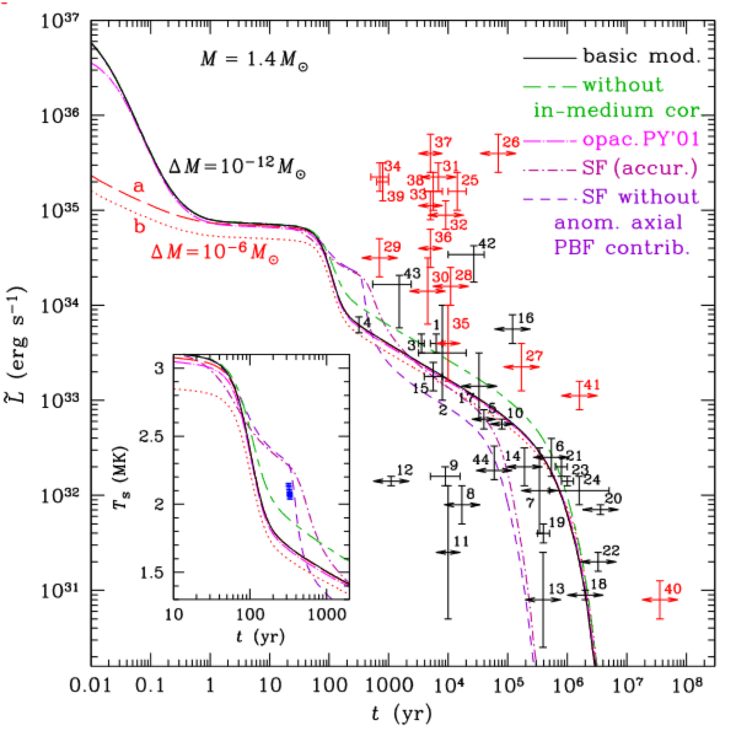
Redshifted temperature evolution





1711.07662

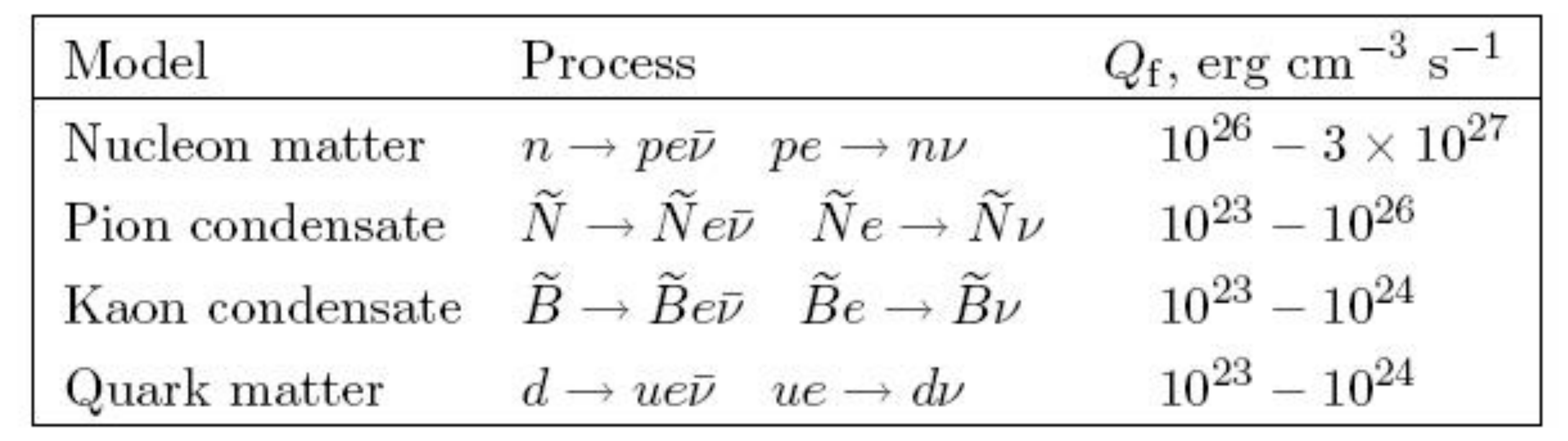
Cooling depends on:

1. Rate of neutrino emission from NS interiors
2. Heat capacity of internal parts of a star
3. Superfluidity
4. Thermal conductivity in the outer layers
5. Possible heating

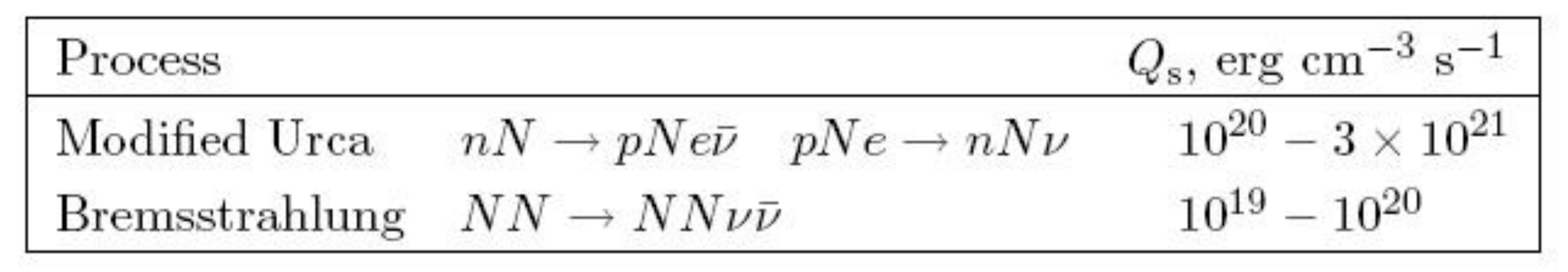
Depend on the EoS
and composition

(see Yakovlev & Pethick 2004)

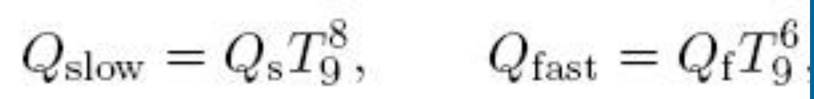
Main neutrino processes



Model	Process	Q_f, erg cm^{-3} s^{-1}
Nucleon matter	$n \to pe\bar{\nu} \quad pe \to n\nu$	$10^{26} - 3 \times 10^{27}$
Pion condensate	$\tilde{N} \to \tilde{N}e\bar{\nu} \quad \tilde{N}e \to \tilde{N}\nu$	$10^{23} - 10^{26}$
Kaon condensate	$\tilde{B} \to \tilde{B}e\bar{\nu} \quad \tilde{B}e \to \tilde{B}\nu$	$10^{23} - 10^{24}$
Quark matter	$d \to ue\bar{\nu} \quad ue \to d\nu$	$10^{23} - 10^{24}$



Process		Q_s, erg cm^{-3} s^{-1}
Modified Urca	$nN \to pNe\bar{\nu} \quad pNe \to nN\nu$	$10^{20} - 3 \times 10^{21}$
Bremsstrahlung	$NN \to NN\nu\bar{\nu}$	$10^{19} - 10^{20}$



$$Q_{\rm slow} = Q_{\rm s} T_9^8, \qquad Q_{\rm fast} = Q_{\rm f} T_9^6$$

(Yakovlev & Pethick astro-ph/0402143)

Fast Cooling (URCA cycle)

$n \longrightarrow p+e^-+\bar{\nu}_e$

$p+e^- \longrightarrow n+\nu_e$

Slow Cooling (modified URCA cycle)

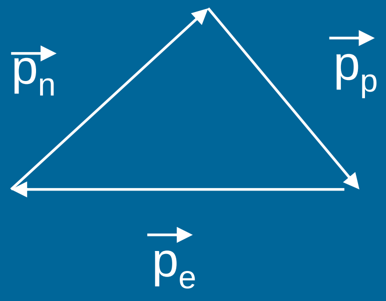
$n+n \longrightarrow n+p+e^-+\bar{\nu}_e$

$n+p+e^- \longrightarrow n+n+\nu_e$

$p+n \longrightarrow p+p+e^-+\bar{\nu}_e$

$p+p+e^- \longrightarrow p+n+\nu_e$

- Fast cooling possible only if $n_p > n_n/8$
- Nucleon Cooper pairing is important
- Minimal cooling scenario (Page et al 2004):
  - no exotica
  - no fast processes
  - pairing included



$p_n<p_p+p_e$

[See the book Haensel, Potekhin, Yakovlev p. 265 (p.286 in the file) and Shapiro, Teukolsky for details: Ch. 2.3, 2.5, 11.]

Equations

Neutrino emissivity

heating

$$\frac{\mathrm{e}^{-\lambda-2\Phi}}{4\pi r^2}\,\frac{\partial}{\partial r}\left(\mathrm{e}^{2\Phi}L_r\right) = -Q + Q_{\mathrm{h}} - \frac{c_T}{\mathrm{e}^{\Phi}}\,\frac{\partial T}{\partial t},$$

$$\frac{L_r}{4\pi\kappa r^2} = \mathrm{e}^{-\lambda-\Phi}\,\frac{\partial}{\partial r}\left(T\mathrm{e}^{\Phi}\right),$$

After thermal relaxation we have in the whole star: $T_i(t)=T(r,t)e^{\Phi(r)}$

$$\mathrm{e}^{-\lambda} = \sqrt{1-2Gm(r)/c^2 r},$$

At the surface we have: $\Phi(R) = -\lambda(R)$

$$C(T_i)\,\frac{\mathrm{d}T_i}{\mathrm{d}t} = -L_\nu^\infty(T_i) + L_{\mathrm{h}}^\infty - L_\gamma^\infty(T_{\mathrm{s}}),$$

$$L_\nu^\infty(T_i) = \int \mathrm{d}V\, Q(T)\,\mathrm{e}^{2\Phi}, \quad \text{and } L_{\mathrm{h}}^\infty = \int \mathrm{d}V\, Q_{\mathrm{h}}\,\mathrm{e}^{2\Phi}, \quad C(T_i) = \int \mathrm{d}V\, c_T(T),$$

$\mathrm{d}V = 4\pi r^2 \mathrm{e}^{\lambda}\,\mathrm{d}r$ is the element of proper volume

L_ν^∞ is the total neutrino luminosity (for a distant observer)

L_{h}^∞ is the total reheating power.

(Yakovlev & Pethick 2004)

Total stellar heat capacity

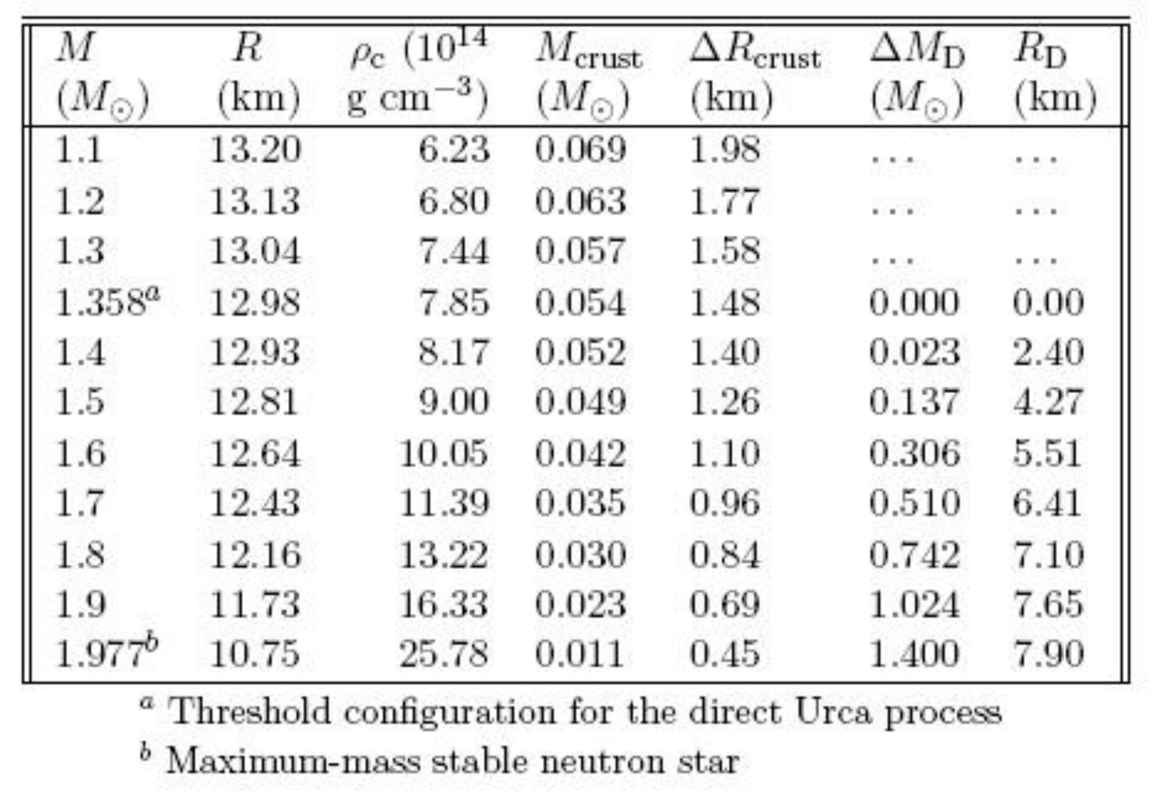
Simplified model of a cooling NS

No superfluidity, no envelopes and magnetic fields, only hadrons.

The most critical moment is the onset of direct URCA cooling.

ρ_D= 7.851 10^{14} g/cm^3.

The critical mass depends on the EoS. For the examples below M_D=1.358 M_{solar}.

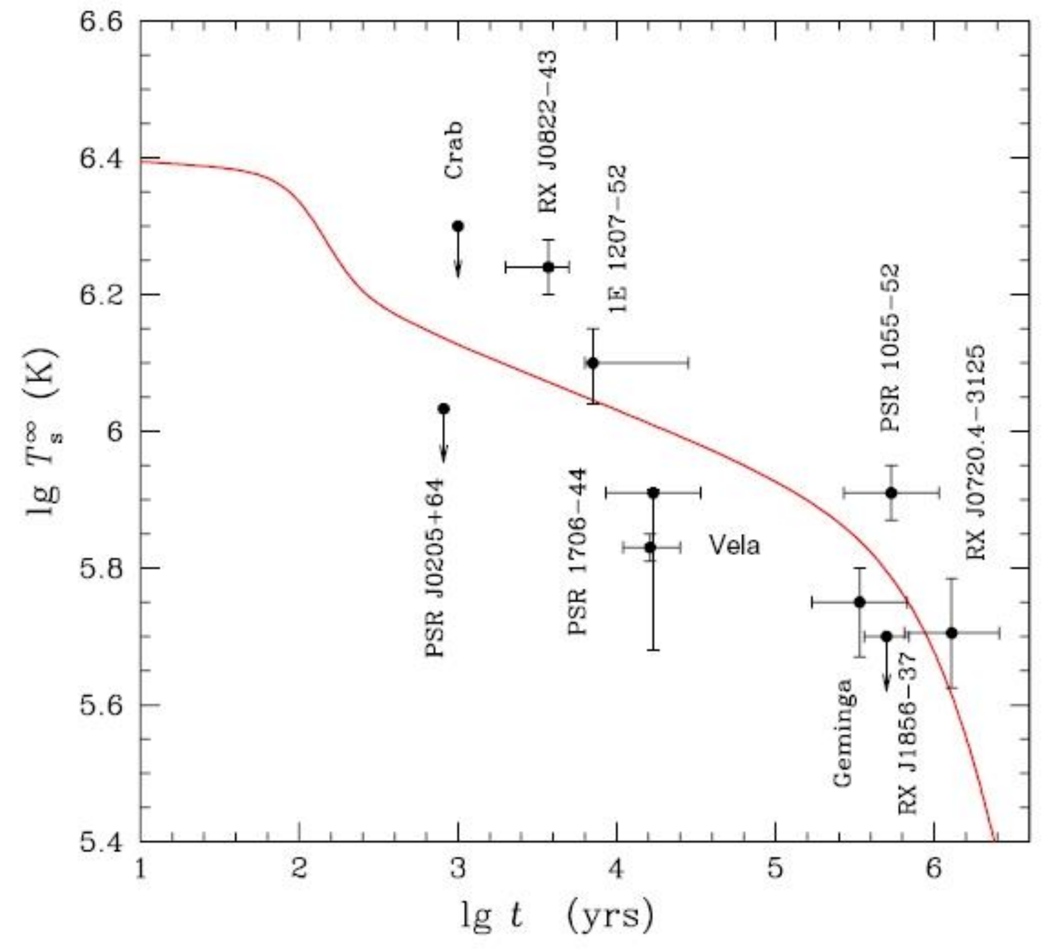


M ($M_\odot$)	R (km)	ρ_c (10^{14} g cm^{-3})	$M_{\rm crust}$ ($M_\odot$)	$\Delta R_{\rm crust}$ (km)	$\Delta M_{\rm D}$ ($M_\odot$)	$R_{\rm D}$ (km)
1.1	13.20	6.23	0.069	1.98	…	…
1.2	13.13	6.80	0.063	1.77	…	…
1.3	13.04	7.44	0.057	1.58	…	…
1.358[a]	12.98	7.85	0.054	1.48	0.000	0.00
1.4	12.93	8.17	0.052	1.40	0.023	2.40
1.5	12.81	9.00	0.049	1.26	0.137	4.27
1.6	12.64	10.05	0.042	1.10	0.306	5.51
1.7	12.43	11.39	0.035	0.96	0.510	6.41
1.8	12.16	13.22	0.030	0.84	0.742	7.10
1.9	11.73	16.33	0.023	0.69	1.024	7.65
1.977[b]	10.75	25.78	0.011	0.45	1.400	7.90

[a] Threshold configuration for the direct Urca process

[b] Maximum-mass stable neutron star

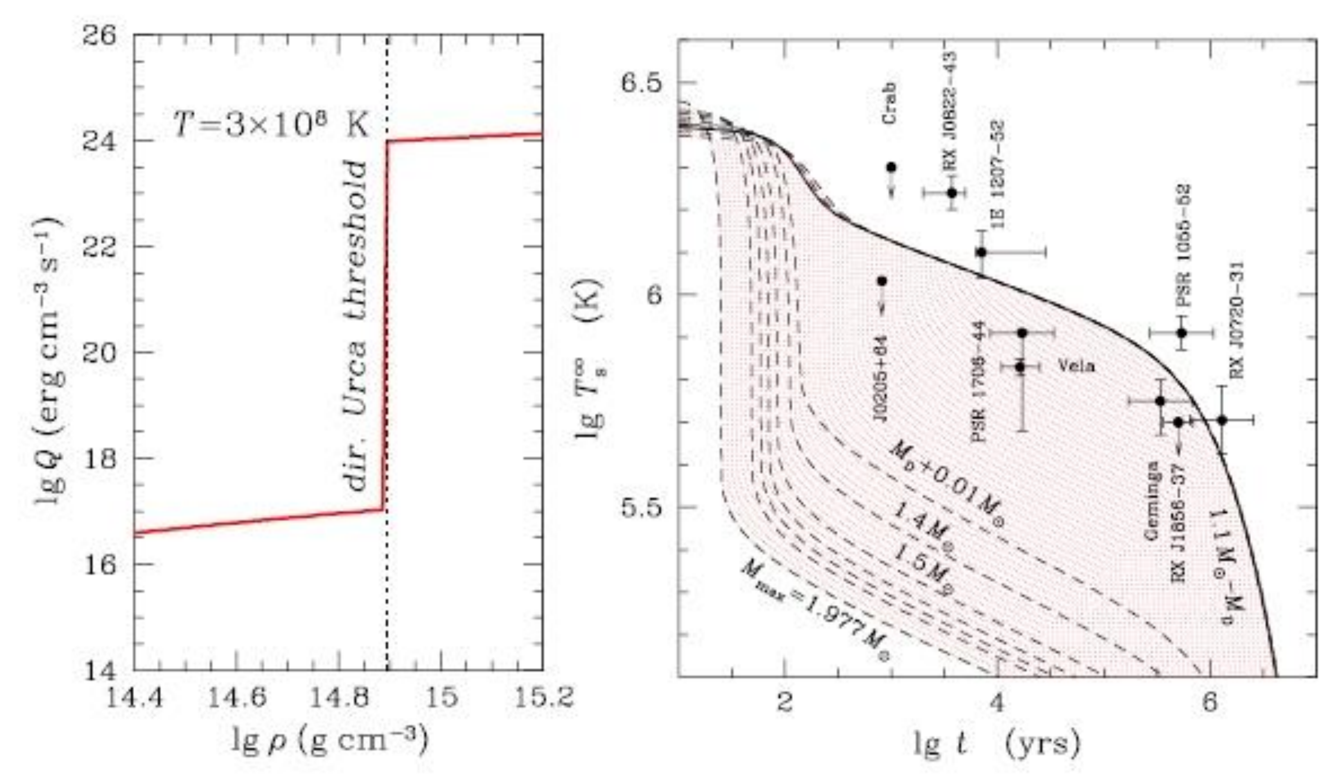
Simple cooling model for low-mass NSs.



Too hot ......
Too cold ....

(Yakovlev & Pethick 2004)

Nonsuperfluid nucleon cores

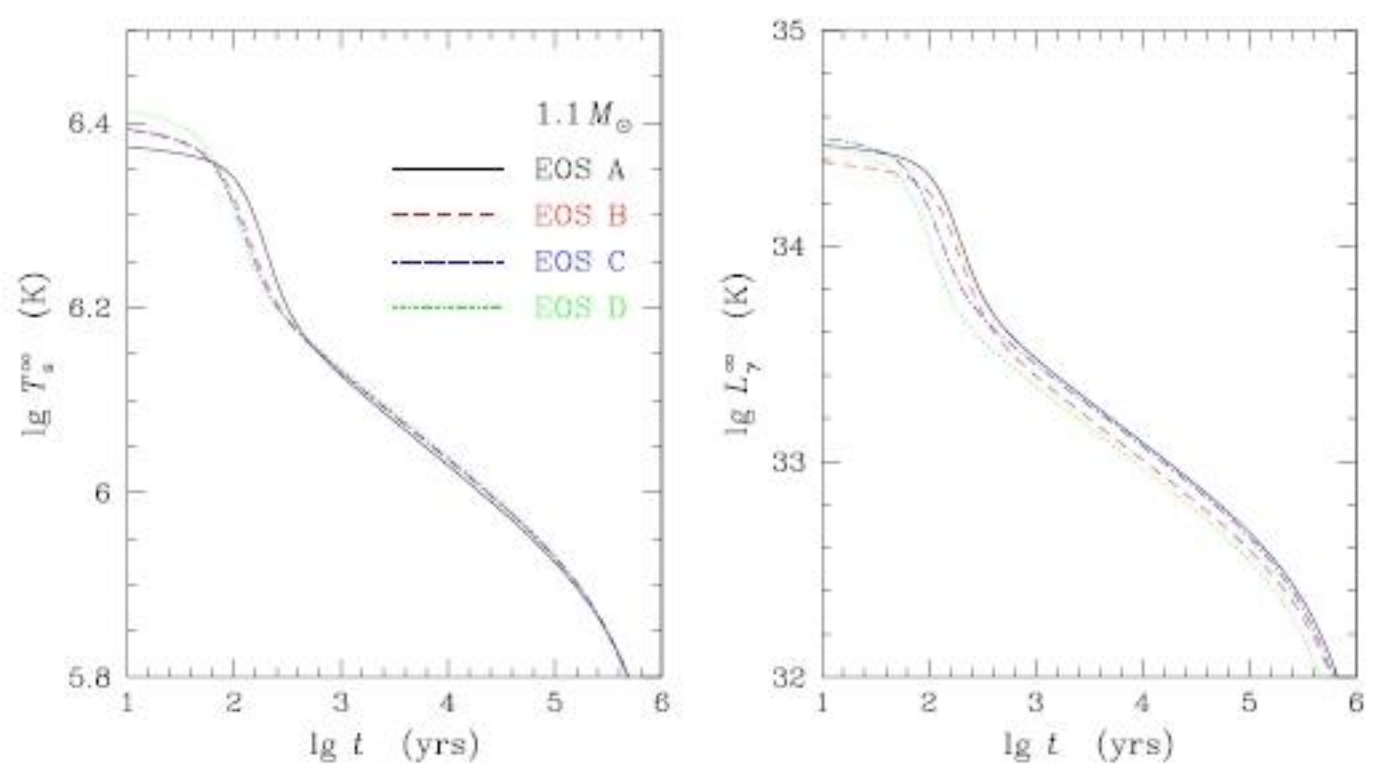


Note “population aspects” of the right plot: too many NSs have to be explained by a very narrow range of mass.

For slow cooling at the neutrino cooling stage t_{slow}~1 yr/T_{i9}^{6}
For fast cooling t_{fast}~ 1 min/T_{i9}^{4}

(Yakovlev & Pethick 2004)

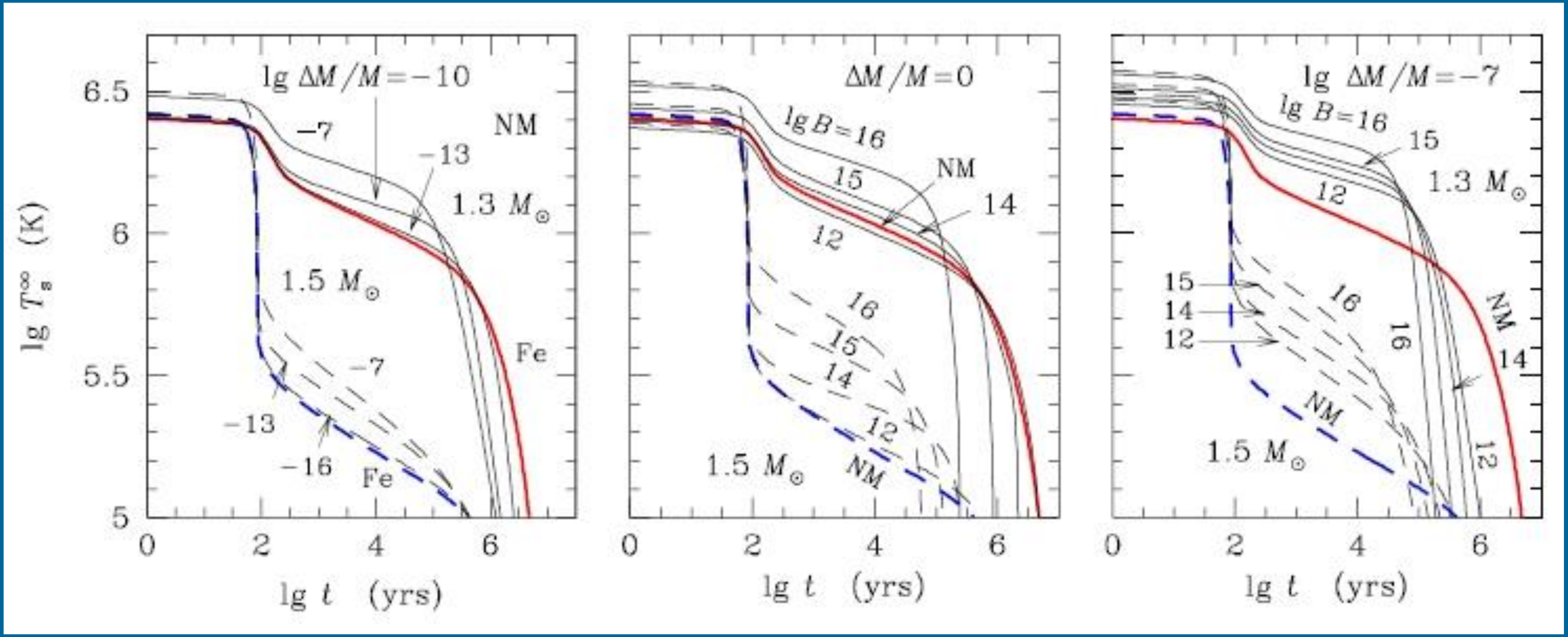
Slow cooling for different EoS



For slow cooling there is nearly no dependence on the EoS.
The same is true for cooling curves for maximum mass for each EoS.

(Yakovlev & Pethick 2004)

Envelopes and magnetic field



Non-magnetic stars

Thick lines – no envelope

No accreted envelopes, different magnetic fields.

Thick lines – non-magnetic

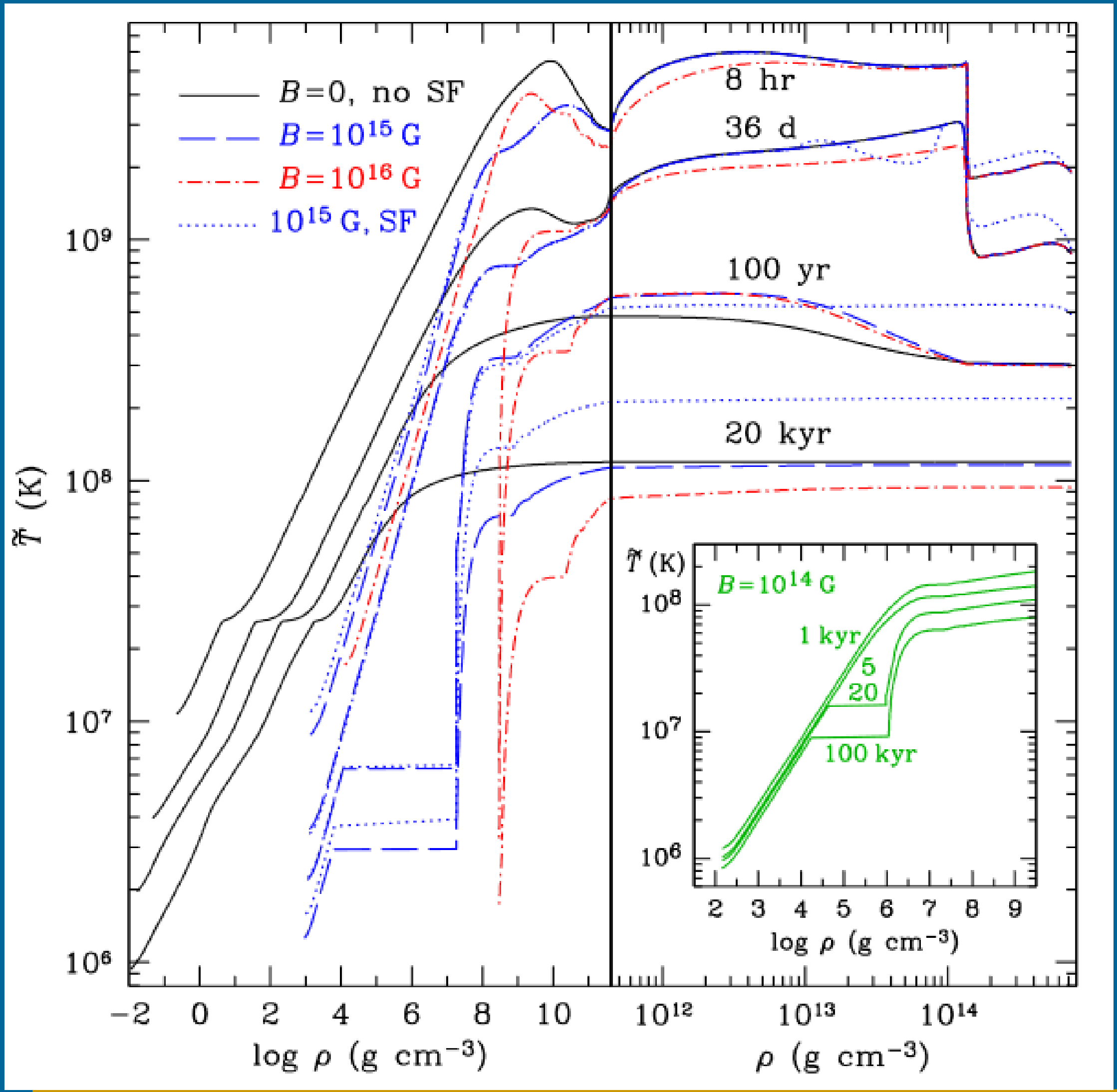
Envelopes + Fields

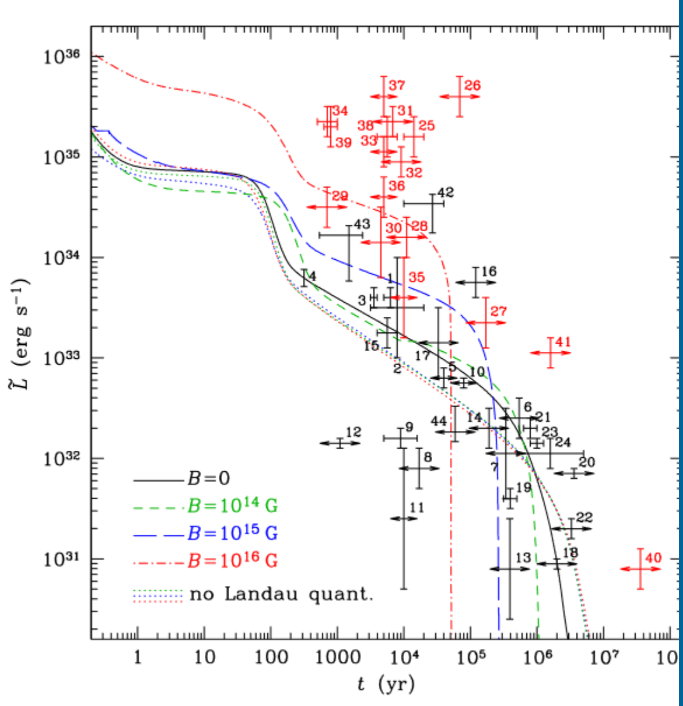
Solid line M=1.3 M_{solar}, Dashed lines M=1.5 M_{solar}

Envelopes can be related to the fact that we see a subpopulation of hot NS in CCOs with relatively long initial spin periods and low magnetic field, but do not observed representatives of this population around us, i.e. in the Solar vicinity.

(Yakovlev & Pethick 2004)

Magnetar cooling





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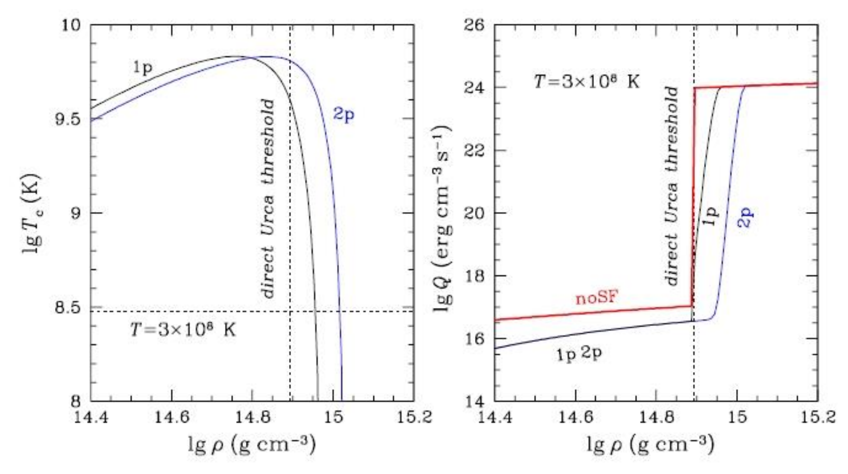
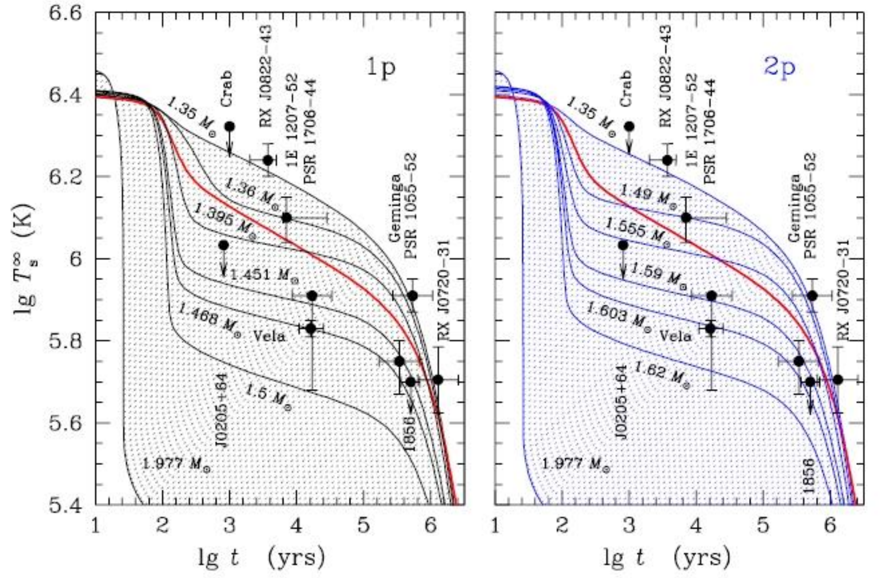
Simplified model: no neutron superfluidity



Superfluidity is an important ingredient of cooling models.
It is important to consider different types of proton and neutron superfluidity.

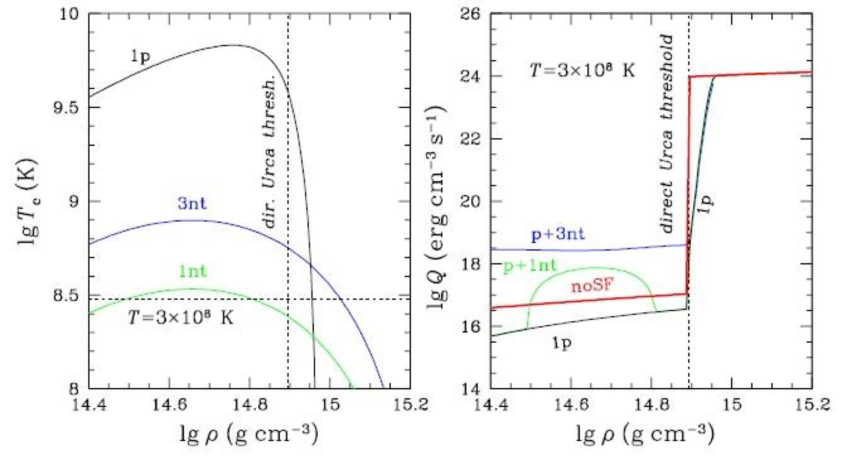
There is no complete microphysical theory which can describe superfluidity in neutron stars.

If proton superfluidity is strong,
but neutron superfluidity
in the core is weak
then it is possible
to explain observations.



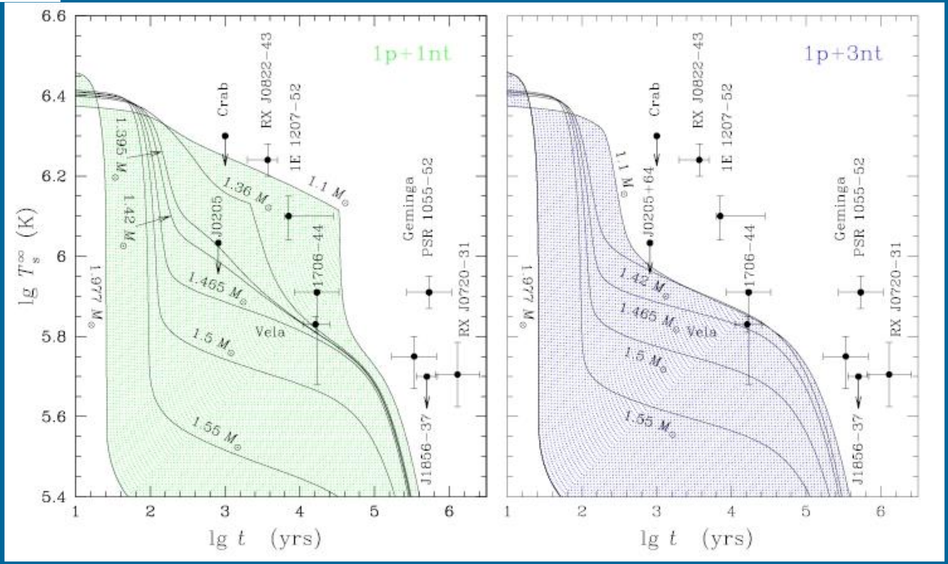
(Yakovlev & Pethick 2004)

Neutron superfluidity and observations



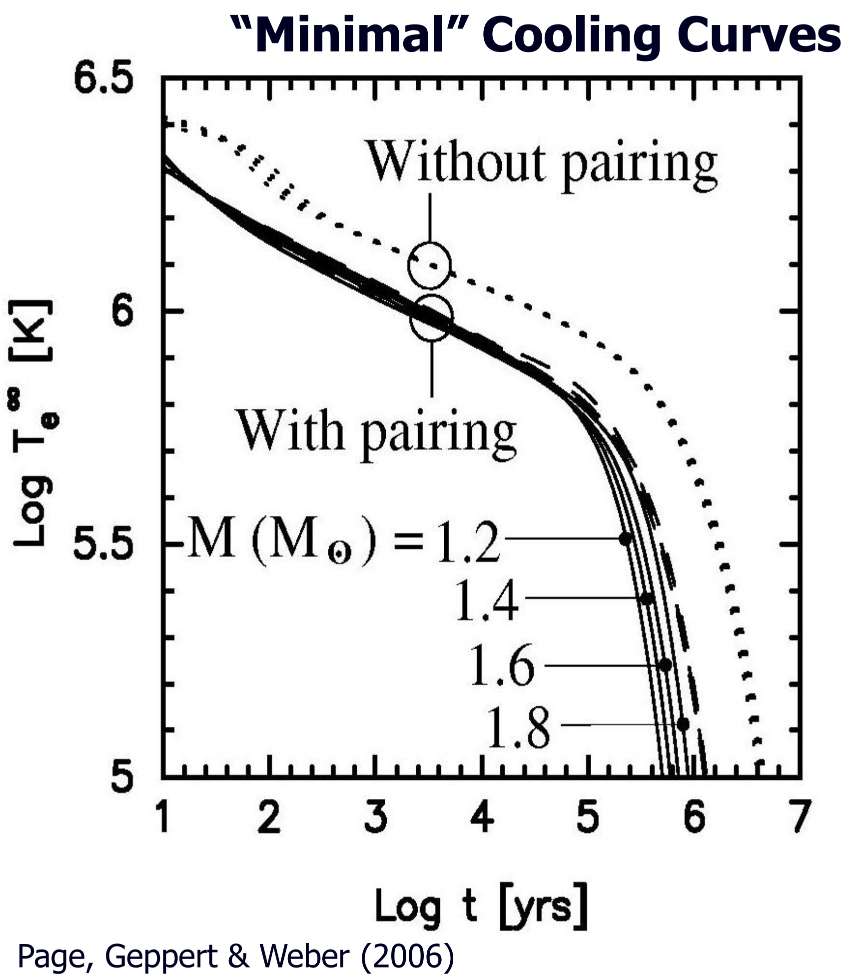
Mild neutron pairing in the core contradicts observations.

See a recent review about superfluidity and its relation to the thermal evolution of NSs in 1206.5011 and a very detailed review about superfluids in NSs in 1302.6626. A brief and more popular review in 1303.3282.



(Yakovlev & Pethick 2004)

Minimal cooling model



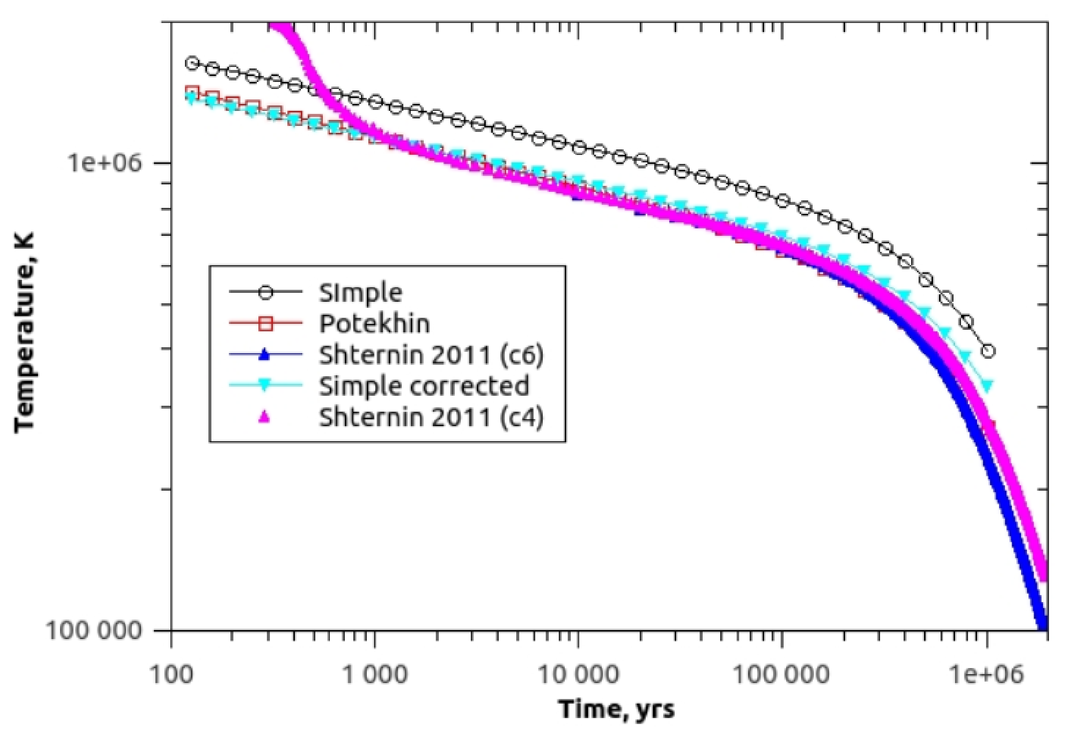
Page, Geppert & Weber (2006)

"minimal" means
without additional cooling
due to direct URCA
and without additional heating

Main ingredients of the minimal model

- EoS
- Superfluid properties
- Envelope composition
- NS mass

Analytical fits



$T_{\mathrm{surface}} \sim T_{\mathrm{core}}^{1/2}$

$$T_{\mathrm{eff6}} = T_* \equiv (7\,T_{\mathrm{b9}}\sqrt{g_{14}})^{1/2}\,.$$

(iron envelope)

$$T^4_{\mathrm{eff6,a}} = g_{14}\,(18.1\,T_{\mathrm{b9}})^{2.42},$$

(accreted envelope)

astro-ph/9706148

astro-ph/0105261

$$T = b\left(\frac{t}{1\mathrm{yr}}\right)^a \exp(-t/\tau_{\mathrm{c}}).$$

1709.10385

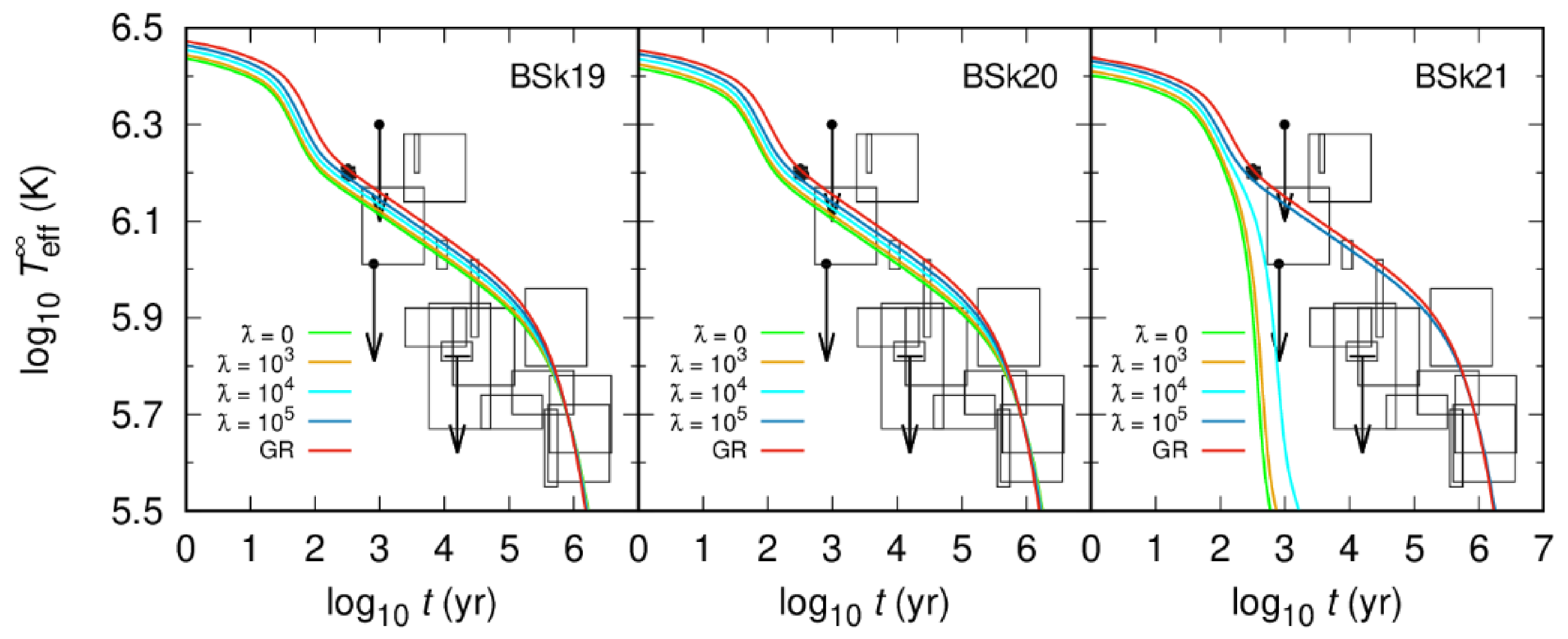
$$T_{\mathrm{s}}^{(0)} \approx 10^6\, g_{14}^{1/4}\left[(7\zeta)^{2.25} + (\zeta/3)^{1.25}\right]^{1/4}\ \mathrm{K}, \qquad (27)$$

where $\zeta \equiv T_{\mathrm{int,9}} - 0.001\, g_{14}^{1/4}\sqrt{7\,T_{\mathrm{int,9}}}$, $T_{\mathrm{int,9}} \equiv T_{\mathrm{int}}/(10^9\ \mathrm{K})$ and $g_{14} \equiv g/10^{14}\ \mathrm{cm\ s^{-2}}$.

NS cooling in modified gravity theories

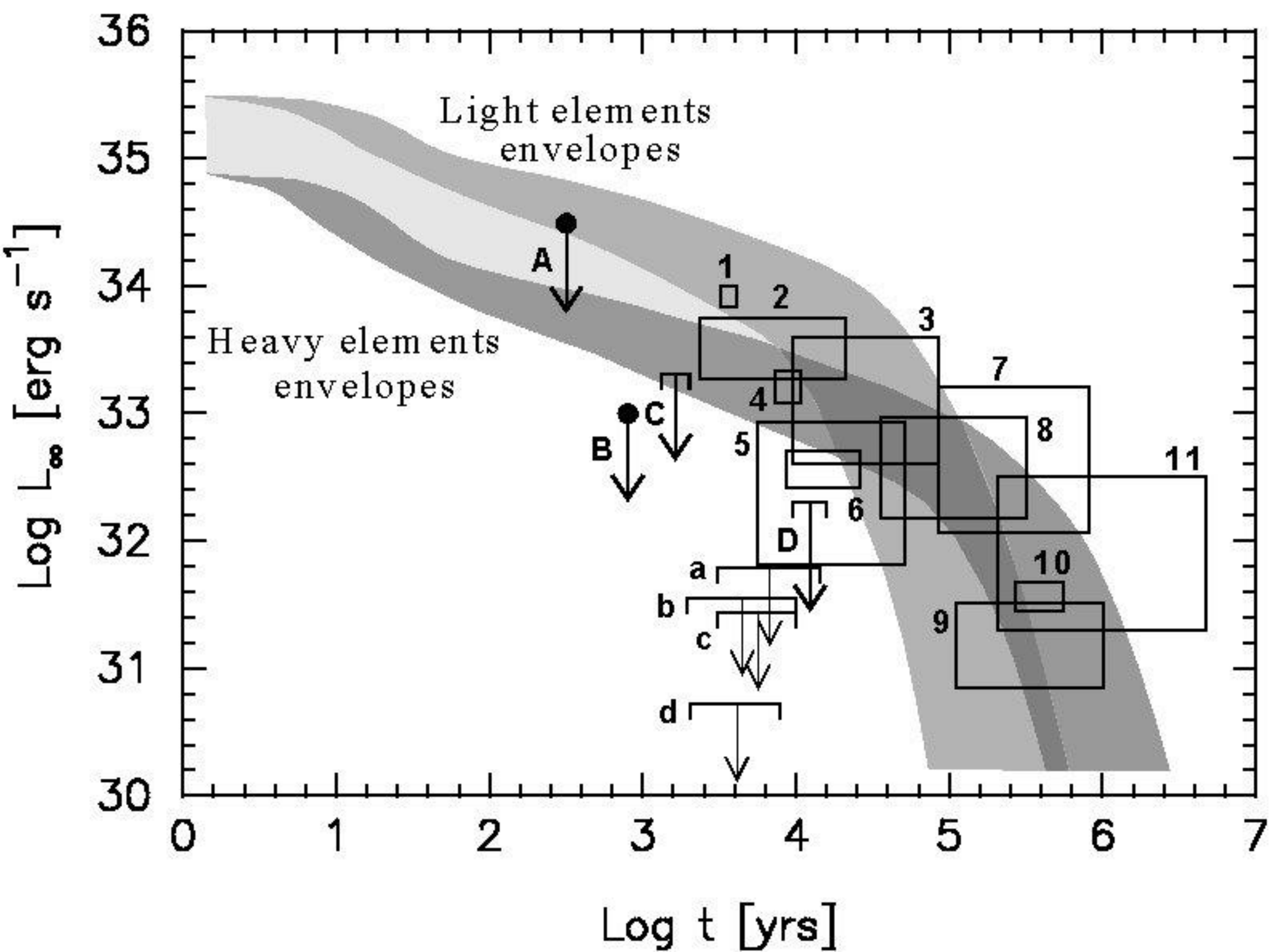
scalar-tensor theories

M=1.4 Msolar



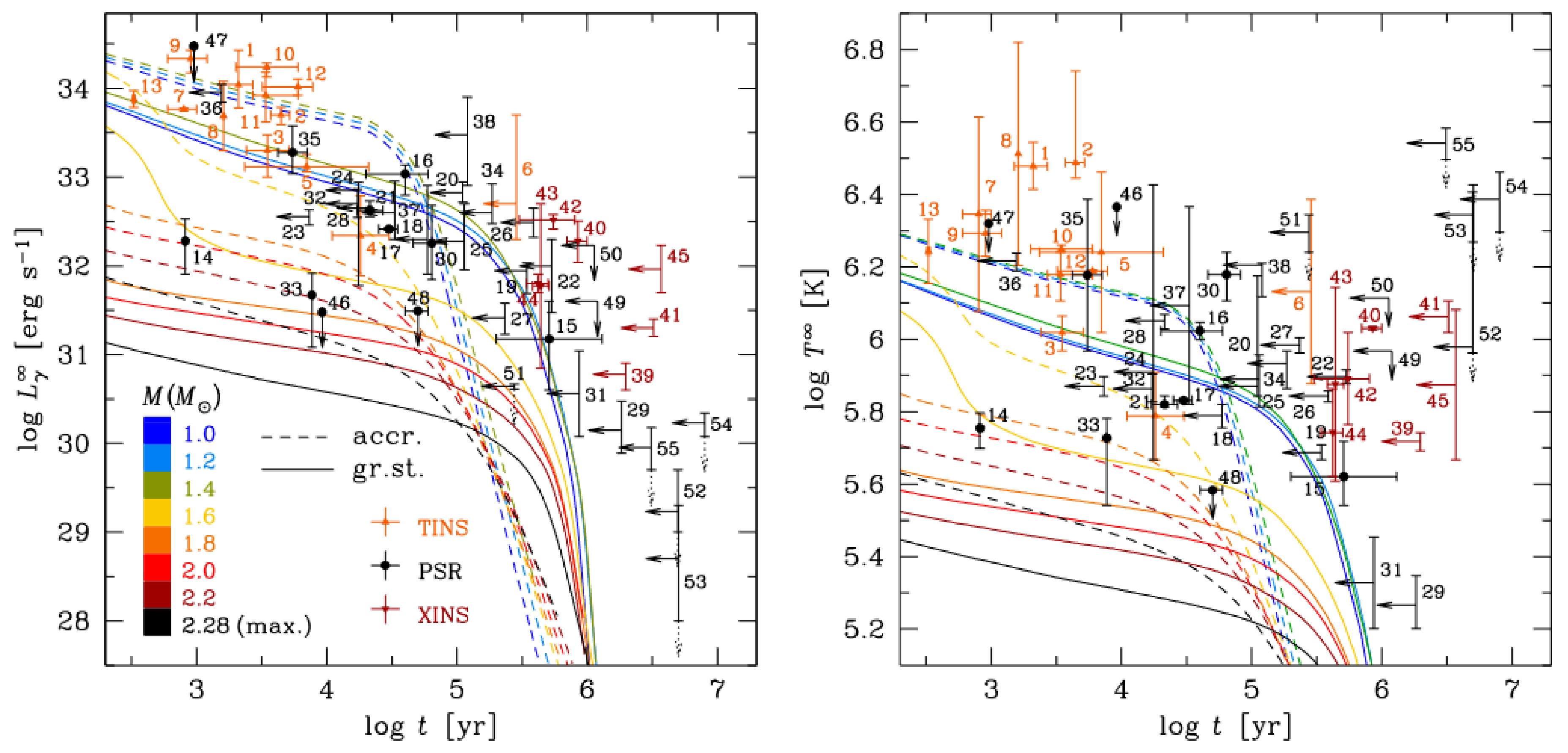
2003.12571

Luminosity and age uncertainties



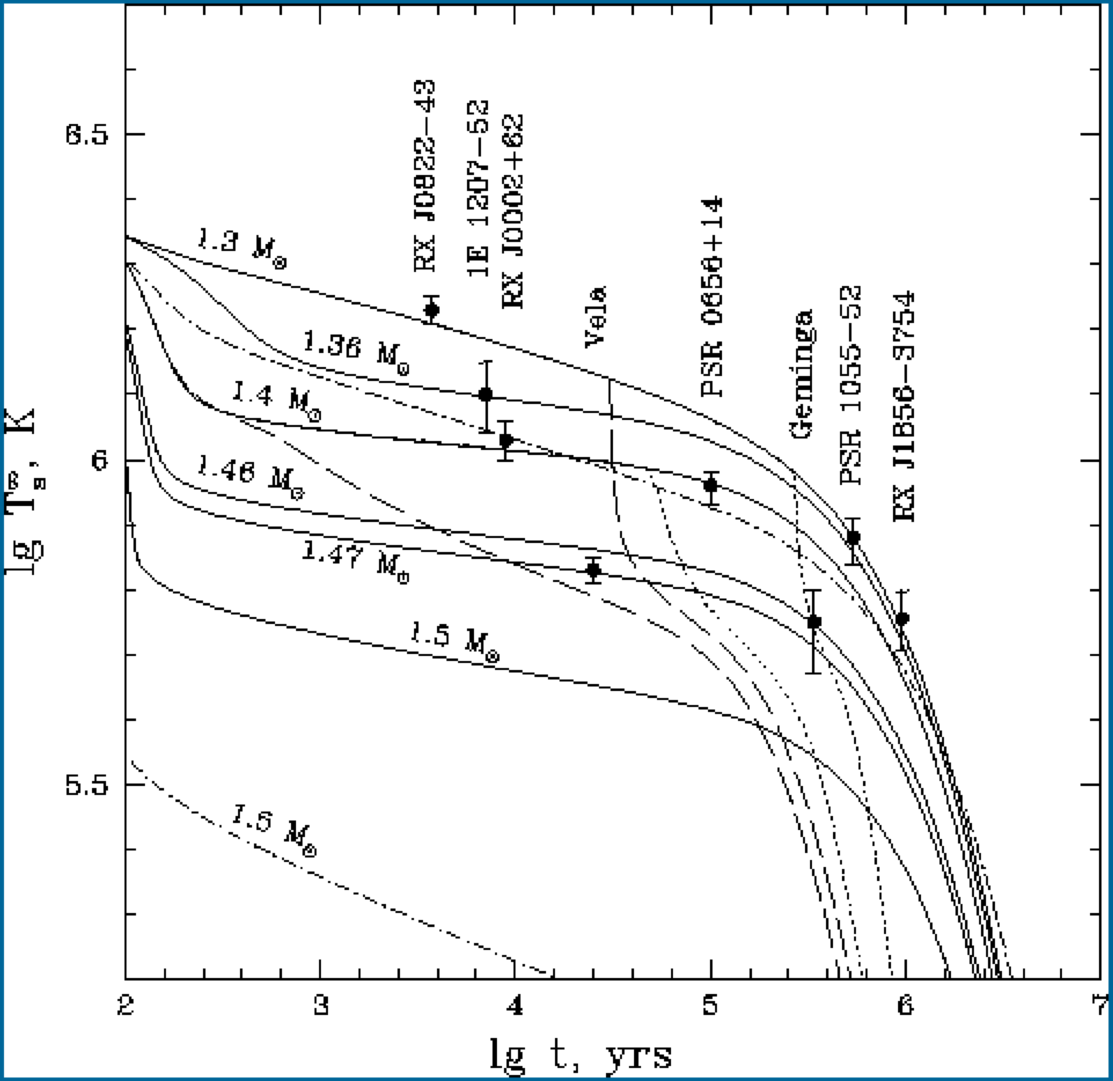
Page, Geppert, Weber
astro-ph/0508056

Large set of data



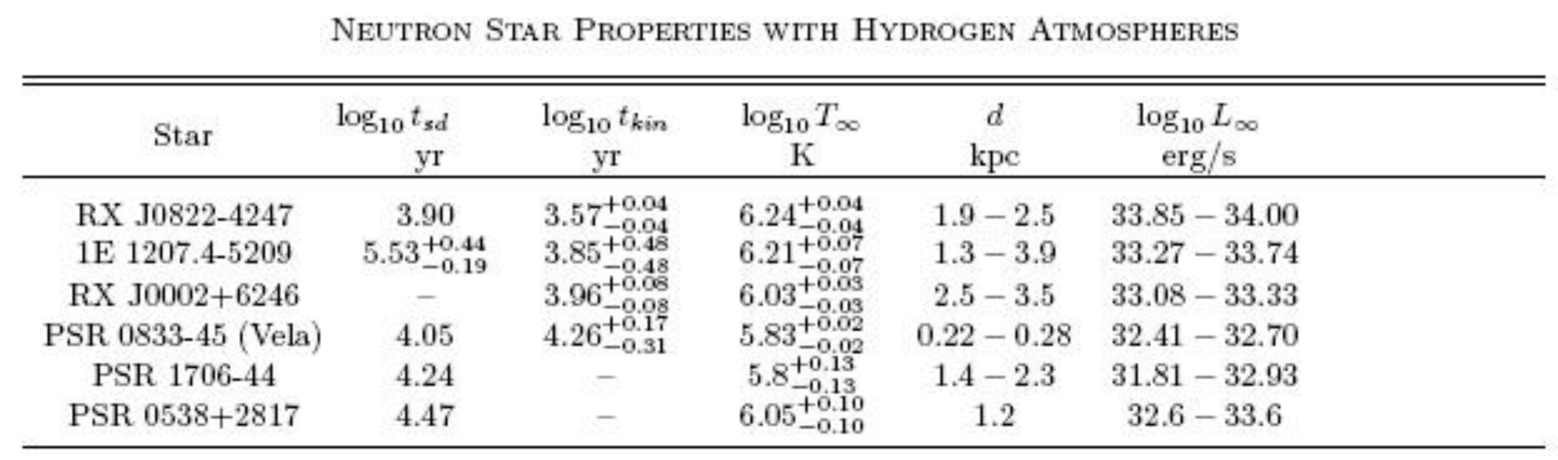
2006.15004

Standard test: temperature vs. age



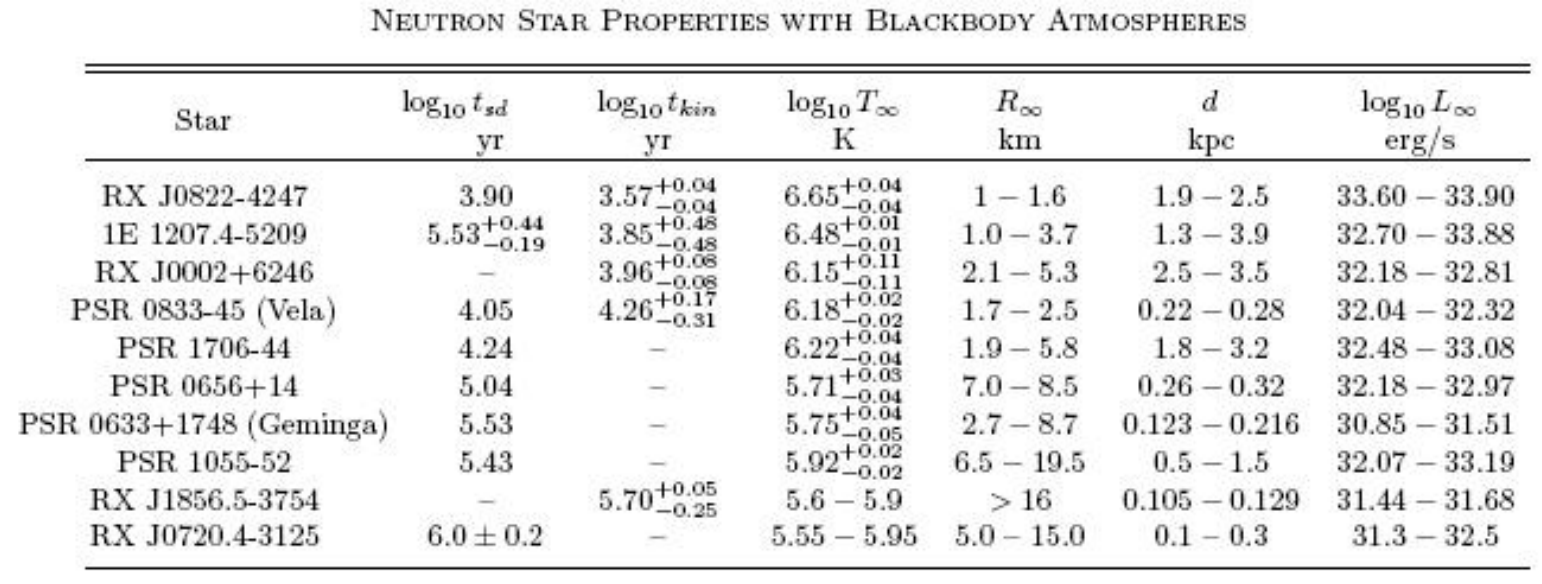
Kaminker et al. (2001)

Data



NEUTRON STAR PROPERTIES WITH HYDROGEN ATMOSPHERES

Star	$\log_{10} t_{sd}$ yr	$\log_{10} t_{kin}$ yr	$\log_{10} T_\infty$ K	d kpc	$\log_{10} L_\infty$ erg/s
RX J0822-4247	3.90	$3.57^{+0.04}_{-0.04}$	$6.24^{+0.04}_{-0.04}$	$1.9 - 2.5$	$33.85 - 34.00$
1E 1207.4-5209	$5.53^{+0.44}_{-0.19}$	$3.85^{+0.48}_{-0.48}$	$6.21^{+0.07}_{-0.07}$	$1.3 - 3.9$	$33.27 - 33.74$
RX J0002+6246	–	$3.96^{+0.08}_{-0.08}$	$6.03^{+0.03}_{-0.03}$	$2.5 - 3.5$	$33.08 - 33.33$
PSR 0833-45 (Vela)	4.05	$4.26^{+0.17}_{-0.31}$	$5.83^{+0.02}_{-0.02}$	$0.22 - 0.28$	$32.41 - 32.70$
PSR 1706-44	4.24	–	$5.8^{+0.13}_{-0.13}$	$1.4 - 2.3$	$31.81 - 32.93$
PSR 0538+2817	4.47	–	$6.05^{+0.10}_{-0.10}$	1.2	$32.6 - 33.6$

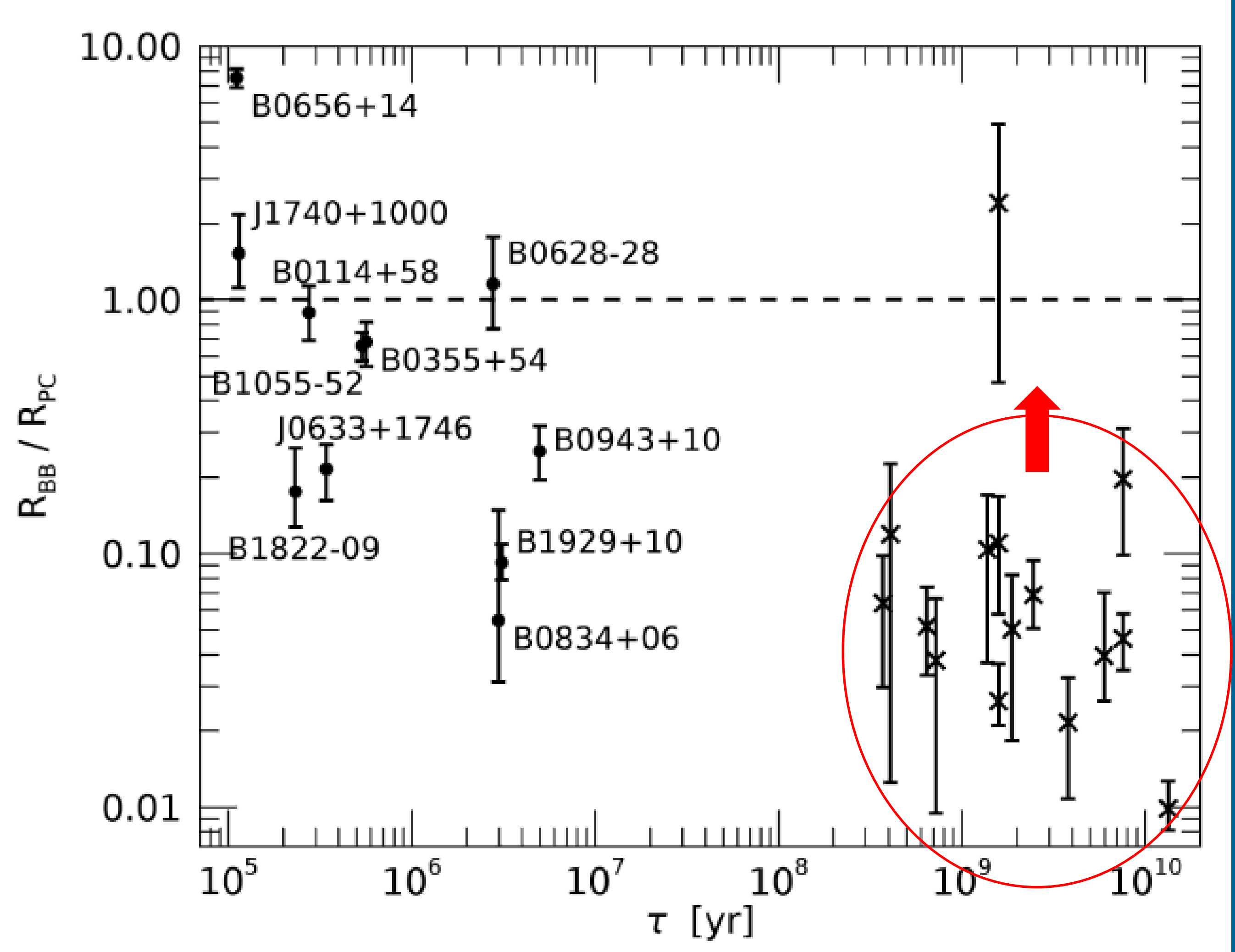


NEUTRON STAR PROPERTIES WITH BLACKBODY ATMOSPHERES

Star	$\log_{10} t_{sd}$ yr	$\log_{10} t_{kin}$ yr	$\log_{10} T_\infty$ K	R_∞ km	d kpc	$\log_{10} L_\infty$ erg/s
RX J0822-4247	3.90	$3.57^{+0.04}_{-0.04}$	$6.65^{+0.04}_{-0.04}$	$1 - 1.6$	$1.9 - 2.5$	$33.60 - 33.90$
1E 1207.4-5209	$5.53^{+0.44}_{-0.19}$	$3.85^{+0.48}_{-0.48}$	$6.48^{+0.01}_{-0.01}$	$1.0 - 3.7$	$1.3 - 3.9$	$32.70 - 33.88$
RX J0002+6246	–	$3.96^{+0.08}_{-0.08}$	$6.15^{+0.11}_{-0.11}$	$2.1 - 5.3$	$2.5 - 3.5$	$32.18 - 32.81$
PSR 0833-45 (Vela)	4.05	$4.26^{+0.17}_{-0.31}$	$6.18^{+0.02}_{-0.02}$	$1.7 - 2.5$	$0.22 - 0.28$	$32.04 - 32.32$
PSR 1706-44	4.24	–	$6.22^{+0.04}_{-0.04}$	$1.9 - 5.8$	$1.8 - 3.2$	$32.48 - 33.08$
PSR 0656+14	5.04	–	$5.71^{+0.03}_{-0.04}$	$7.0 - 8.5$	$0.26 - 0.32$	$32.18 - 32.97$
PSR 0633+1748 (Geminga)	5.53	–	$5.75^{+0.04}_{-0.05}$	$2.7 - 8.7$	$0.123 - 0.216$	$30.85 - 31.51$
PSR 1055-52	5.43	–	$5.92^{+0.02}_{-0.02}$	$6.5 - 19.5$	$0.5 - 1.5$	$32.07 - 33.19$
RX J1856.5-3754	–	$5.70^{+0.05}_{-0.25}$	$5.6 - 5.9$	> 16	$0.105 - 0.129$	$31.44 - 31.68$
RX J0720.4-3125	6.0 ± 0.2	–	$5.55 - 5.95$	$5.0 - 15.0$	$0.1 - 0.3$	$31.3 - 32.5$

(Page et al. astro-ph/0403657)

Not to mix with polar caps heating!



$R_{cap}=R_{NS}(R_{NS}/R_l)^{1/2}$

Might be shifted upwards due to account for hydrogen non-magnetic atmosphere

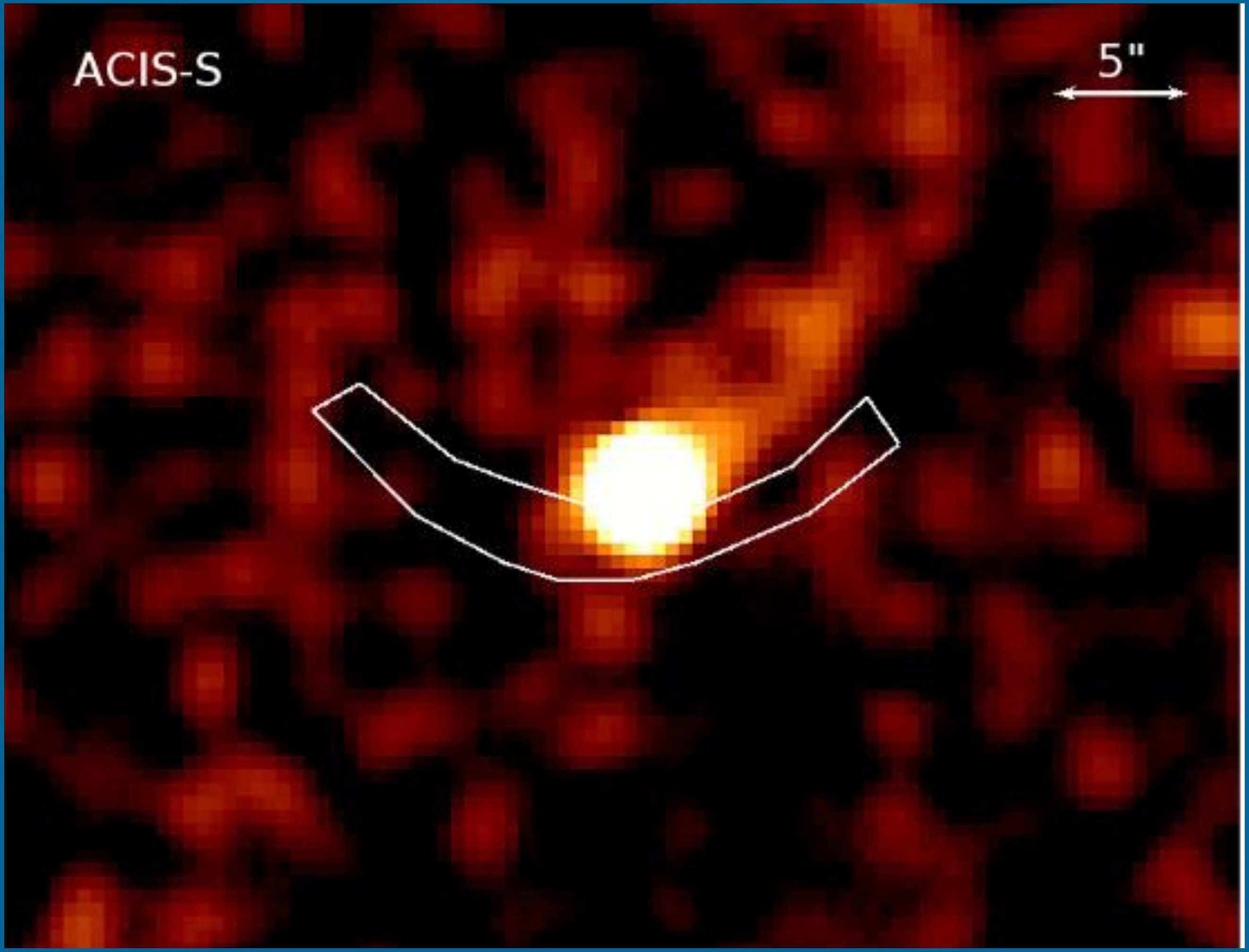
1802.09454

A puzzling source

Millisecond Pulsar J2124−3358

Characteristic age 3.4 Gyr

T~(0.5-2)x10^5K



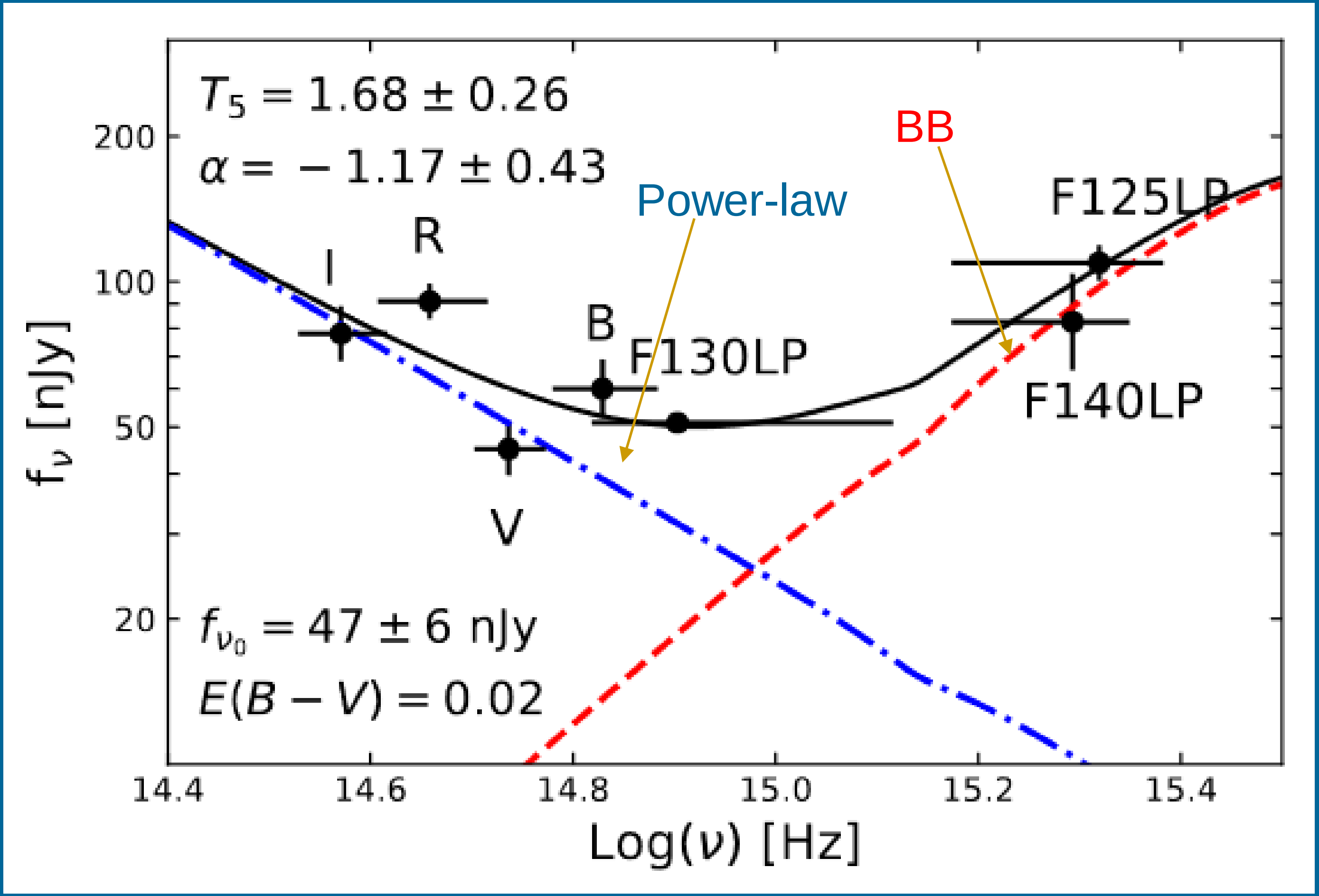
Rangelov et al. (2017)
1701.00002

Another old, but hot

PSR B0950+08

Characteristic age 17.5 Myr

T~(1-3)x10^5K



1710.06448

Rotochemical heating

Contraction due to spin down.
Thus, no beta-equilibrium.

$$C\frac{dT^\infty}{dt} = -L_\nu^\infty - L_\gamma^\infty + L_H^\infty,$$

$$L_H^\infty = \sum_{\ell=e,\mu}\sum_{N=n,p}\int dV\,\eta_\ell\cdot\Delta\Gamma_{M,N\ell}\,e^{2\Phi(r)}$$

$$\eta_\ell \equiv \mu_n - \mu_p - \mu_\ell$$

$$\ell = e,\mu.$$

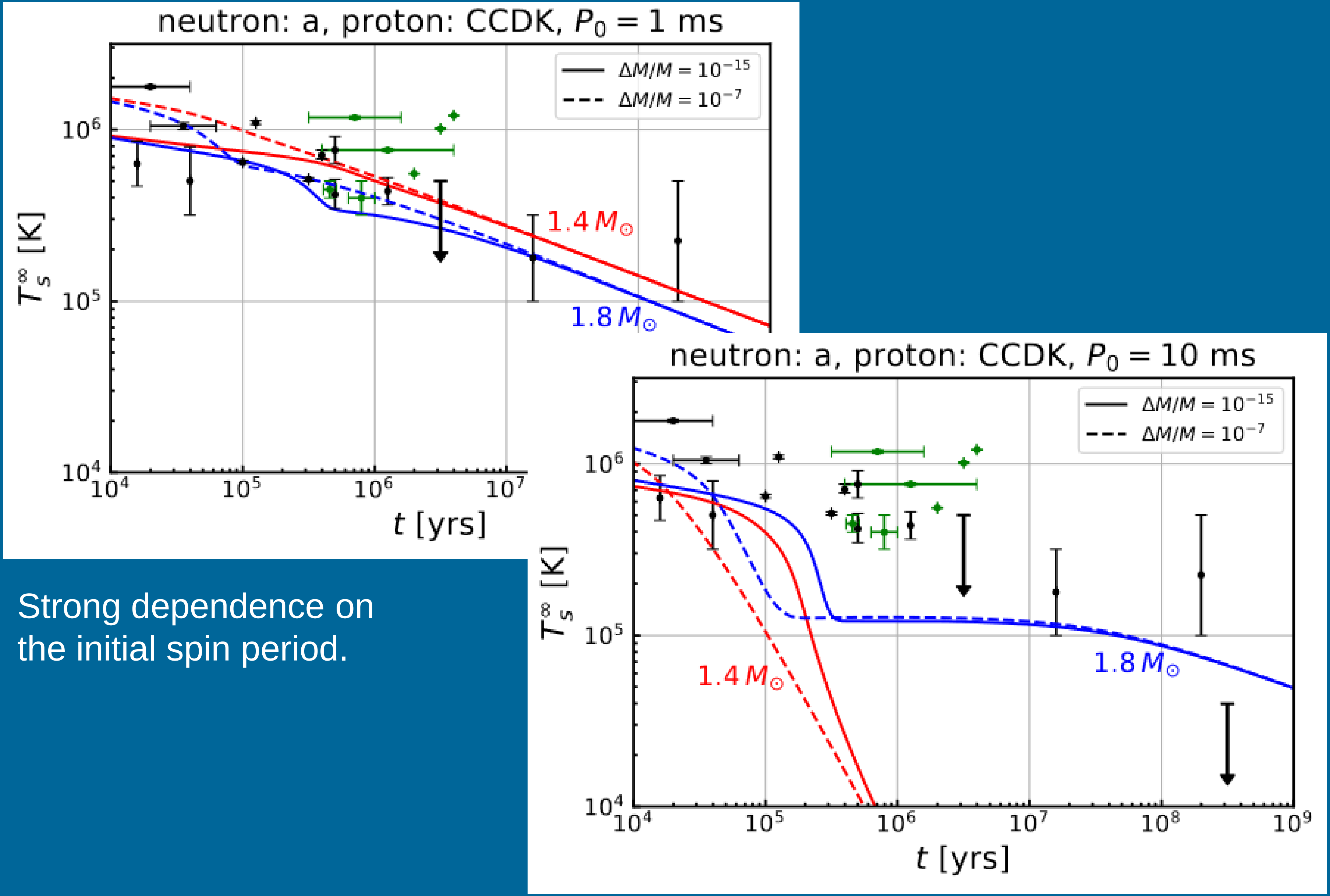
$$\Delta\Gamma_{M,N\ell} = \int\left[\prod_{j=1}^{4}\frac{d^3p_j}{(2\pi)^3}\right]\frac{d^3p_\ell}{(2\pi)^3}\frac{d^3p_\nu}{(2\pi)^3}(2\pi)^4\delta^4(P_f - P_i)\cdot\frac{1}{2}\sum_{\rm spin}|\mathcal{M}_{M,N\ell}|^2$$
$$\times\left[f_1 f_2(1-f_3)(1-f_4)(1-f_\ell) - (1-f_1)(1-f_2)f_3 f_4 f_\ell\right] .$$

$$\frac{d\eta_e^\infty}{dt} = -\sum_{N=n,p}\int dV\,(Z_{npe}\Delta\Gamma_{M,Ne} + Z_{np}\Delta\Gamma_{M,N\mu})\,e^{\Phi(r)} + 2W_{npe}\Omega\dot{\Omega},$$

$$\frac{d\eta_\mu^\infty}{dt} = -\sum_{N=n,p}\int dV\,(Z_{np}\Delta\Gamma_{M,Ne} + Z_{np\mu}\Delta\Gamma_{M,N\mu})\,e^{\Phi(r)} + 2W_{np\mu}\Omega\dot{\Omega},$$

$$\eta_\ell^\infty \equiv \eta_\ell e^{\Phi(r)}$$

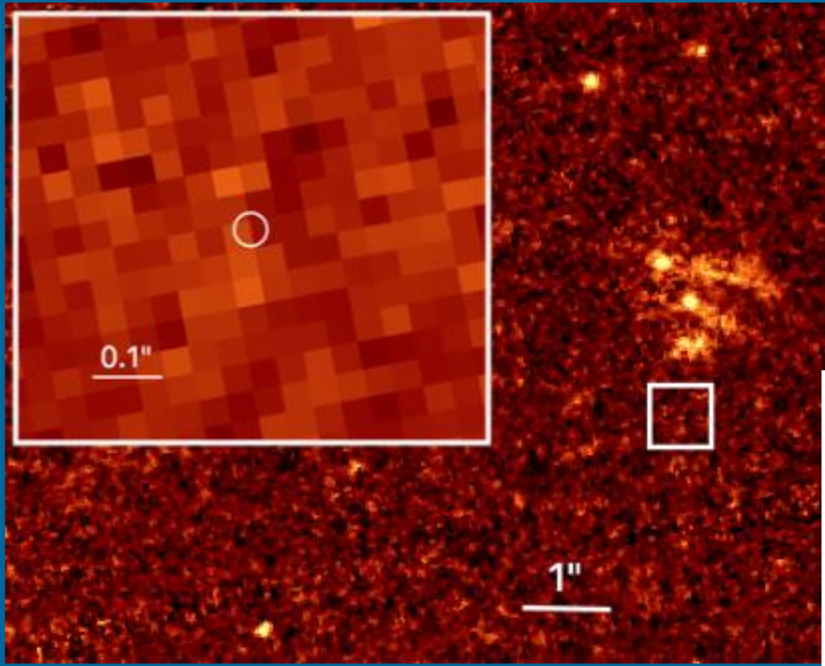
1904.04667



Strong dependence on the initial spin period.

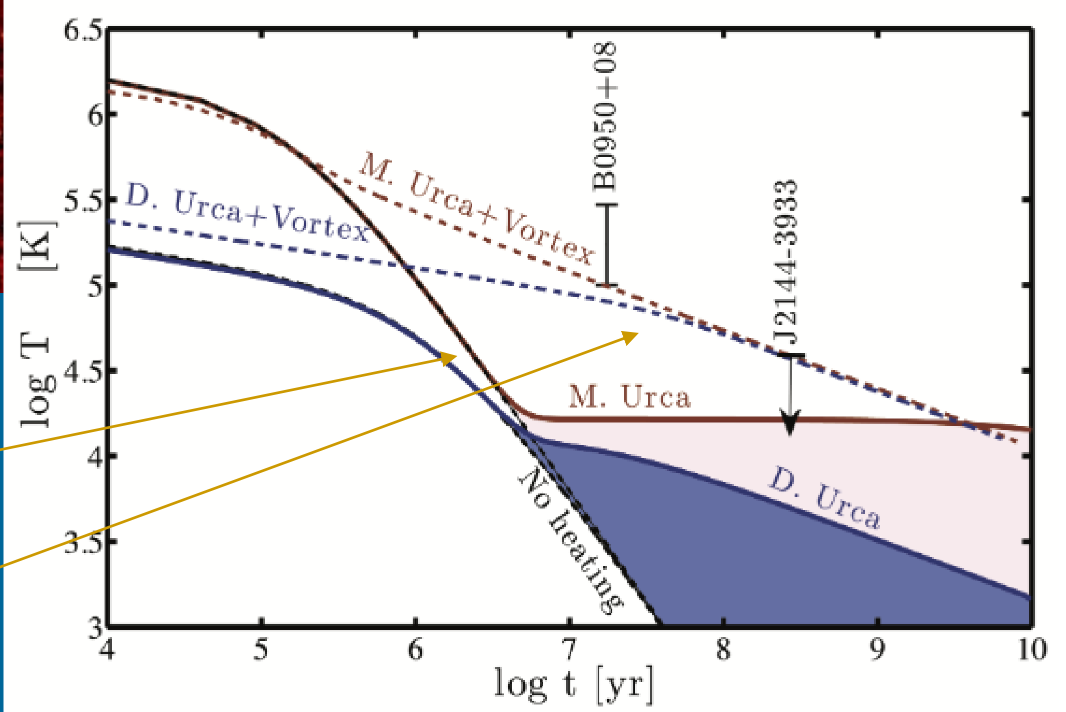
1904.04667

The coldest NS known



PSR J2144–3933
P=8.5 sec
B=2 10^{12} G
d=160-200 pc

Limit: T<42000K



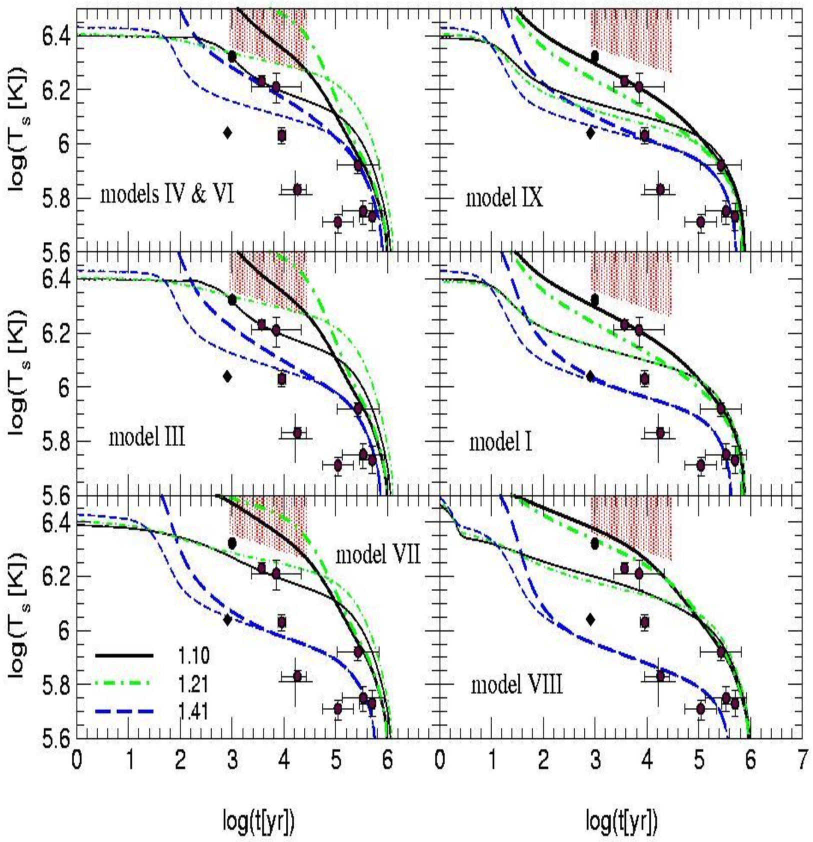
Rotochemical

Rotochemical+vortex friction

1901.07998

Brightness constraint

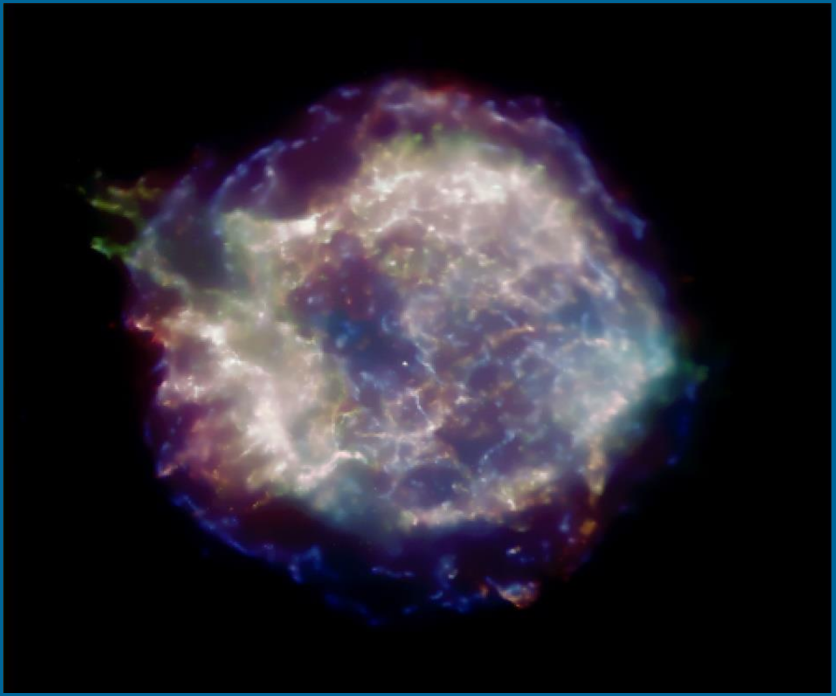
Different tests and constraints
are sensitive to different parameters,
so, typically it is better to use
several different tests



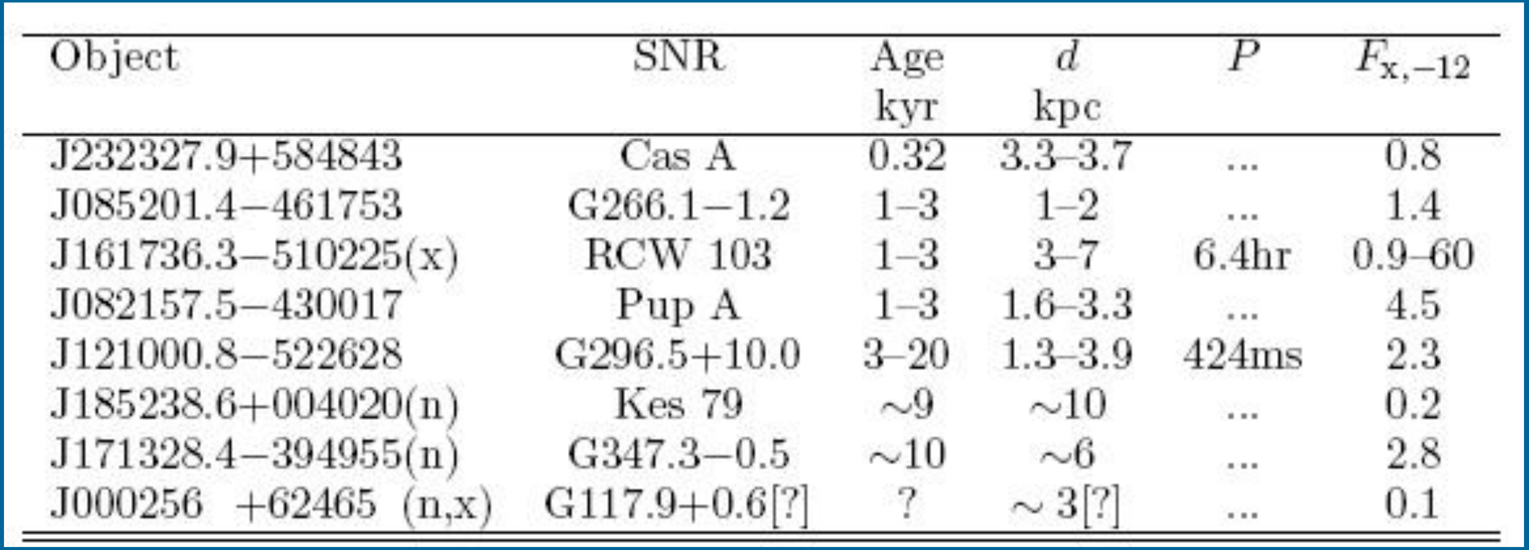
(H. Grigorian astro-ph/0507052)

CCOs

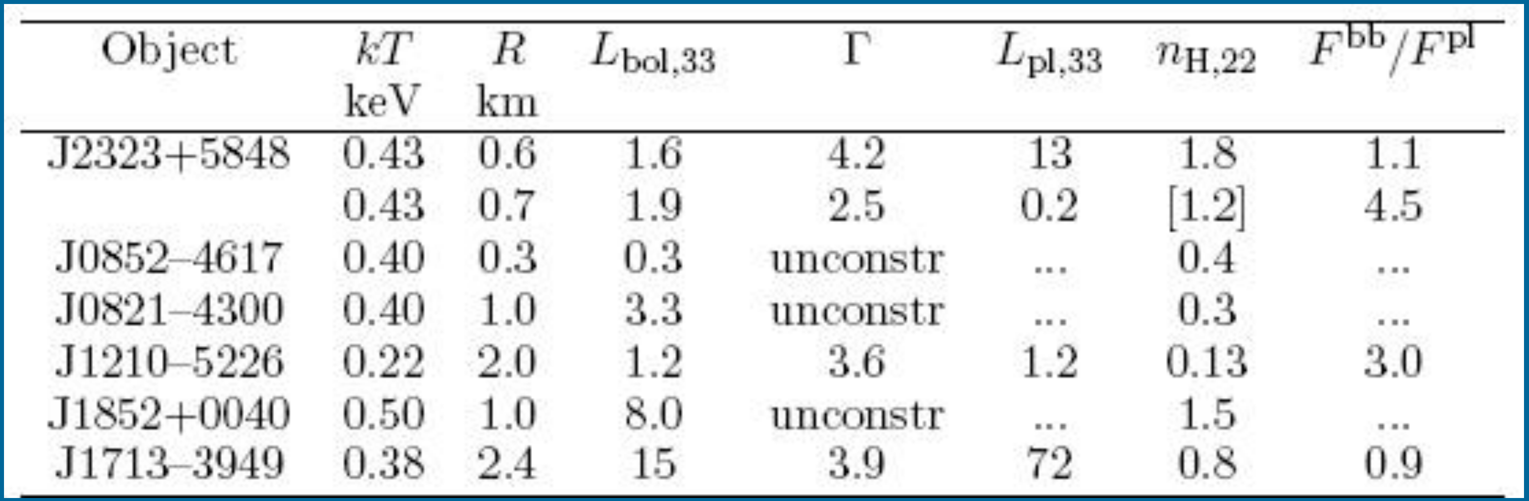
1. Found in SNRs
2. Have no radio or gamma-ray counterpats
3. No pulsar wind nebula (PWN)
4. Have soft thermal-like spectra



Known objects



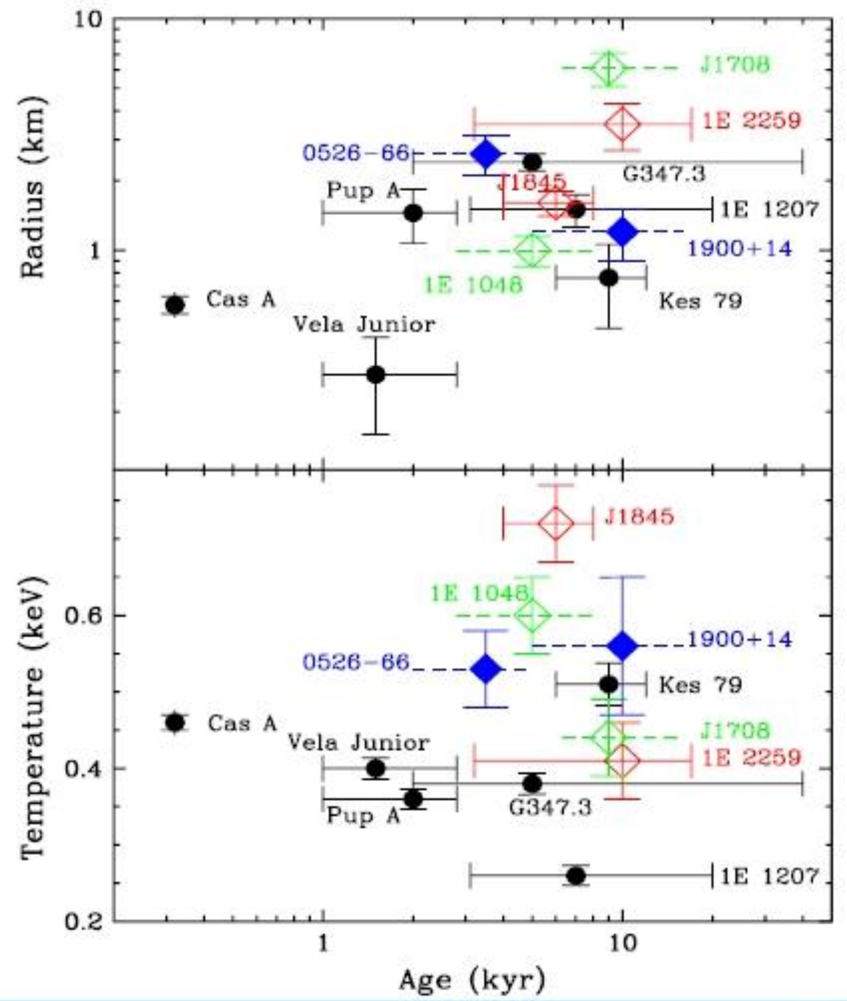
Object	SNR	Age kyr	d kpc	P	$F_{\mathrm{x},-12}$
J232327.9+584843	Cas A	0.32	3.3–3.7	...	0.8
J085201.4−461753	G266.1−1.2	1–3	1–2	...	1.4
J161736.3−510225(x)	RCW 103	1–3	3–7	6.4hr	0.9–60
J082157.5−430017	Pup A	1–3	1.6–3.3	...	4.5
J121000.8−522628	G296.5+10.0	3–20	1.3–3.9	424ms	2.3
J185238.6+004020(n)	Kes 79	$\sim$9	$\sim$10	...	0.2
J171328.4−394955(n)	G347.3−0.5	$\sim$10	$\sim$6	...	2.8
J000256 +62465 (n,x)	G117.9+0.6[?]	?	~ 3[?]	...	0.1

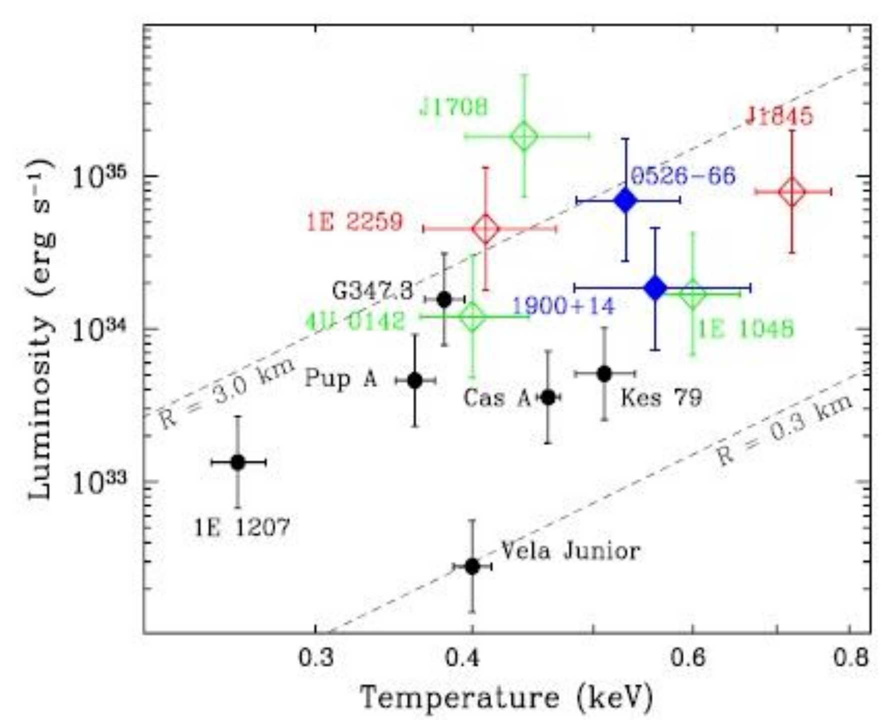


Object	kT keV	R km	$L_{\mathrm{bol},33}$	Γ	$L_{\mathrm{pl},33}$	$n_{\mathrm{H},22}$	$F^{\mathrm{bb}}/F^{\mathrm{pl}}$
J2323+5848	0.43	0.6	1.6	4.2	13	1.8	1.1
	0.43	0.7	1.9	2.5	0.2	[1.2]	4.5
J0852–4617	0.40	0.3	0.3	unconstr	...	0.4	...
J0821–4300	0.40	1.0	3.3	unconstr	...	0.3	...
J1210–5226	0.22	2.0	1.2	3.6	1.2	0.13	3.0
J1852+0040	0.50	1.0	8.0	unconstr	...	1.5	...
J1713–3949	0.38	2.4	15	3.9	72	0.8	0.9

(Pavlov et al. astro-ph/0311526)

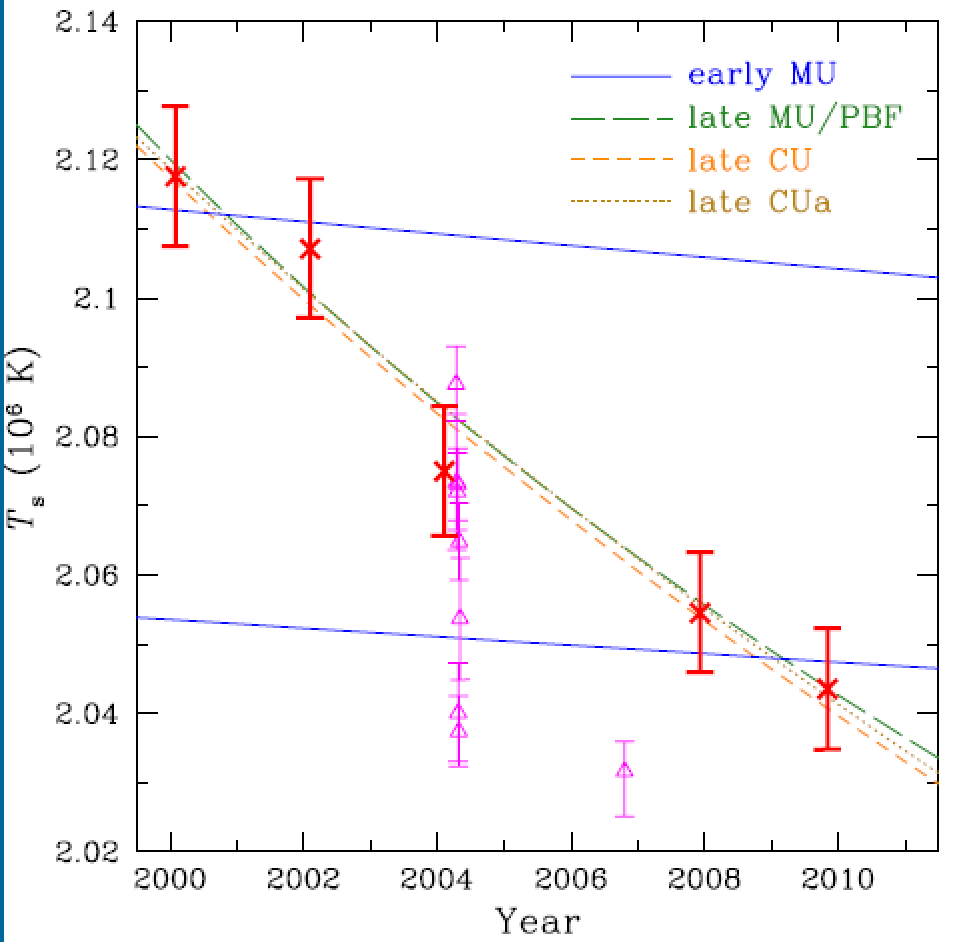
Correlations





(Pavlov et al. astro-ph/0311526)

Cas A peculiar cooling

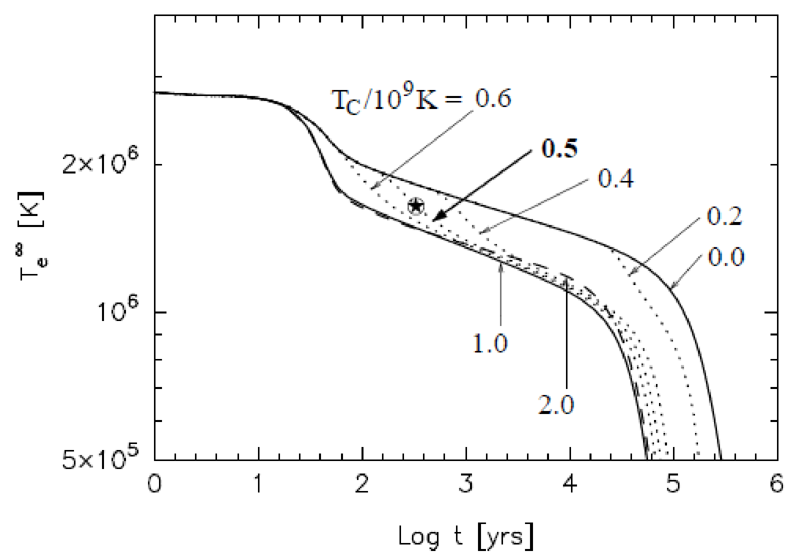


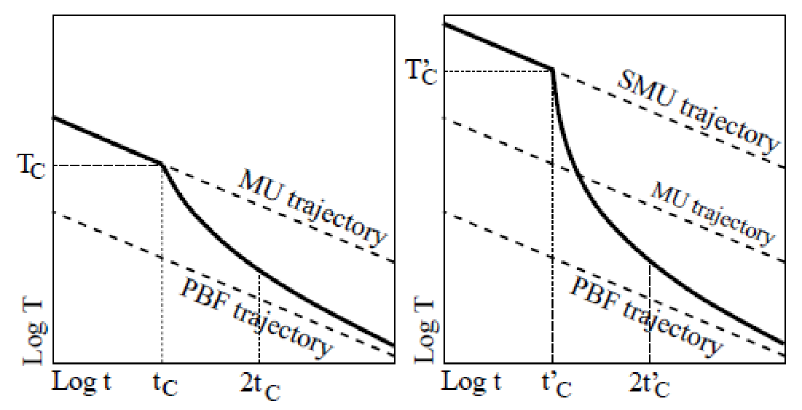
330 years
~3.5 kpc
Carbon atmosphere
The youngest cooler known

Temperature steadily goes down
by ~4% in 10 years:
2.12 10^6K in 2000 – 2.04 10^6K in 2009

1007.4719

Onset of neutron $^3P_2$ superfluidity in the core





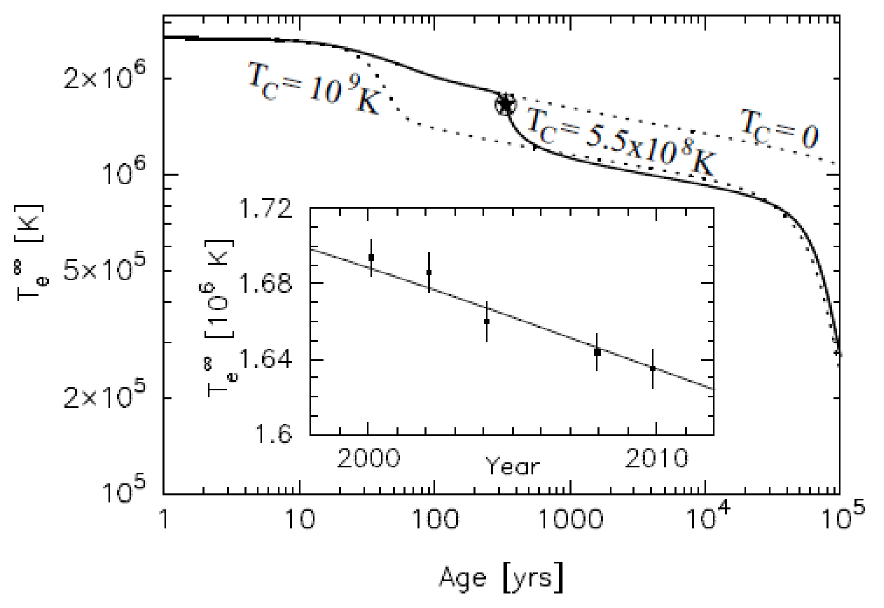
The idea is that we see the result of the onset of neutron $^3P_2$ superfluidity in the core.

The NS just cooled down enough to have this type of neutron superfluidity in the core.

This gives an opportunity to estimate the critical temperature: 0.5 10^9 K

1011.6142

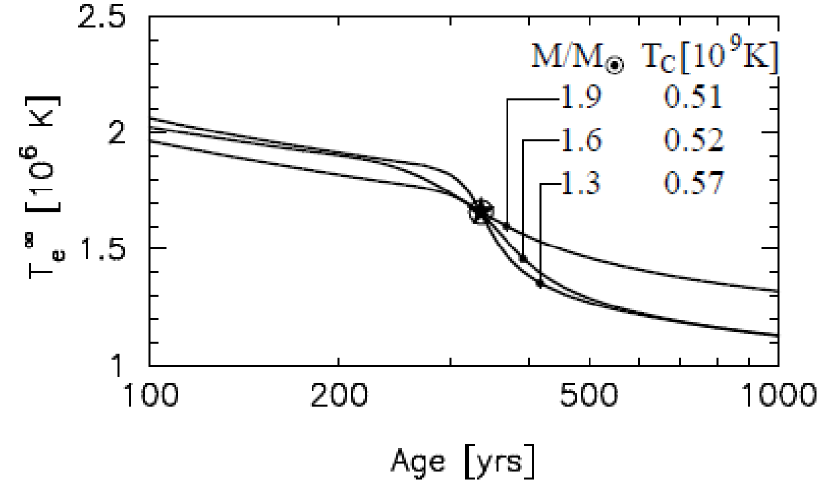
The best fit model



To explain a quick cooling it is necessary to assume suppression of cooling by proton $^1S_0$ superfluidity in the core.

Rapid cooling will proceed for several tens of years more.

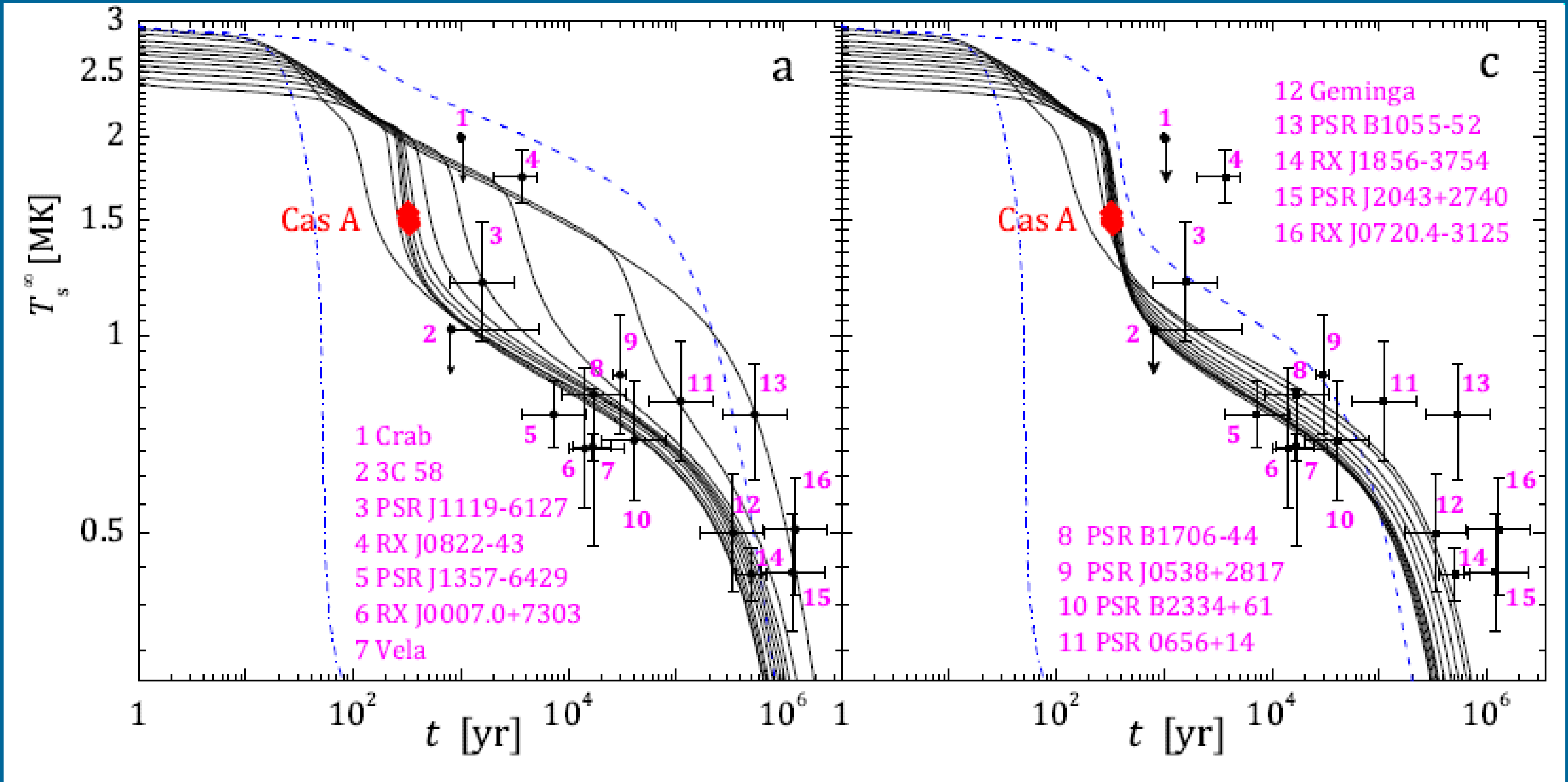
The plot is made for $M=1.4M_O$



Cooling curves depend on masses, but the estimate of the critical temper. depends on M just slightly.

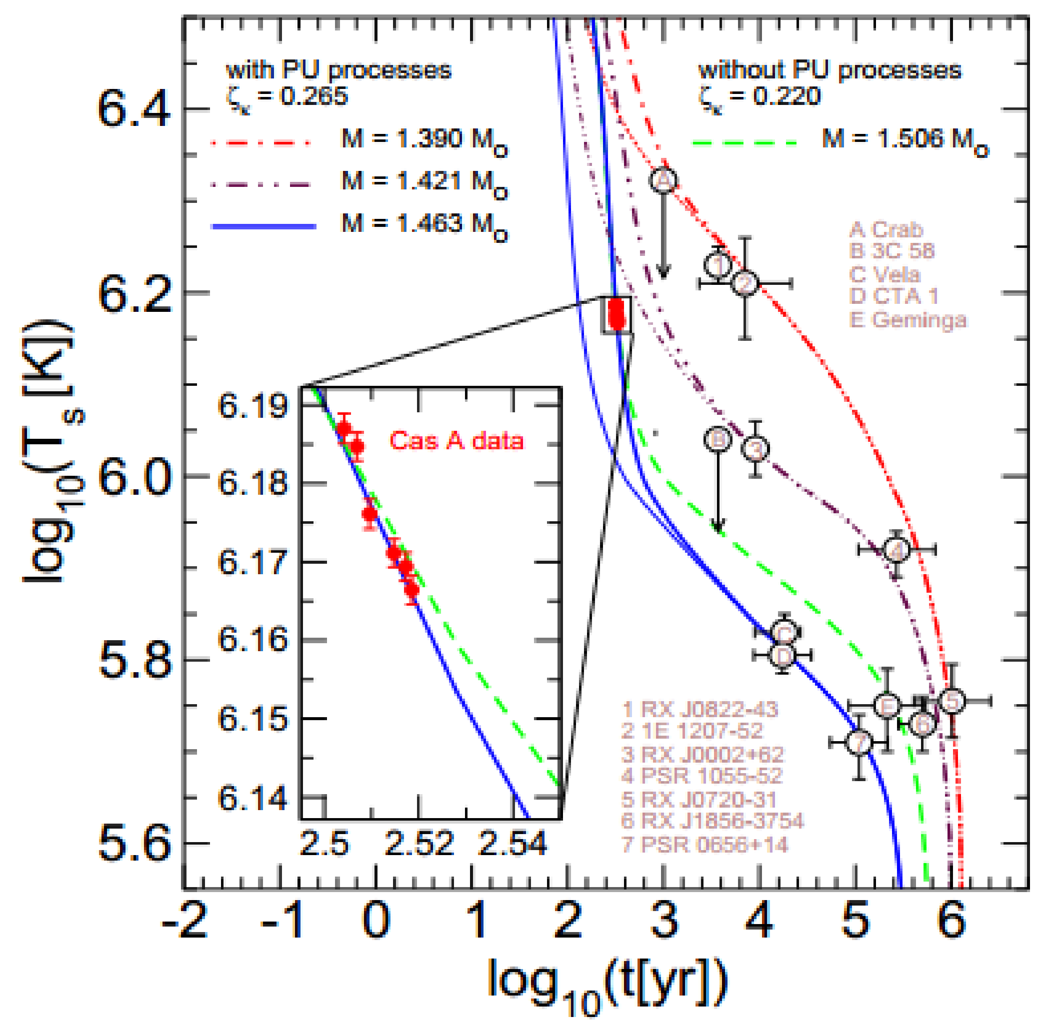
1011.6142, see many details in 1110.5116

Different superfluidity models



1012.0045

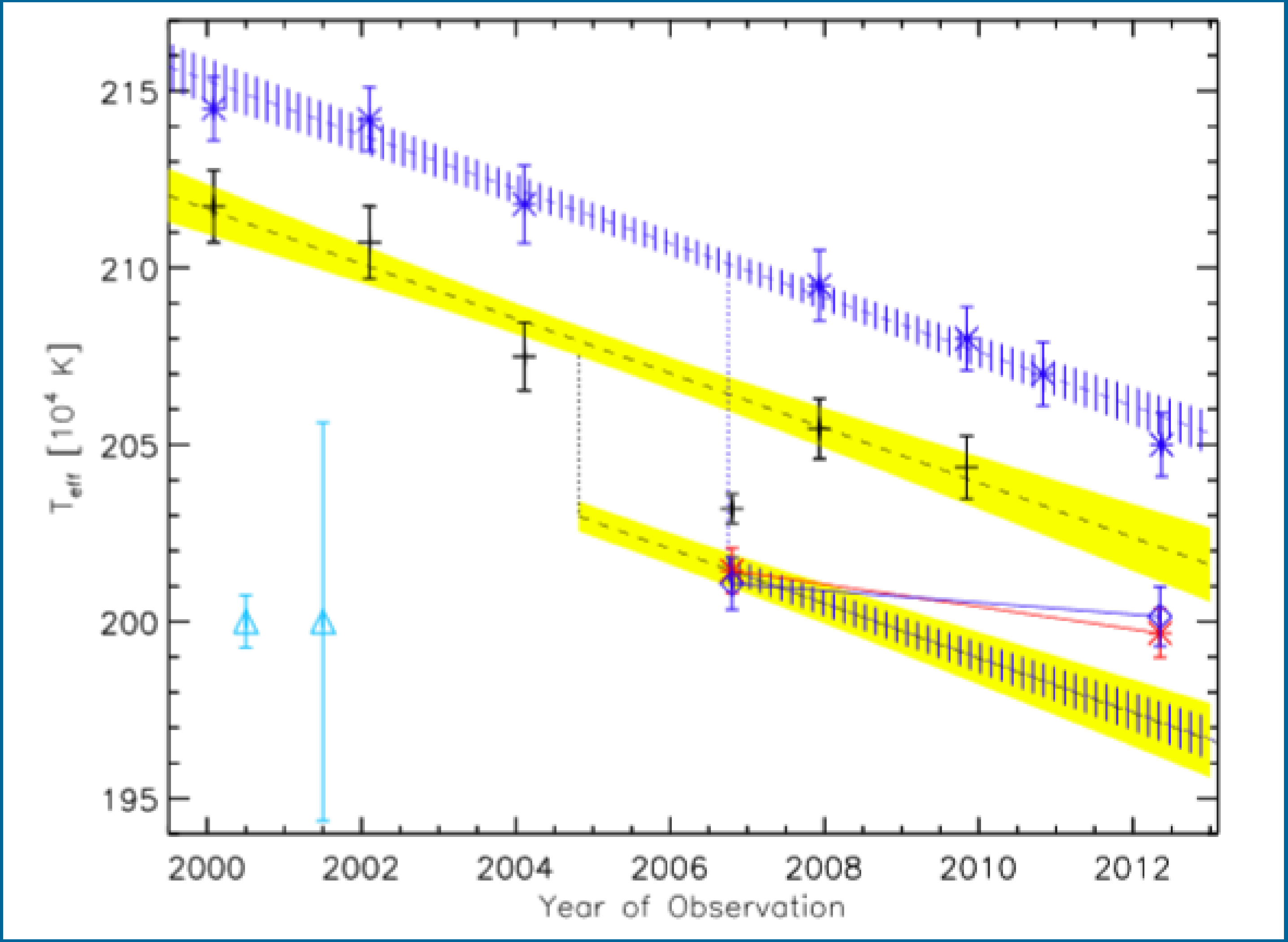
Nuclear medium cooling



Crucial for the successful
description of
the observed data
is a substantial reduction
of the thermal conductivity,
resulting from
a suppression of both
the electron and
nucleon contributions
to it by medium effects.

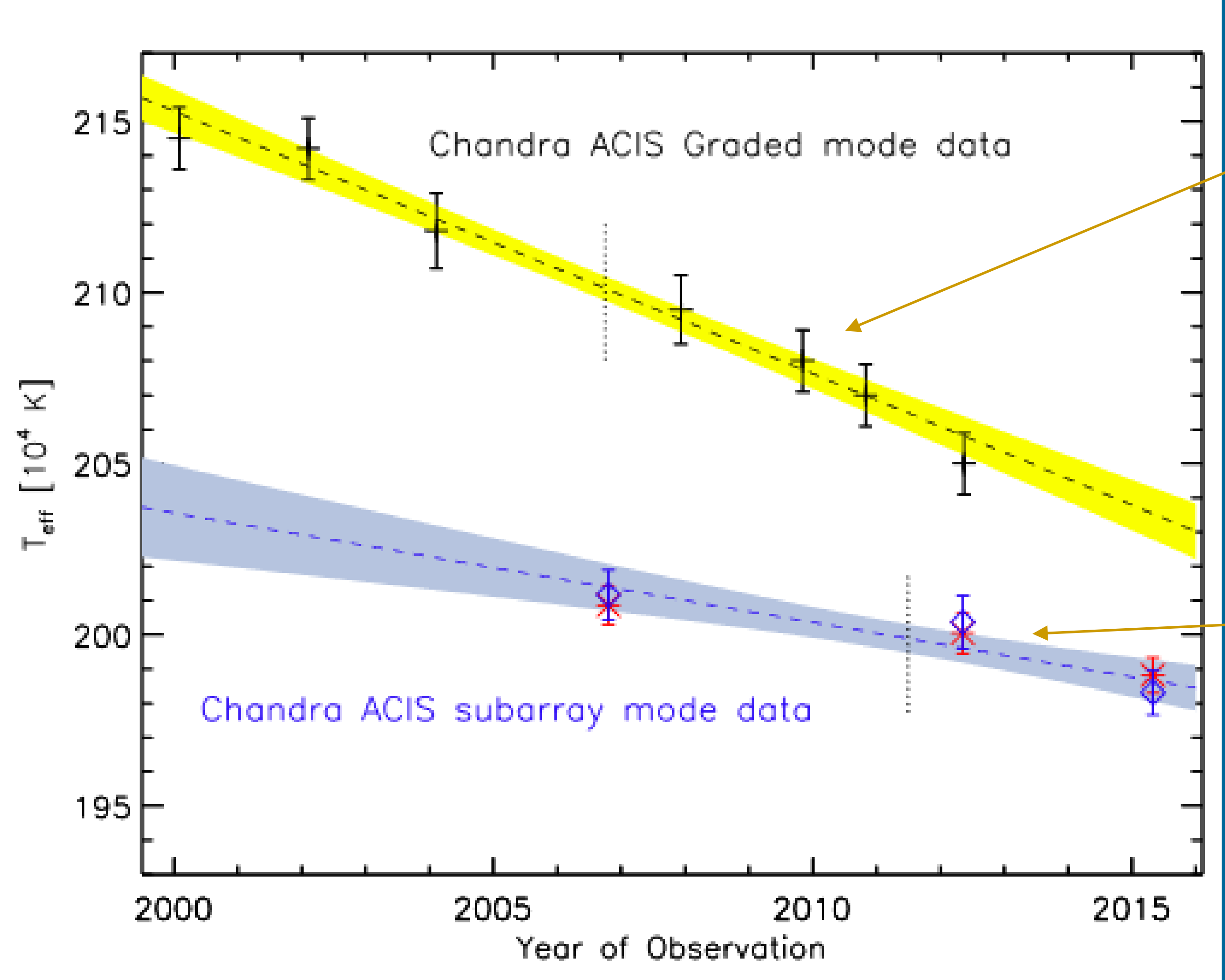
1108.4125

New twist: no cooling!



1311.0888

New data: still no cooling?

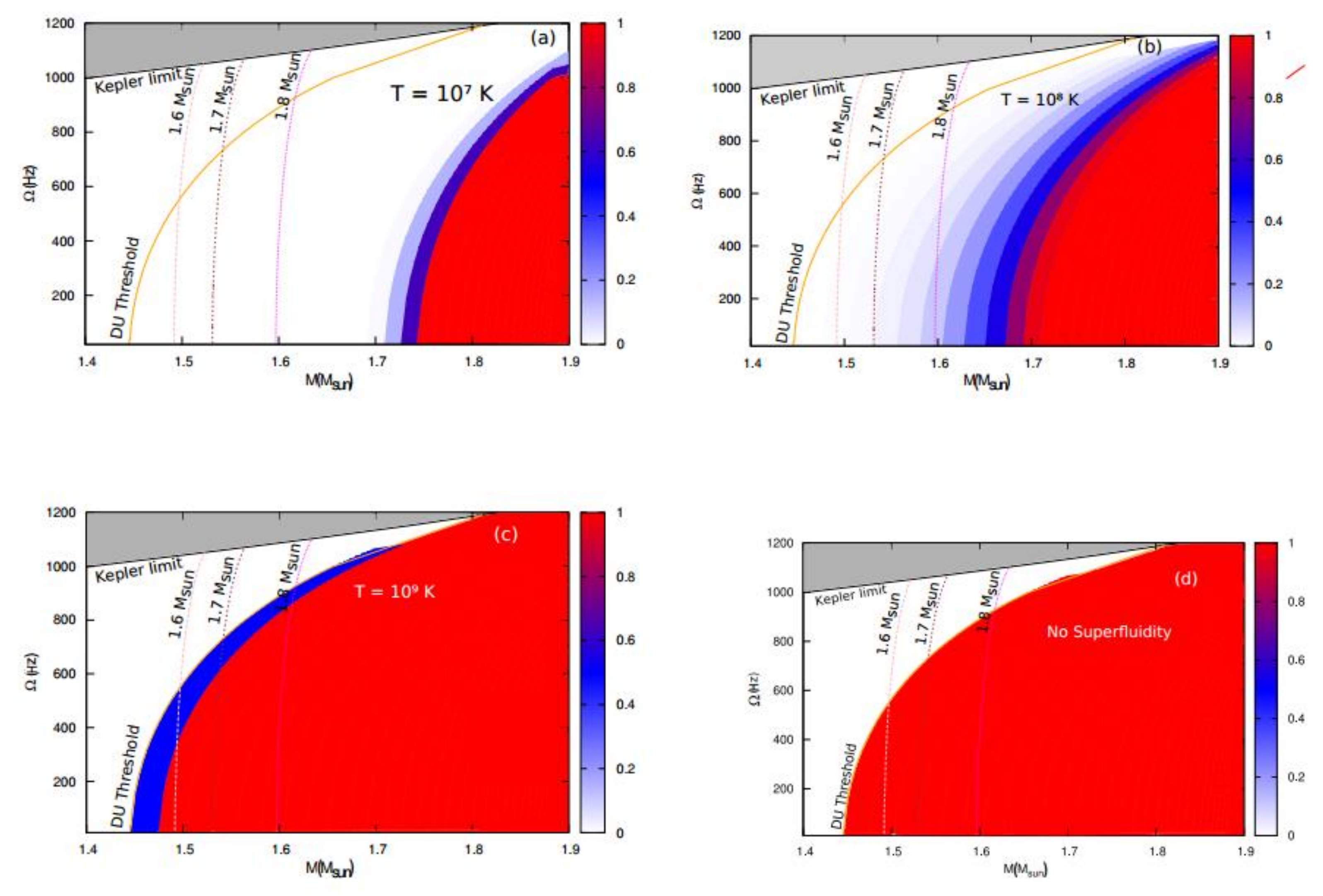


Elshamouty et al. 2013

Posselt et al.

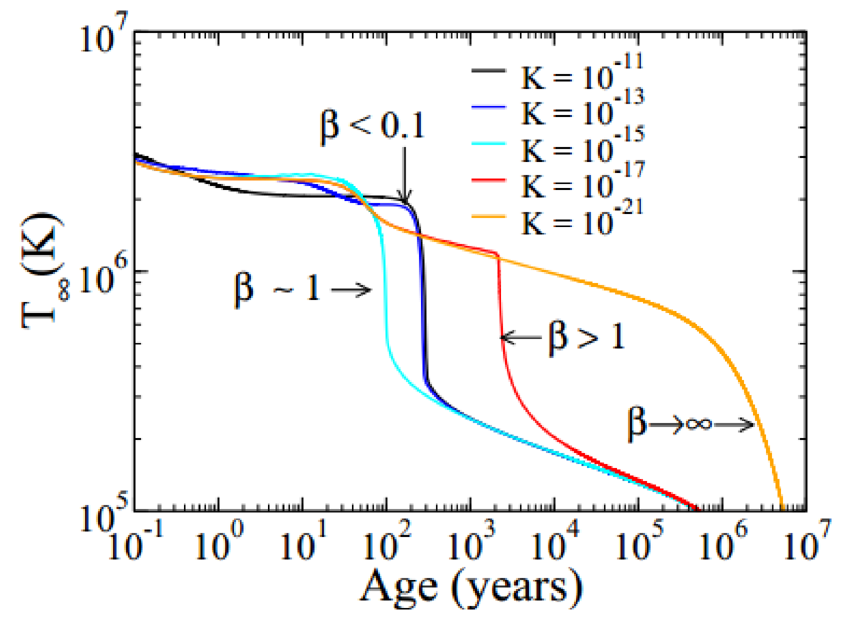
1808.00531

Cooling and rotation

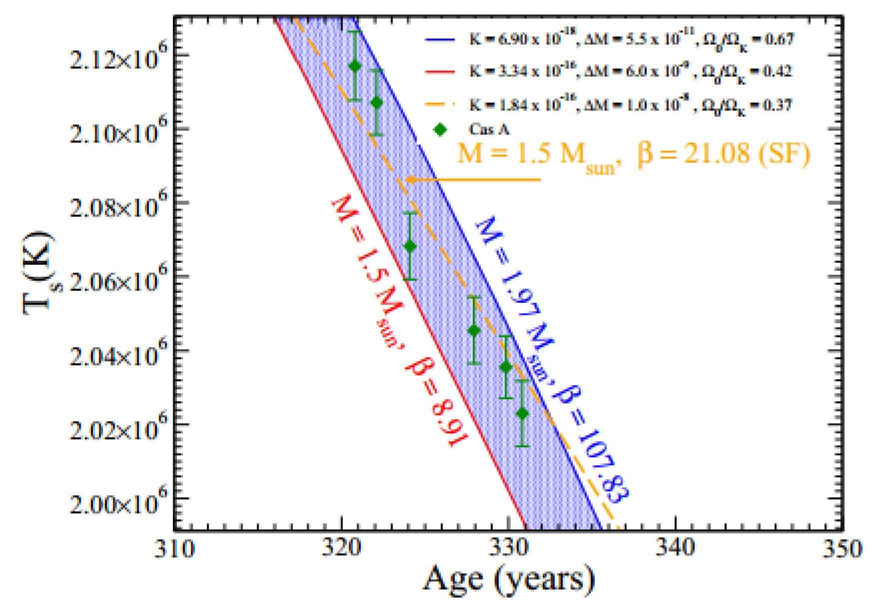


1103.3870

Cas A case



P$_0$=0.0025-0.00125 sec
B~10^{11} G

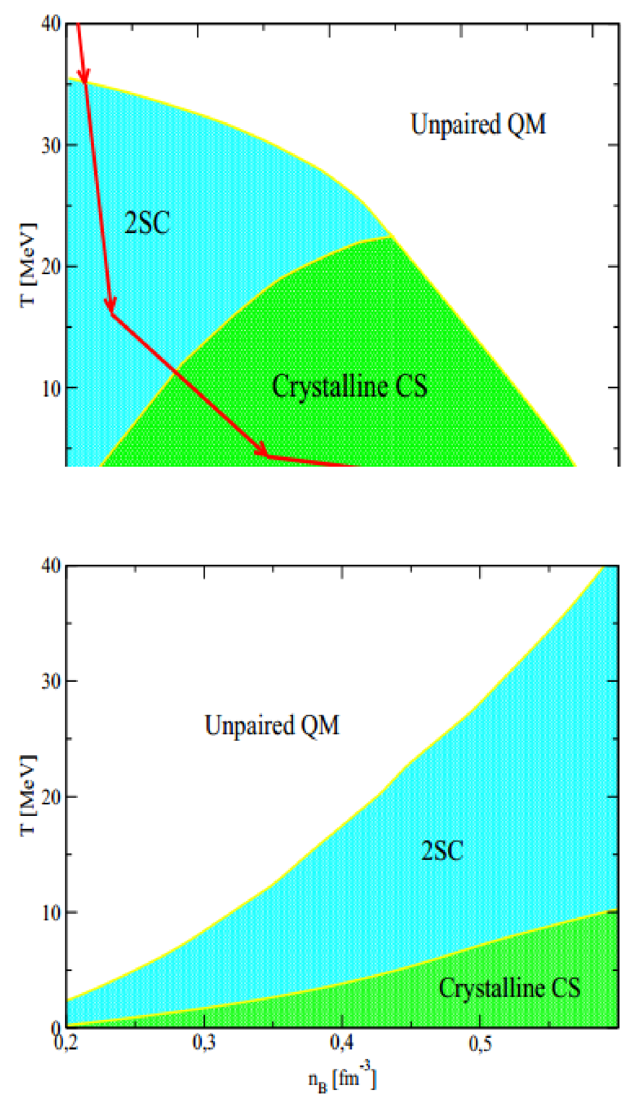


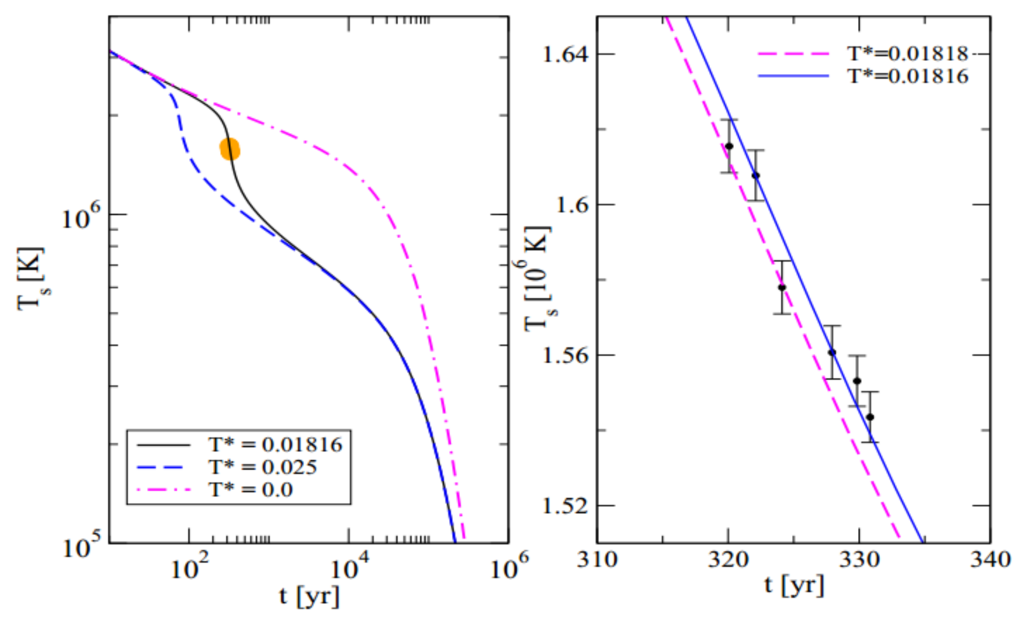
Other studies of the influence of effects of rotation see in 1201.2381

1103.3870

Exotic phase transition

Rapid cooling of Cas A can be understood
as a phase transition from the perfect 2SC phase
to a crystalline/gapless color-superconducting state

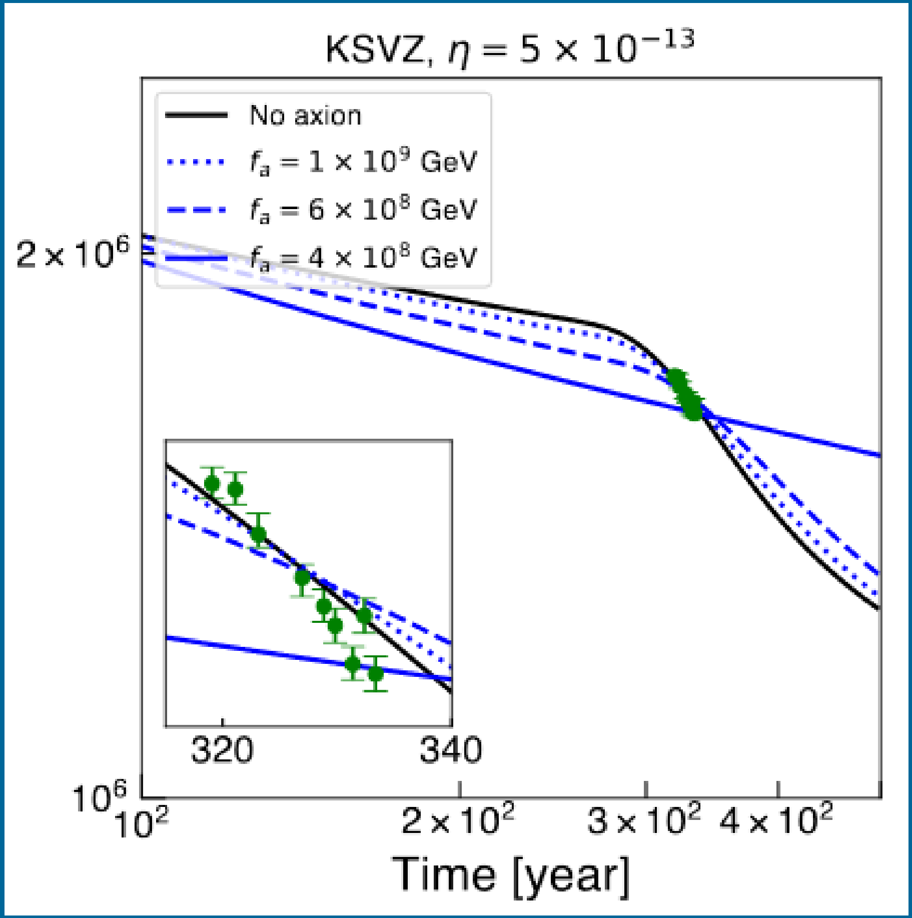




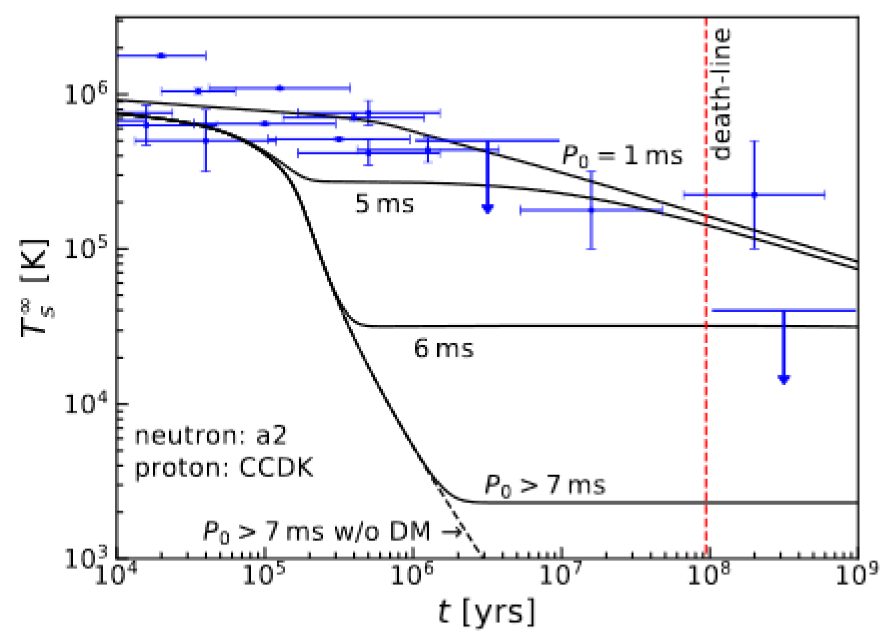
1301.2814

Physics behind Standard Model

Axions cooling and dark matter accumulation and annihilation



If dark matter is accumulated and there is annihilation, then the surface redshifted temperature does not fall down below 2200K.



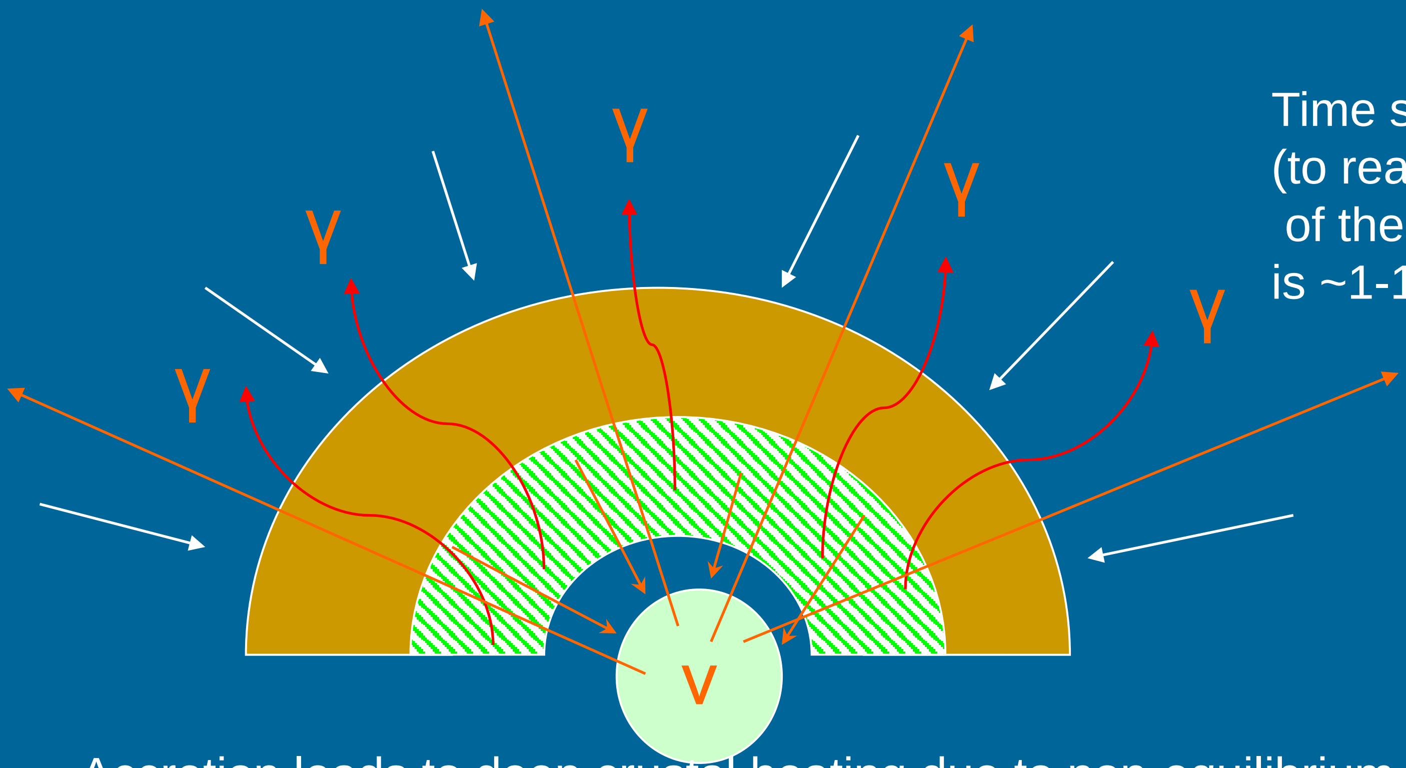
2003.08199

Cooling of X-ray transients

“Many neutron stars in close X-ray binaries are transient accretors (transients);
They exhibit X-ray bursts separated by long periods (months or even years) of quiescence.
It is believed that the quiescence corresponds to a low level, or even halted, accretion onto the neutron star.
During high-state accretion episodes,
the heat is deposited by nonequilibrium processes in the deep layers (10^{12} -10^{13} g cm^{-3}) of the crust.
This deep crustal heating can maintain the temperature of the neutron star interior at a sufficiently high level to explain a persistent thermal X-ray radiation in quiescence (Brown *et al.*, 1998).”

(quotation from the book by Haensel, Potekhin, Yakovlev)

Deep crustal heating and cooling



Time scale of cooling
(to reach thermal equilibrium
of the crust and the core)
is ~1-100 years.

To reach the
state “before”
takes ~10^3-10^4 yrs

Accretion leads to deep crustal heating due to non-equilibrium nuclear reactions.
After accretion is off:

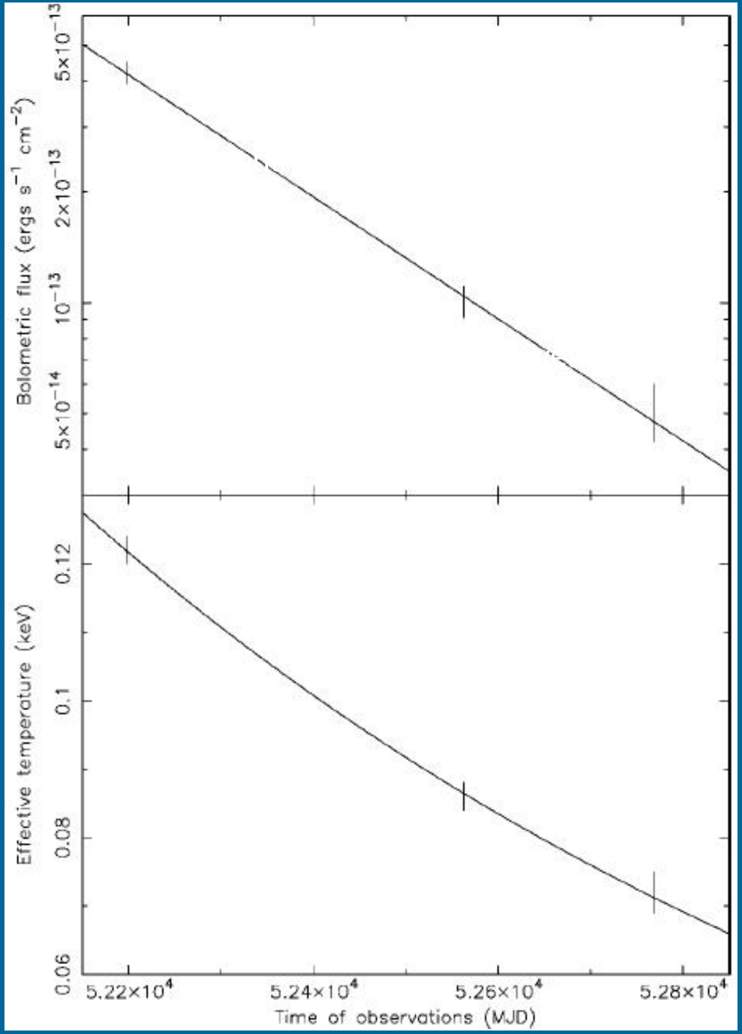
- heat is transported inside and emitted by neutrinos
- heat is slowly transported out and emitted by photons

$\rho \sim 10^{12}$-10^{13} g/cm^3

See, for example, Haensel, Zdunik arxiv:0708.3996

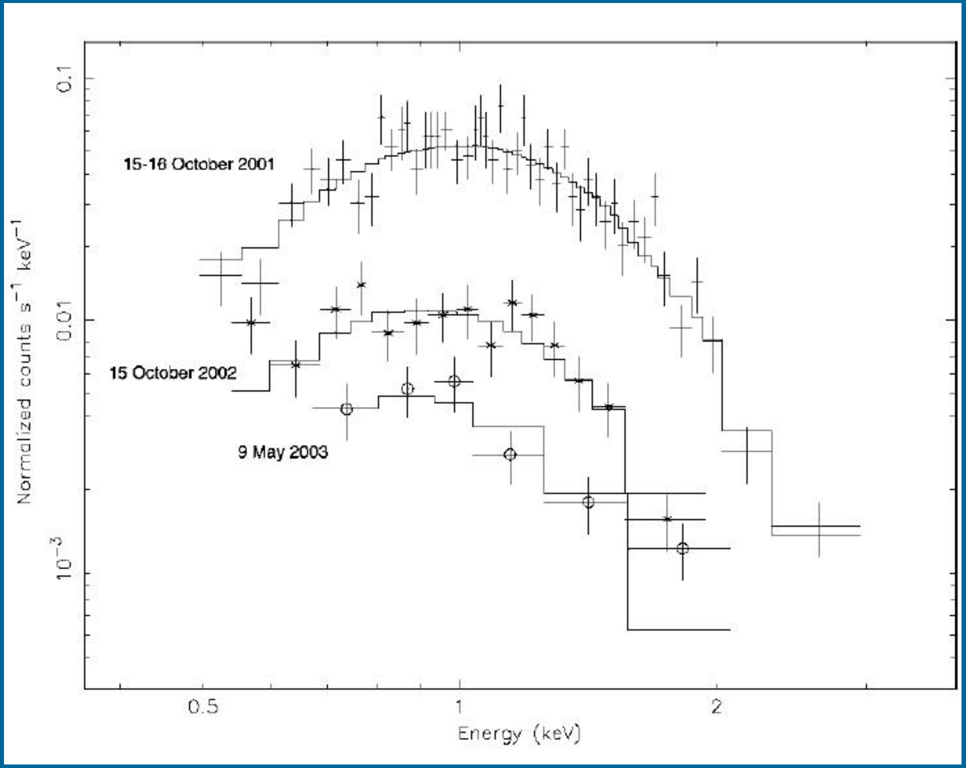
New calculations appeared very recently 0811.1791 Gupta et al.

Cooling in soft X-ray transients



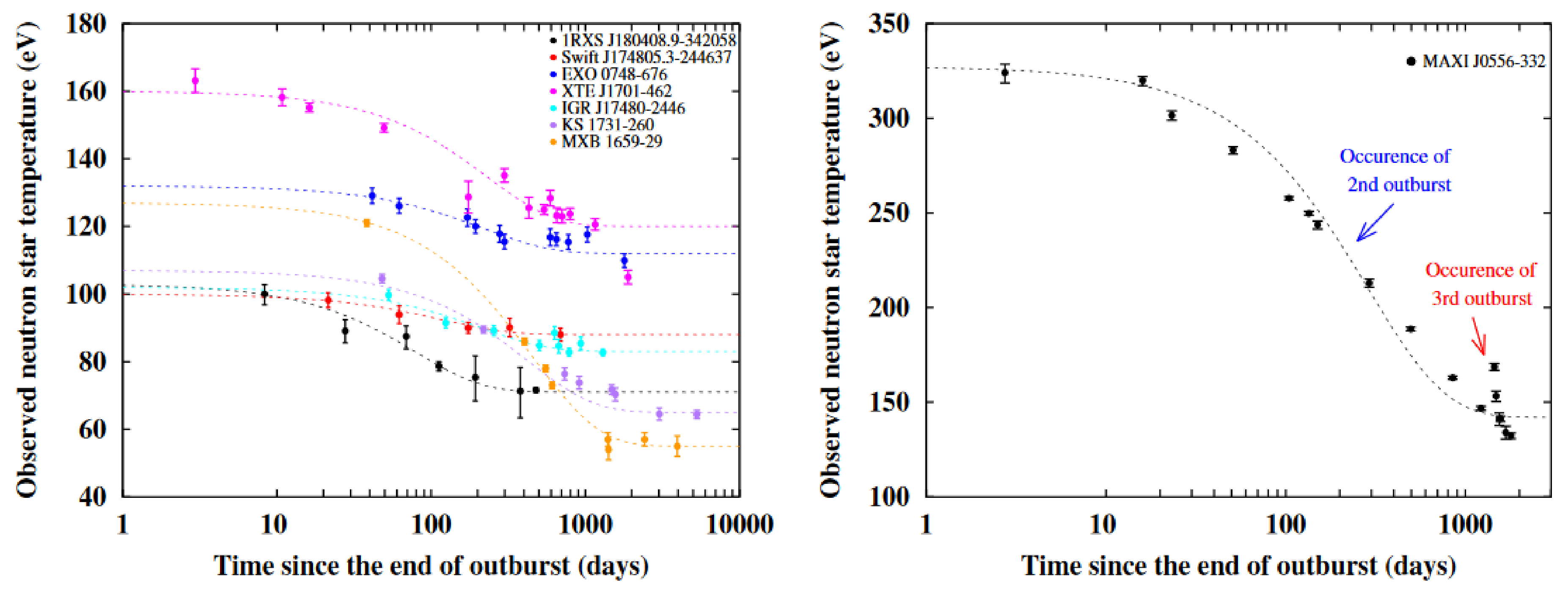
MXB 1659-29

~2.5 years outburst



[Wijnands et al. 2004]

Fitting cooling of known sources



Different systems allow to probe different regimes of cooling and different layers of the crust.

1709.07034

Pycnonuclear reactions

Let us give an example from Haensel, Zdunik (1990)

We start with ^{56}Fe
Density starts to increase

$^{56}\mathrm{Fe} \rightarrow {}^{56}\mathrm{Cr}$
$^{56}\mathrm{Fe} + e^- \rightarrow {}^{56}\mathrm{Mn} + v_e$
$^{56}\mathrm{Mn} + e^- \rightarrow {}^{56}\mathrm{Cr} + v_e$

At ^{56}Ar: neutron drip
$^{56}\mathrm{Ar} + e^- \rightarrow {}^{56}\mathrm{Cl} + v_e$
$^{56}\mathrm{Cl} \rightarrow {}^{55}\mathrm{Cl} + n$
$^{55}\mathrm{Cl} + e^- \rightarrow {}^{55}\mathrm{S} + v_e$
$^{55}\mathrm{S} \rightarrow {}^{54}\mathrm{S} + n$
$^{54}\mathrm{S} \rightarrow {}^{52}\mathrm{S} + 2n$

Then from ^{52}S we have a chain:
$^{52}\mathrm{S} \rightarrow {}^{46}\mathrm{Si} + 6n - 2e^- + 2v_e$

As Z becomes smaller
the Coulomb barrier decreases.
Separation between
nuclei decreases, vibrations grow.
$^{40}\mathrm{Mg} \rightarrow {}^{34}\mathrm{Ne} + 6n - 2e^- + 2v_e$

At Z=10 (Ne) pycnonuclear reactions start.

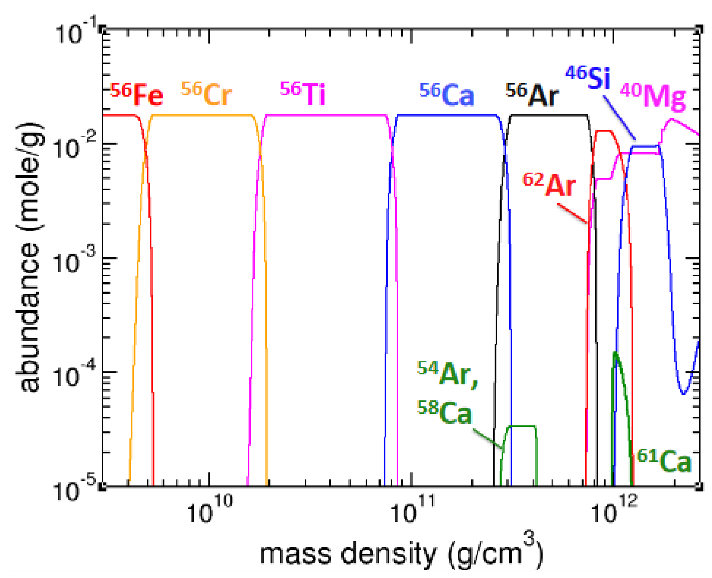
$^{34}\mathrm{Ne} + {}^{34}\mathrm{Ne} \rightarrow {}^{68}\mathrm{Ca}$
$^{36}\mathrm{Ne} + {}^{36}\mathrm{Ne} \rightarrow {}^{72}\mathrm{Ca}$

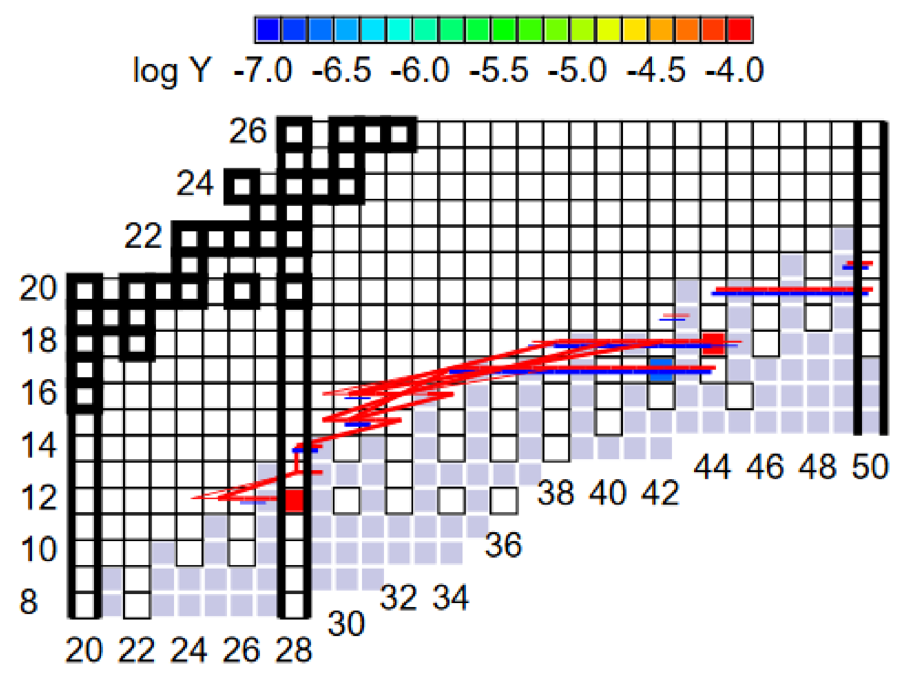
Then a heavy nuclei can react again:
$^{72}\mathrm{Ca} \rightarrow {}^{66}\mathrm{Ar} + 6n - 2e^- + 2v_e$

$^{48}\mathrm{Mg} + {}^{48}\mathrm{Mg} \rightarrow {}^{96}\mathrm{Cr}$
$^{96}\mathrm{Cr} \rightarrow {}^{88}\mathrm{Ti} + 8n - 2e^- + 2v_e$

See a review in 1803.03818

Crust composition and reactions





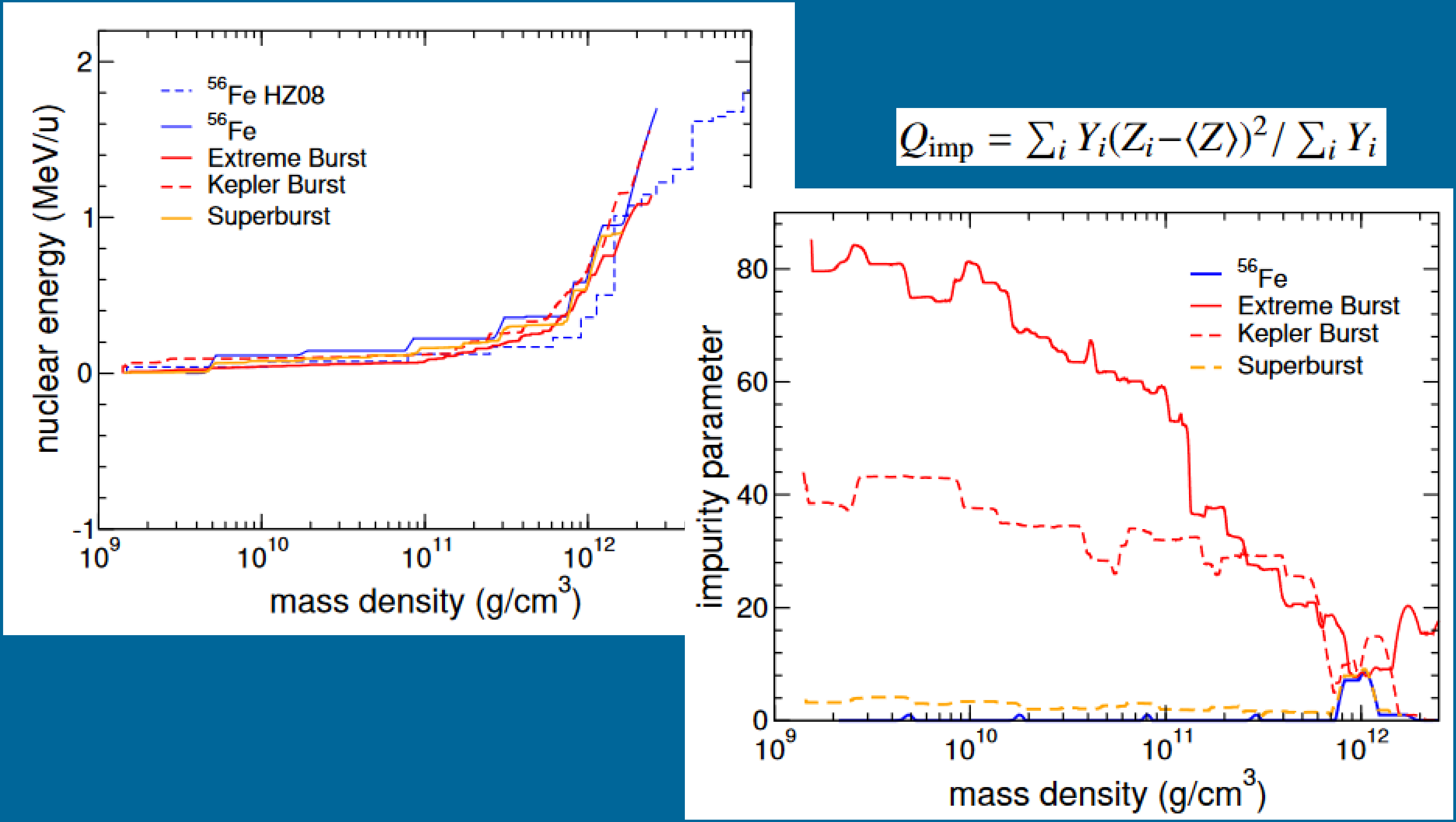
Transition	P^a	ρ^b	μ_e^c	X_n^d
^{56}Fe→^{56}Cr	3.4×10^{27}	4.9×10^{9}	6.2	$<10^{-25}$
^{56}Cr→^{56}Ti	1.7×10^{28}	1.8×10^{10}	9.6	$<10^{-25}$
^{56}Ti→^{56}Ca	1.1×10^{29}	8.1×10^{10}	15.6	$<10^{-25}$
^{56}Ca→^{56}Ar,^{54}Ar,^{58}Ca	5.5×10^{29}	2.9×10^{11}	23.3	1.2×10^{-18}
^{56}Ar,^{54}Ar,^{58}Ca→^{56}Ar	8.3×10^{29}	4.2×10^{11}	25.9	7.2×10^{-20}
^{56}Ar→^{40}Mg,^{62}Ar	1.8×10^{30}	7.8×10^{11}	31.6	5.4×10^{-8}
^{40}Mg,^{62}Ar→^{40}Mg,^{48}Si	2.3×10^{30}	1.1×10^{12}	33.5	0.13
^{40}Mg,^{48}Si→^{40}Mg	4.2×10^{30}	2.8×10^{12}	37.1	0.54

New calculations show that
the model with one-to-one correspondence
nuclei-density is an oversimplication,
2004.00997, 2004.04195

1803.03818

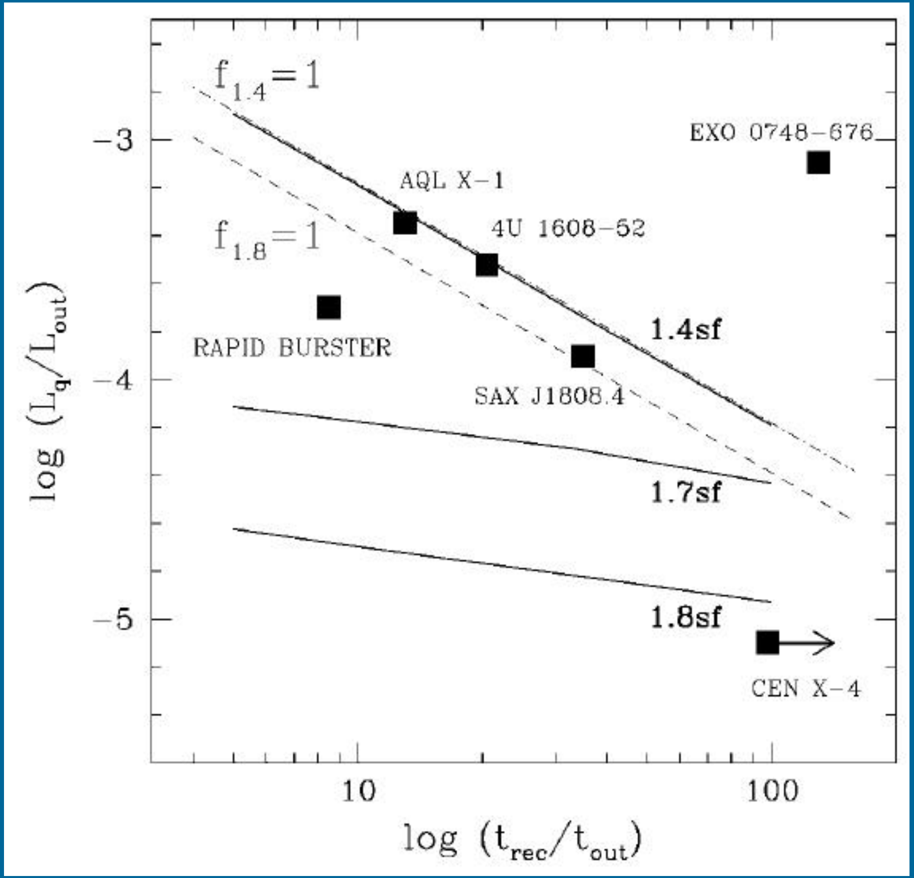
Energy release vs. density and impurity

$$Q_{\text{imp}} = \sum_i Y_i (Z_i - \langle Z \rangle)^2 / \sum_i Y_i$$



1803.03818

A simple model



t_{rec} – time interval between outbursts
t_{out} – duration of an outburst
L_q – quiescent luminosity
L_{out} – luminosity during an outburst

Dashed lines corresponds to the case when all energy is emitted from a surface by photons.

$$L_q \sim \frac{Q_{\rm nuc}}{m_u}\langle \dot{M}\rangle \sim 6\times 10^{32}\,\frac{\langle \dot{M}\rangle}{10^{-11}\,M_\odot\ {\rm yr}^{-1}}\ {\rm ergs\ s^{-1}}$$

Average Mdot = Mdot * $(t_{out}/(t_{rec}+t_{out}))$
L_{out}~Mdot

L_q/L_{out}= (Q/m_u) (Mdot t_{out})/(t_{rec} *L_{out}), $t_{rec}>>t_{out}$

L_q/L_{out}= (Q/m_u) * $(1/10^{20})$ * (t_{out}/t_{rec})

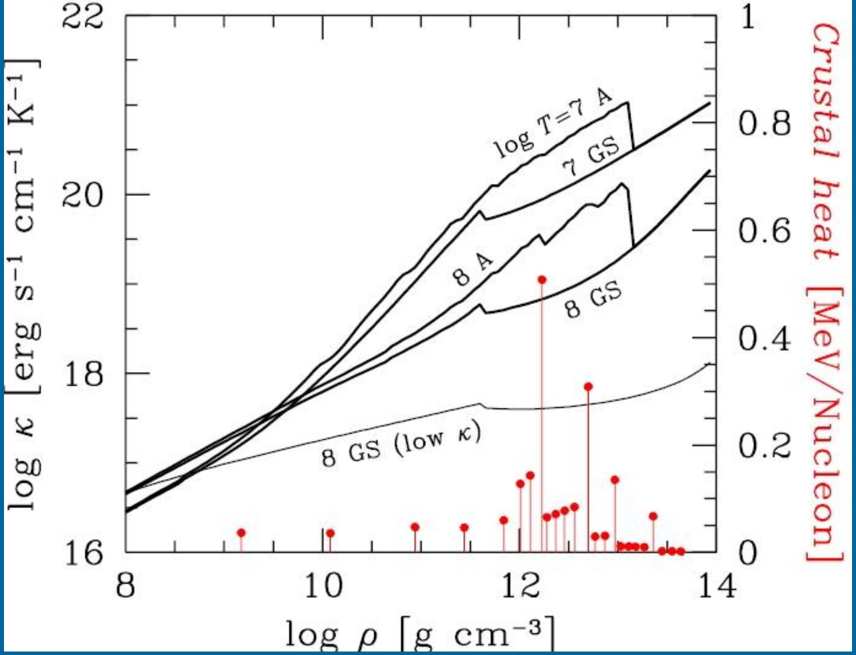
[Colpi et al. 2001] astro-ph/0010572

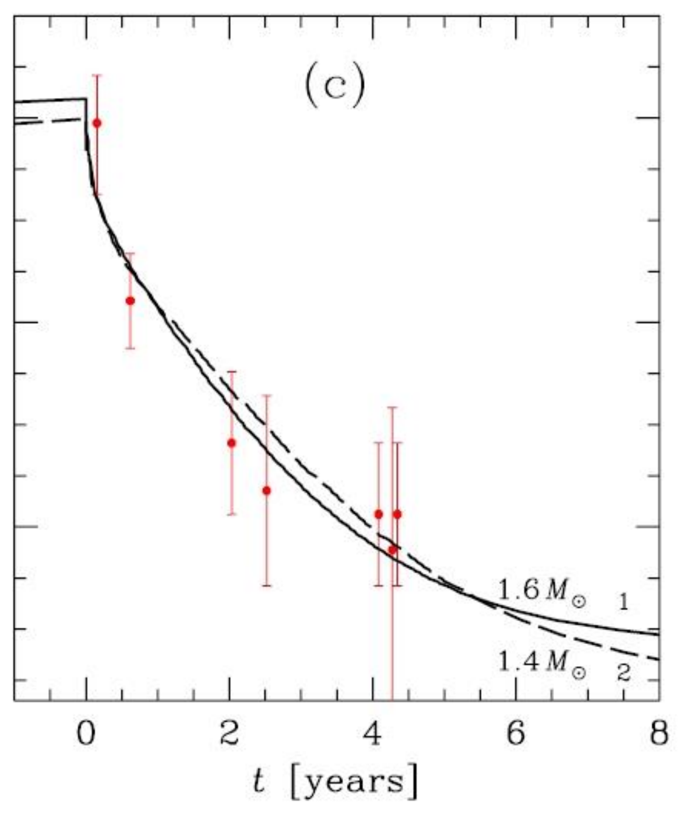
Deep crustal heating

~1.9 Mev per accreted nucleon
Crust is not in thermal equilibrium with the core.
After accretion is off the crust cools down and finally reach equilibrium with the core.

(see a more recent model in 1202.3378)

KS 1731-260

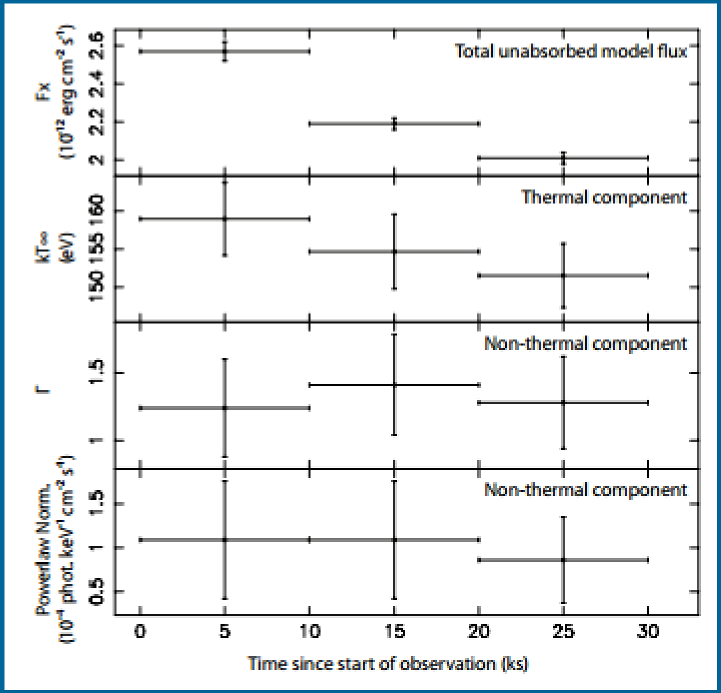




[Shternin et al. 2007]

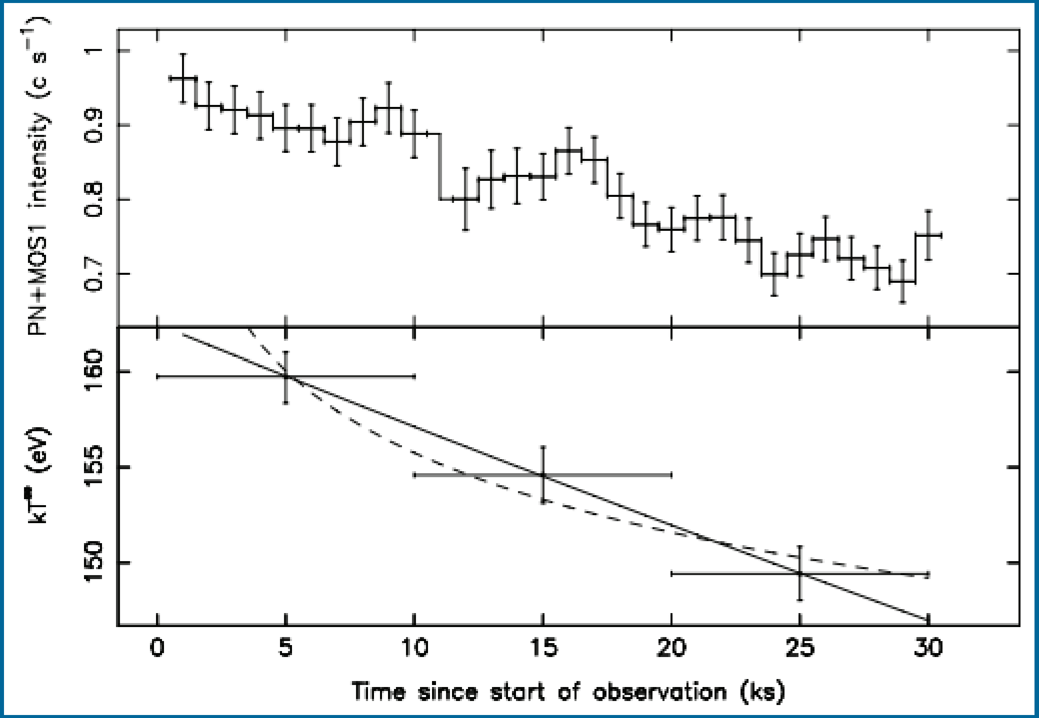
See new results and discussion in 1702.08452

Visible cooling of a NS in a binary



The authors interpret this as cooling of a layer located at a column density of y ≃ 5×10^{12} g cm^{-2} (≃50 m inside the neutron star), which is just below the ignition depth of superbursts.

XTE J1709–267



1212.1453

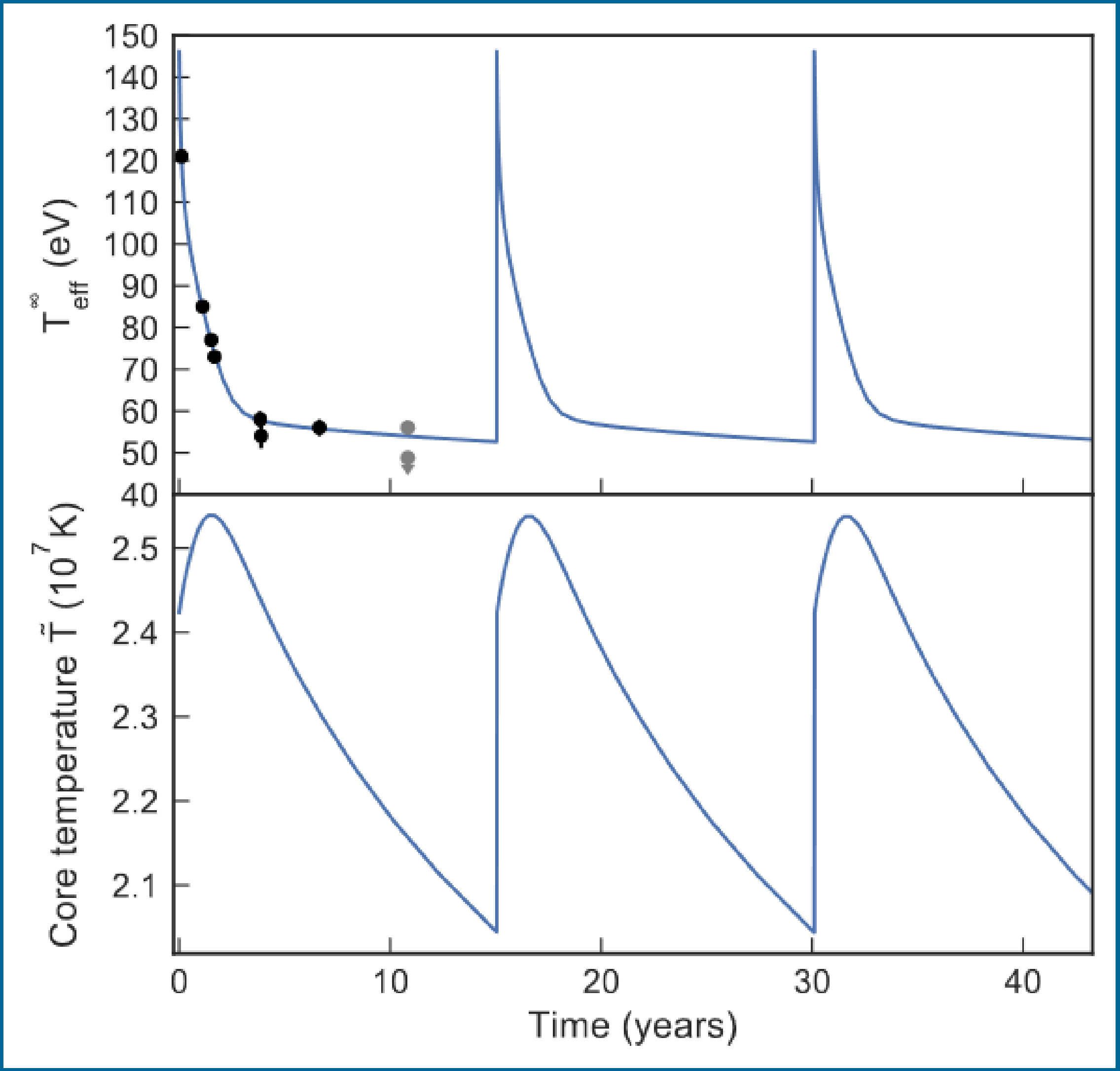
Direct Urca in a cooling NS

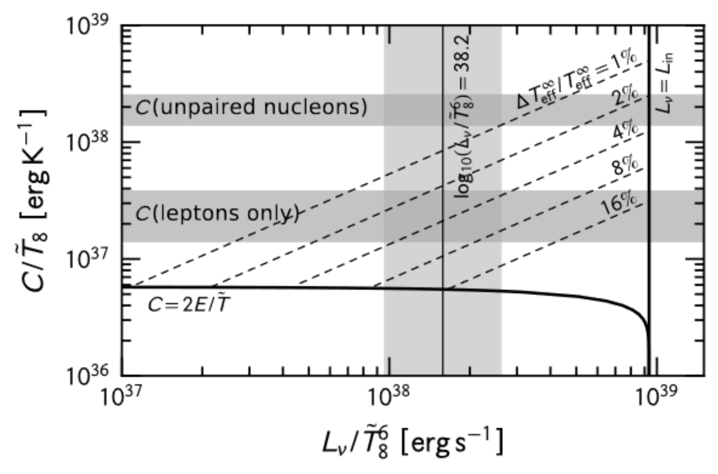
MXB 1659-29

$$2.1 \times 10^{38}\ \mathrm{erg\ s^{-1}}\ \tilde{T}_8^6$$

$$C = 10^{37}\ \mathrm{erg\ K^{-1}}\ \tilde{T}_8.$$

About 1% of the core volume available for direct URCA.





1801.00041

Cooling and crustal properties

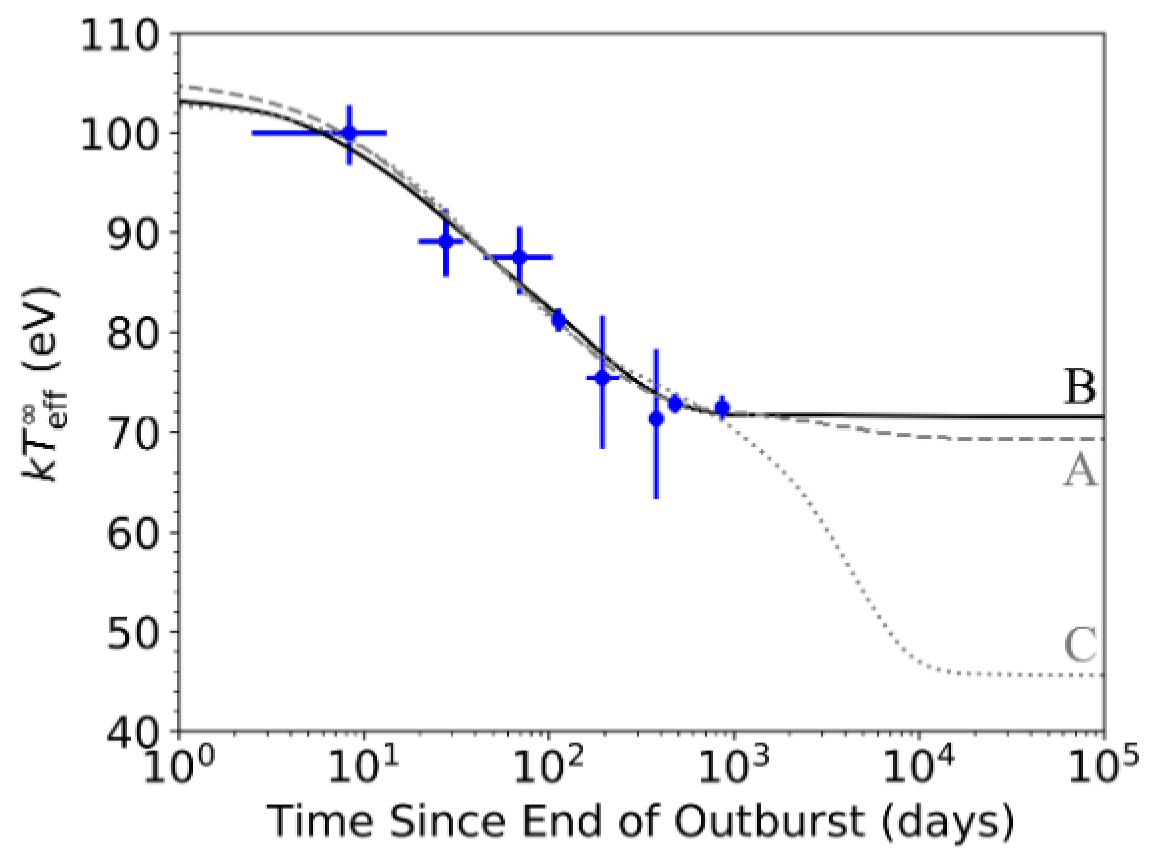
RXS J180408.9−342058
LMXB

Rapid cooling down to thermal equilibrium between the core and the crust.

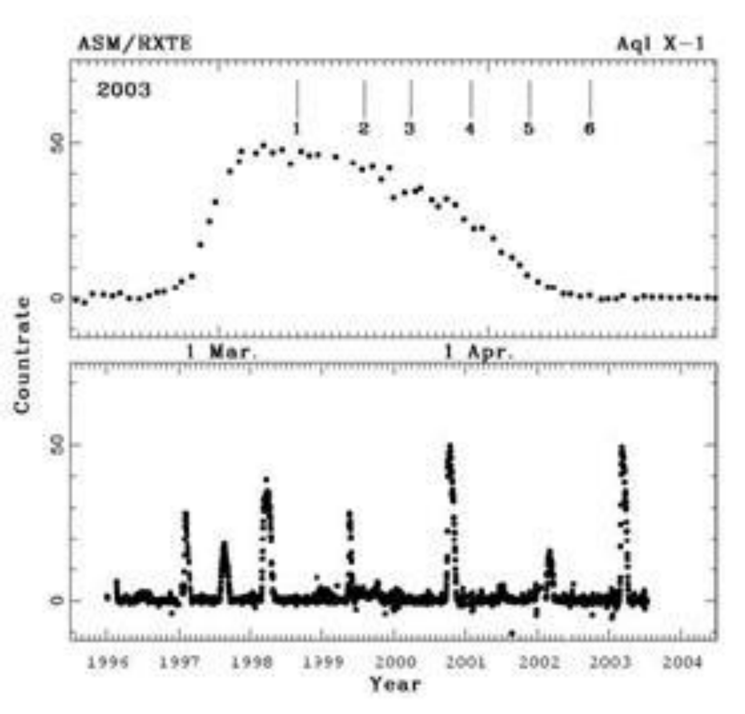
Deep crustal heating + shallow heat source.

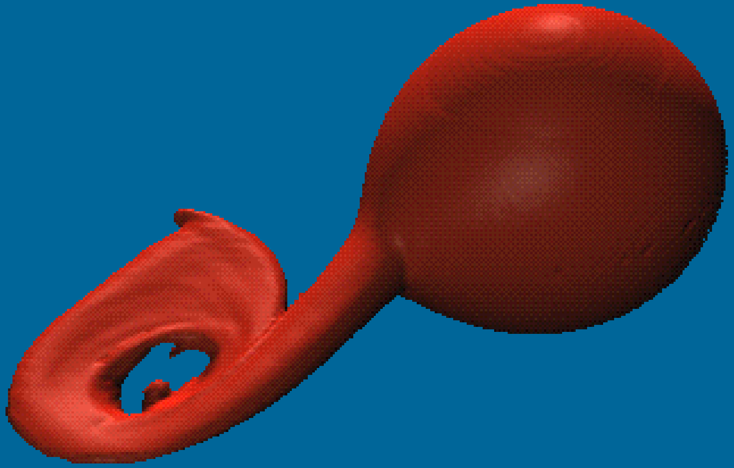
The origin of the shallow heating is unknown.

No DURCA.

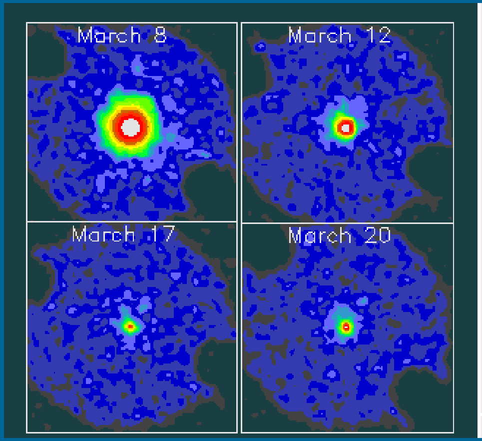


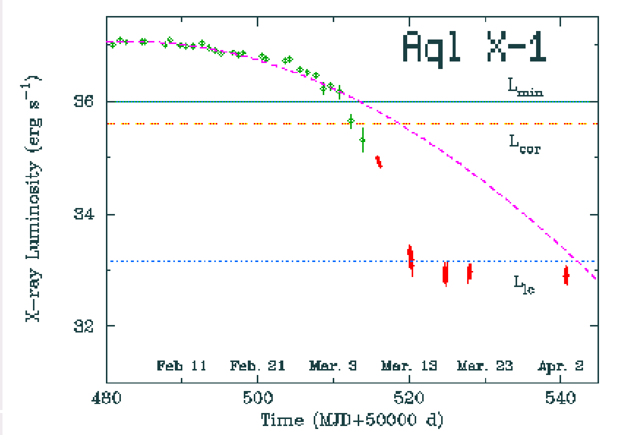
Aql X-1 transient



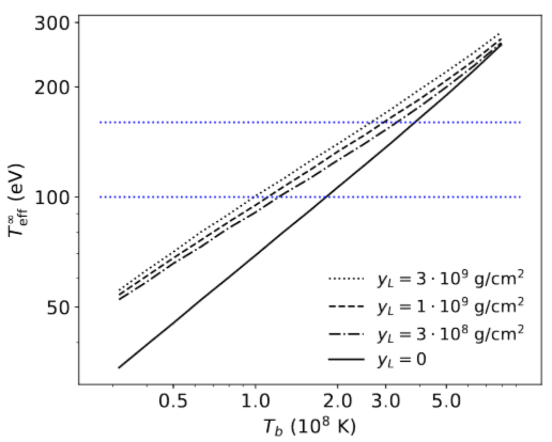


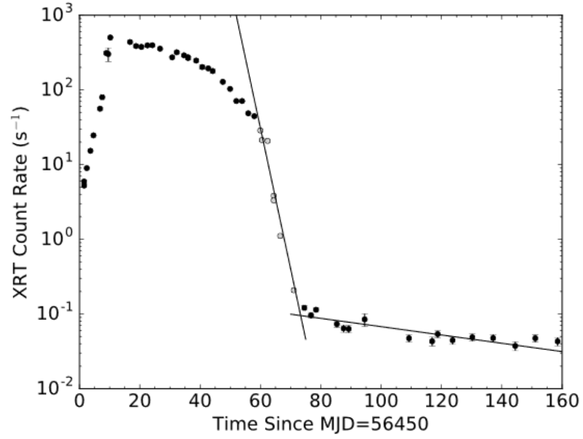
A NS with a K star.
The NS is the hottest among SXTs.

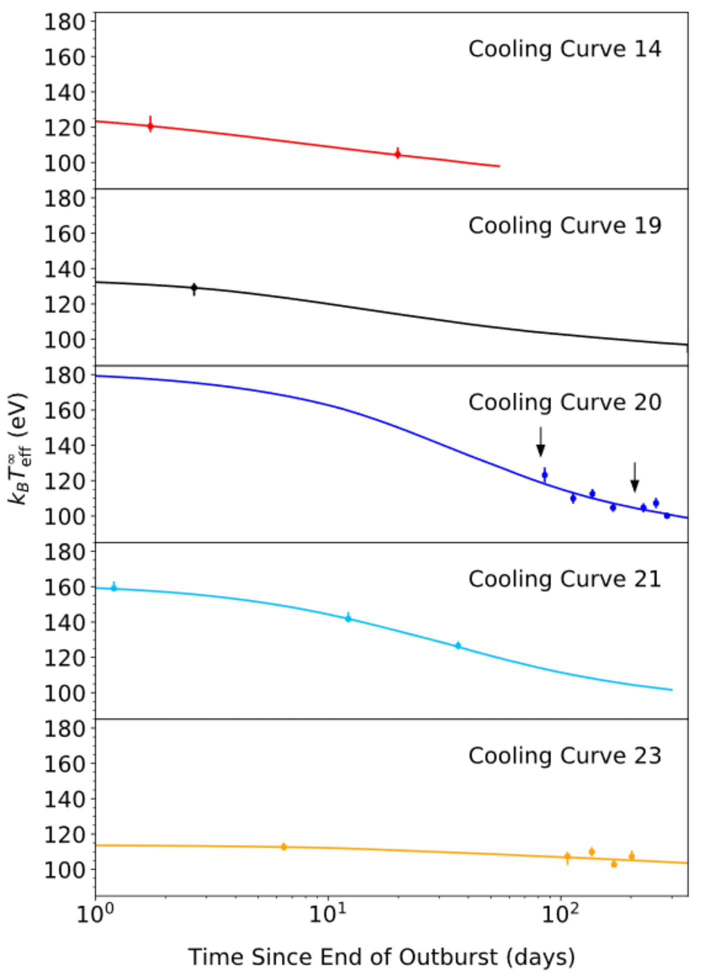


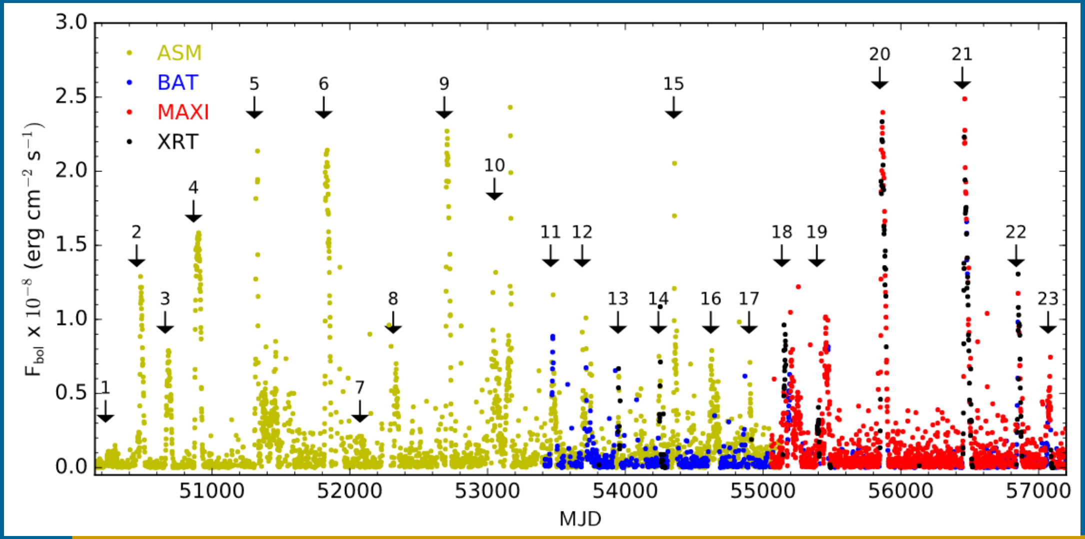


Aql X-1 modeling



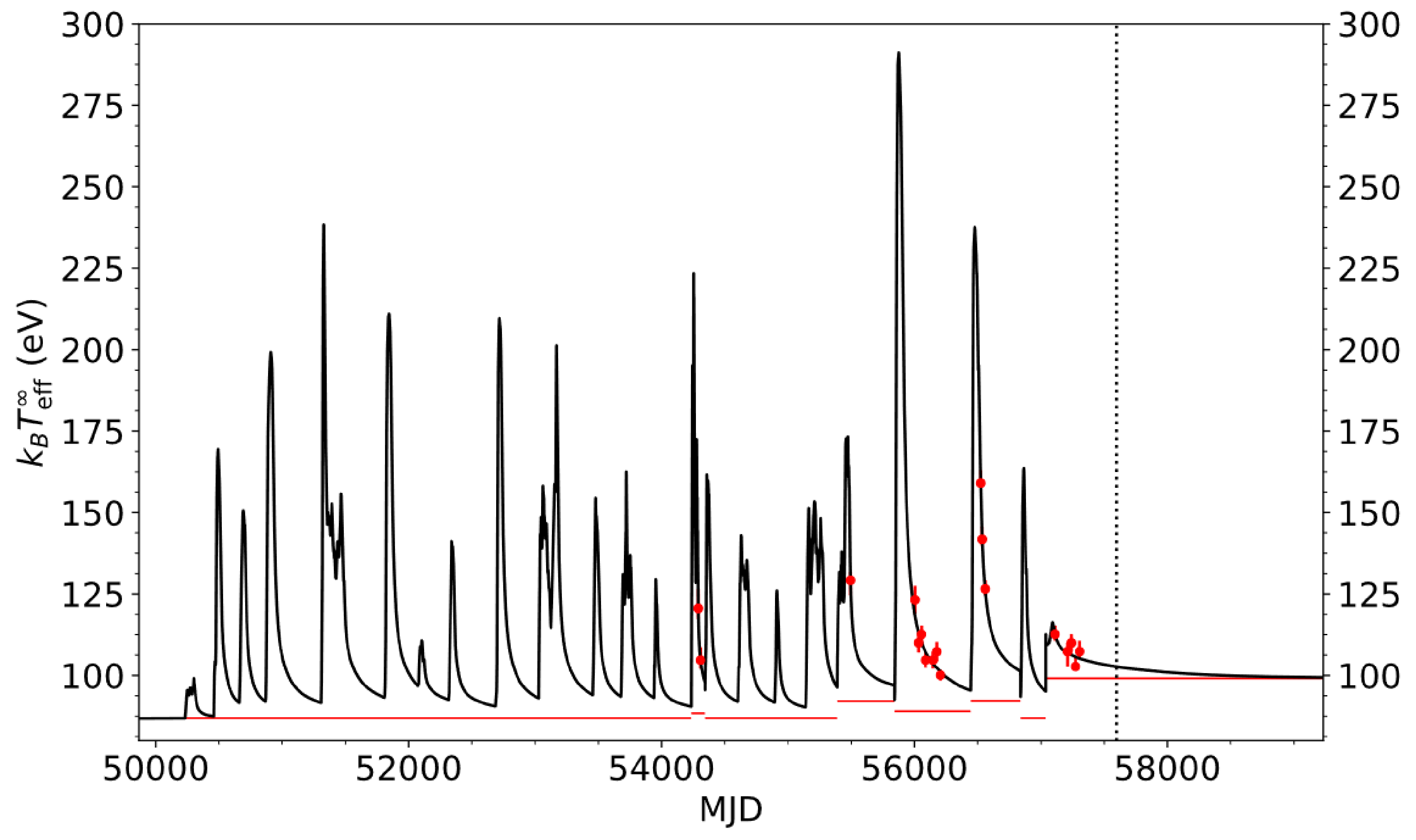






1802.06081

Fitting multiple bursts and decays



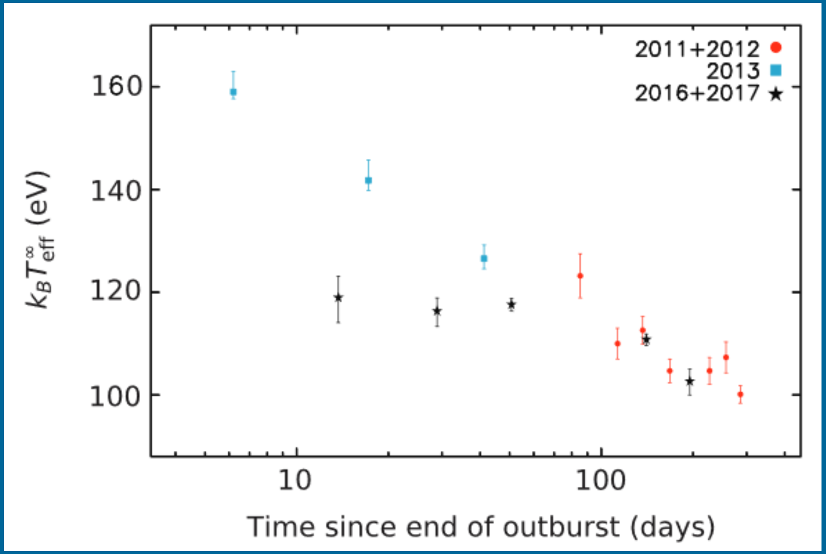
1802.06081

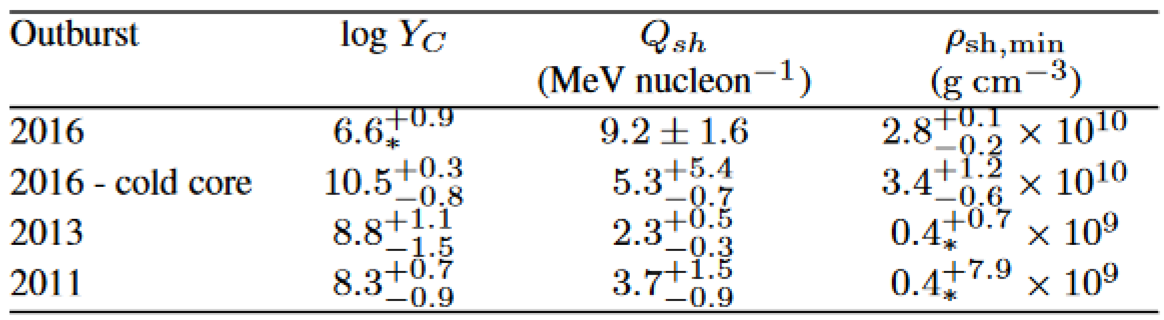
New questions from Aql X-1

Strong energy release is necessary not deep in the crust (intensive shallow heating).

$\sim 5 - 10\ \mathrm{MeV\ nucleon}^{-1}$ $\sim 10^{10}\ \mathrm{g\ cm}^{-3}$

In general, observations of different sources require different energy release and different depths, which are typically smaller than in the classical deep crustal heating ($10^{12} - 10^{13}$ g/cm^3).

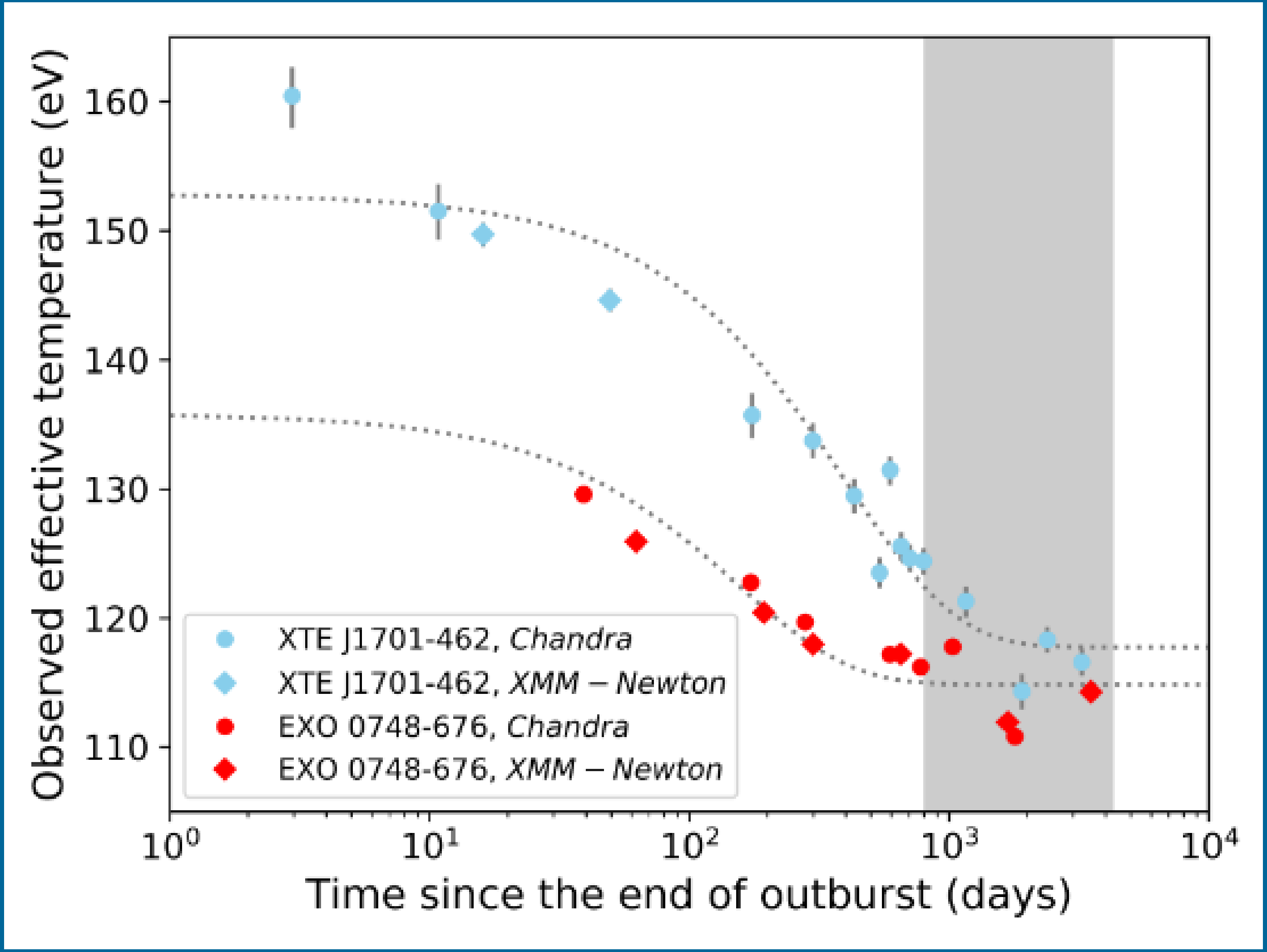




Outburst	log Y_C	Q_{sh} (MeV nucleon^{-1})	$\rho_{\mathrm{sh,min}}$ (g cm^{-3})
2016	$6.6^{+0.9}_{*}$	9.2 ± 1.6	$2.8^{+0.1}_{-0.2} \times 10^{10}$
2016 - cold core	$10.5^{+0.3}_{-0.8}$	$5.3^{+5.4}_{-0.7}$	$3.4^{+1.2}_{-0.6} \times 10^{10}$
2013	$8.8^{+1.1}_{-1.5}$	$2.3^{+0.5}_{-0.3}$	$0.4^{+0.7}_{*} \times 10^{9}$
2011	$8.3^{+0.7}_{-0.9}$	$3.7^{+1.5}_{-0.9}$	$0.4^{+7.9}_{*} \times 10^{9}$

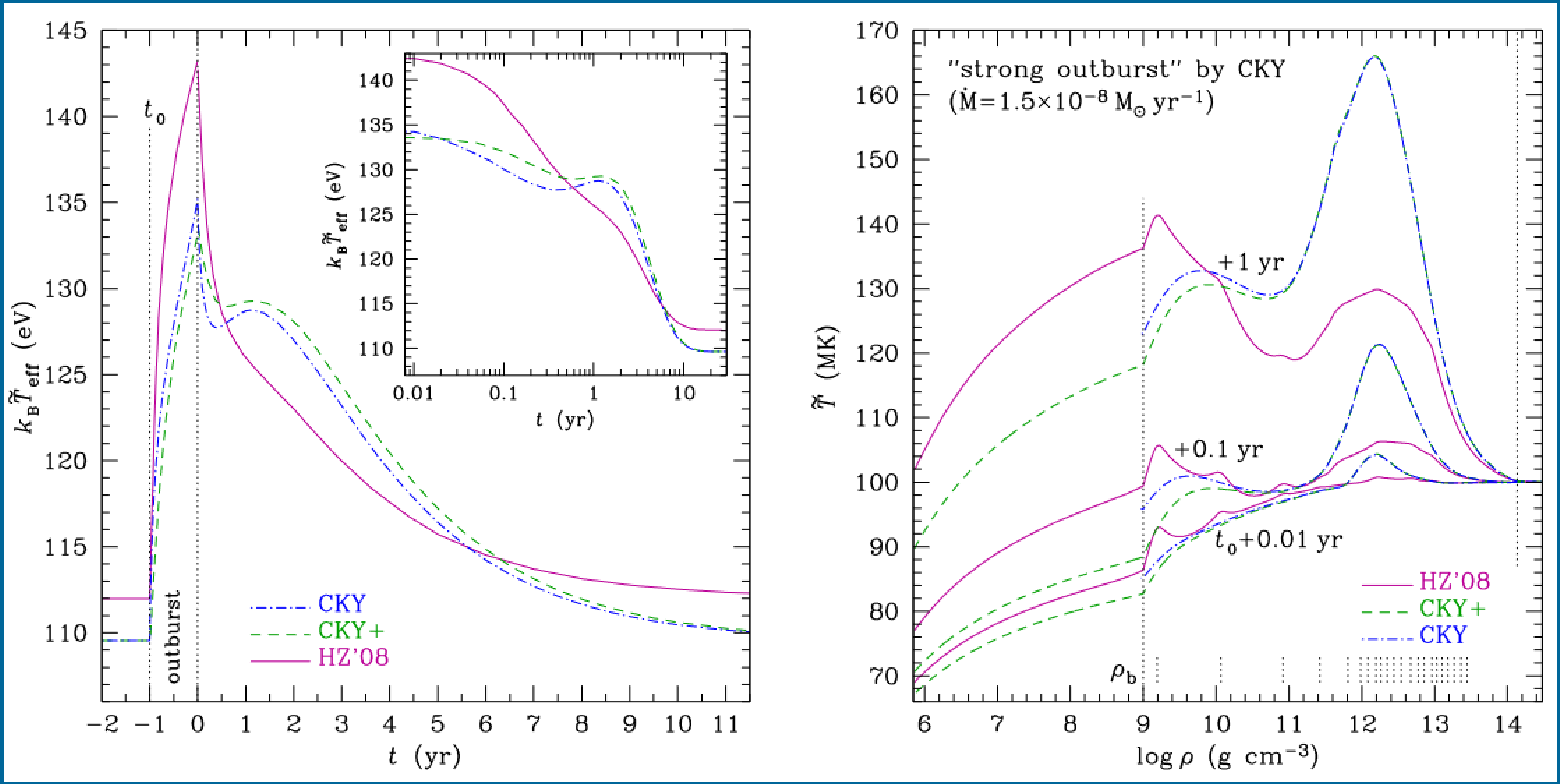
1907.07205

Unexpected temperature rise



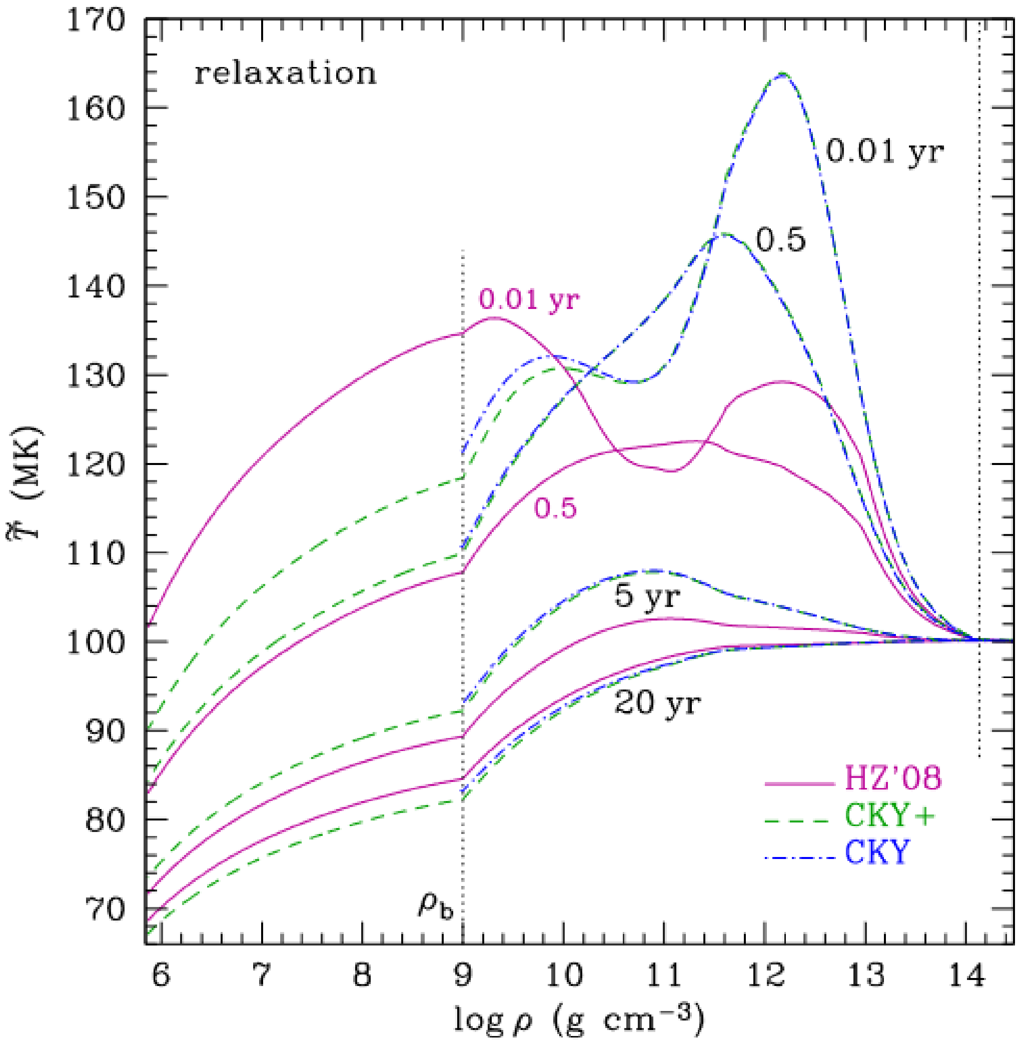
2005.12014

Detailed models



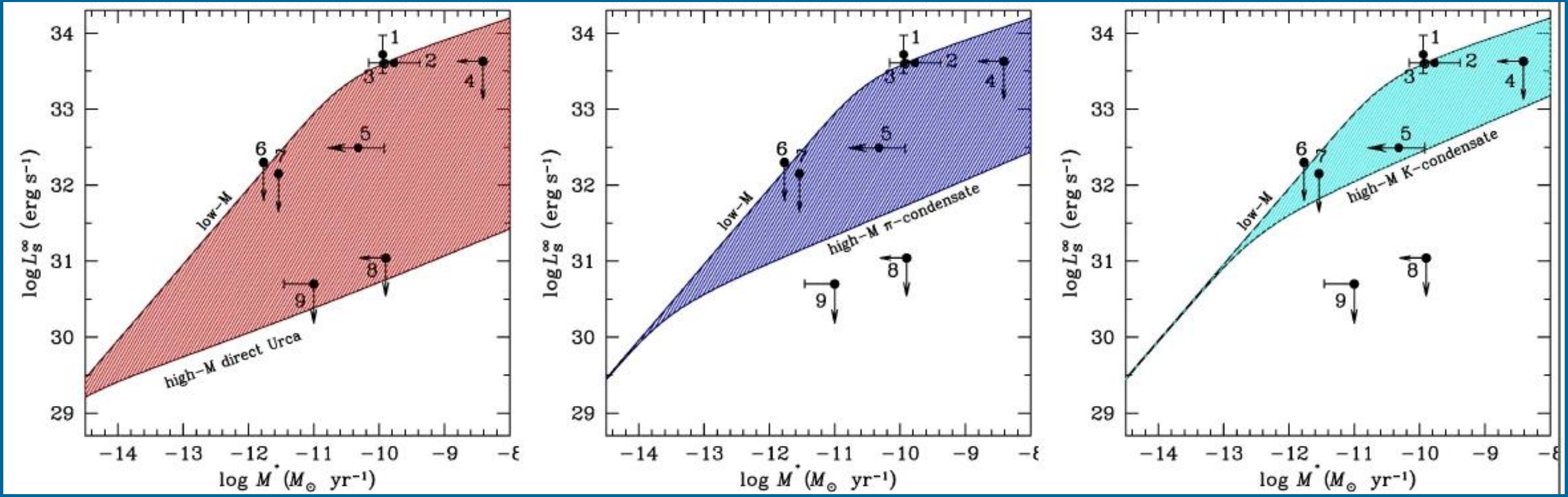
2011.10490

Quiescent state: temperature evolution



2011.10490

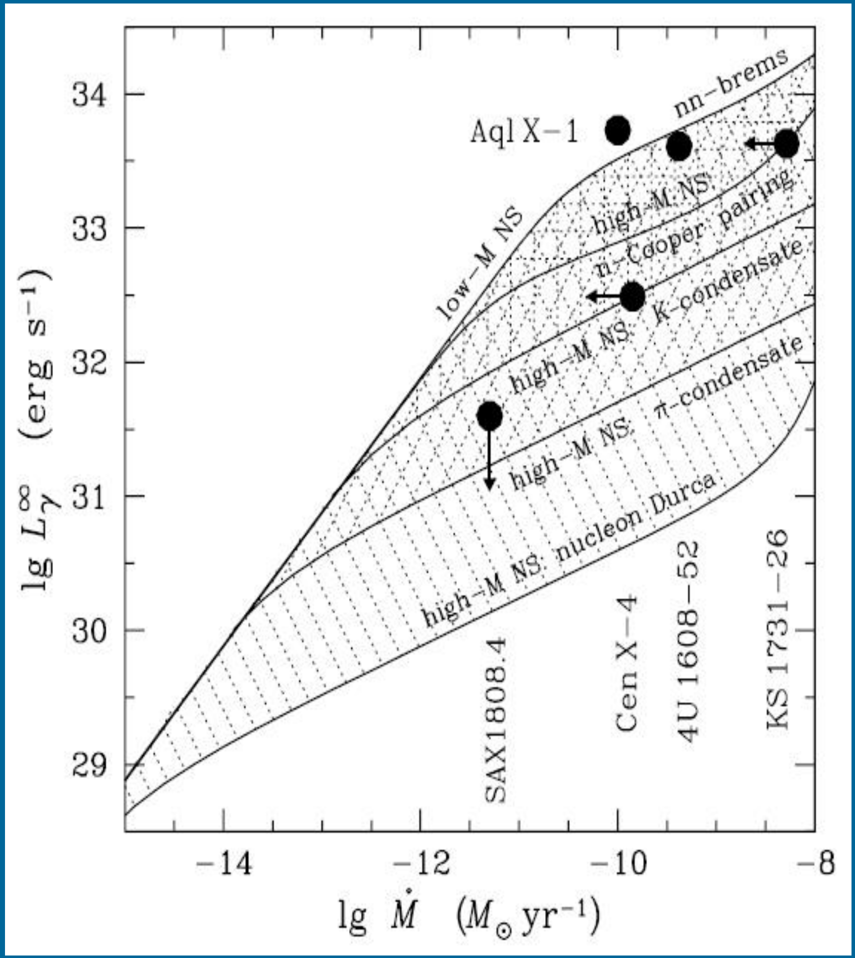
Testing models with SXT

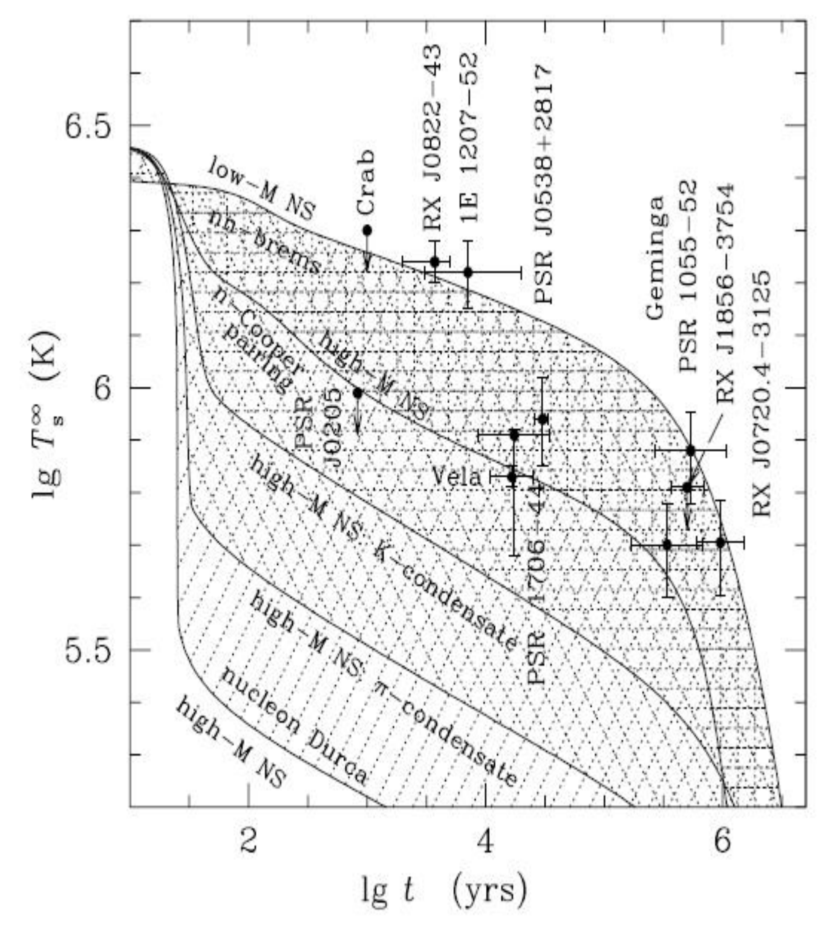


SXTs can be very important in confronting theoretical cooling models with data.

[from a presentation by Haensel, figures by Yakovlev and Levenfish]

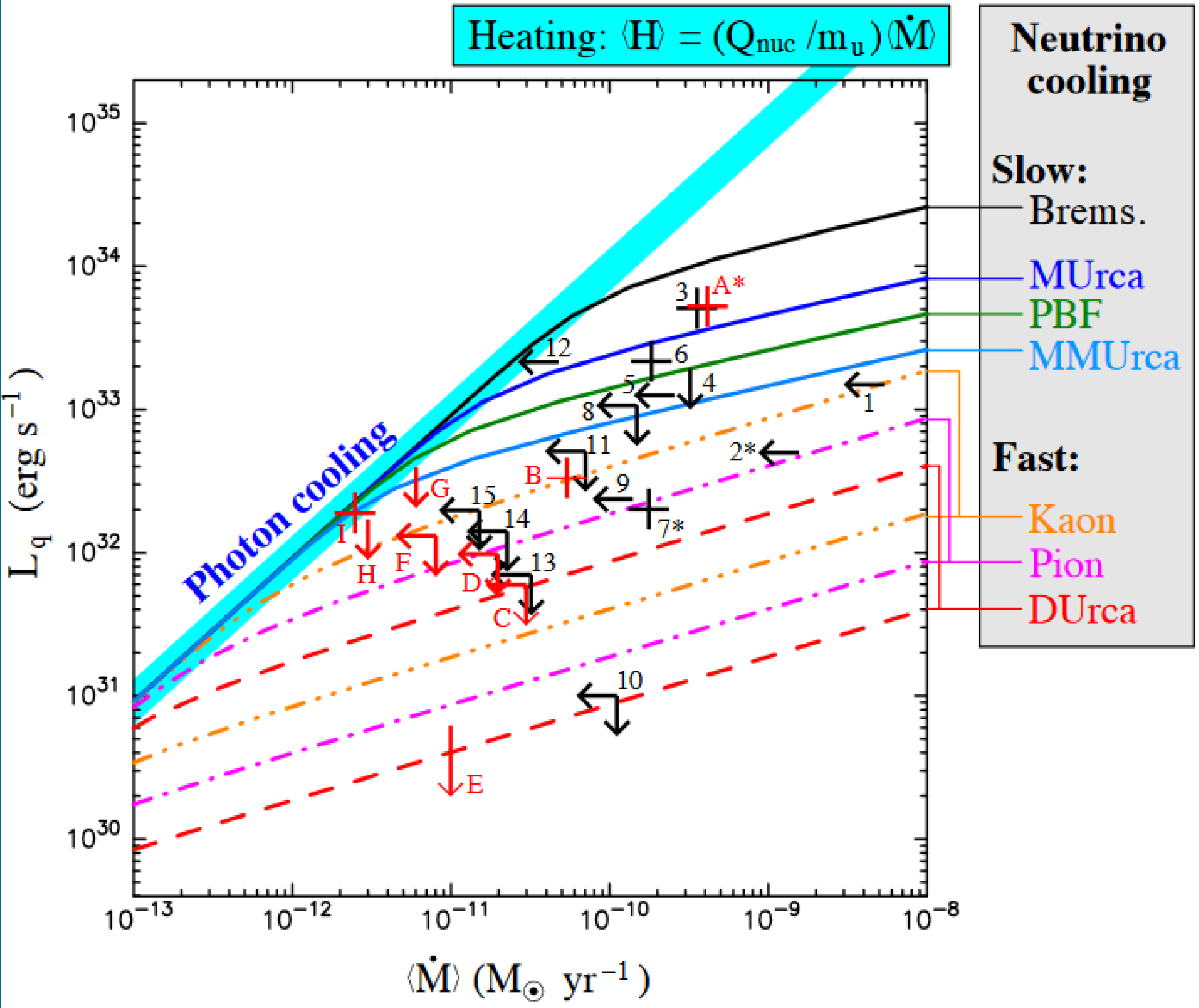
Theory vs. Observations: SXT and isolated cooling NSs





[Yakovlev et al. astro-ph/0501653]

Systems with deep crustal heating



Accreting ms X–ray pulsars

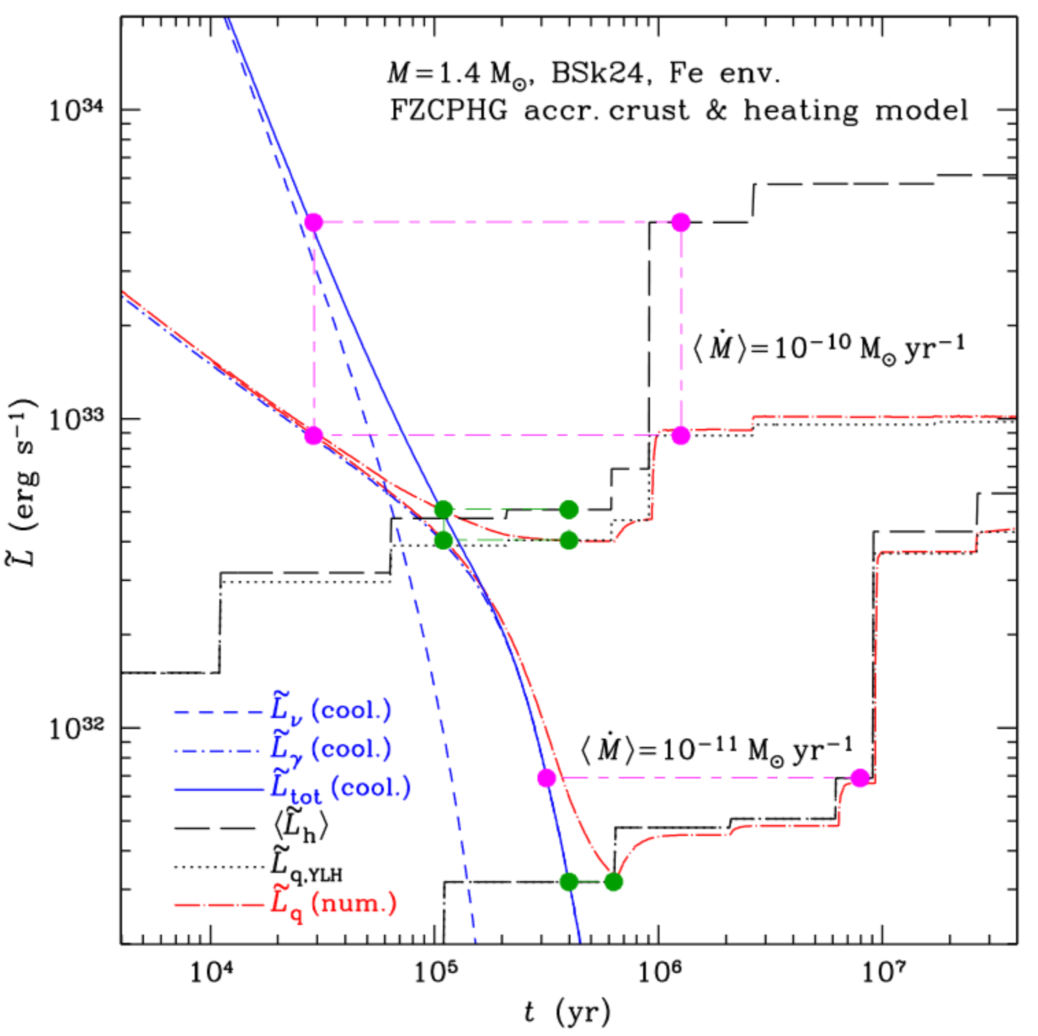
A: Aql X–1*
B: SAX J1748.9–2021
C: NGC 6440 X–2
D: XTE J0929–314
E: SAX J1808.4–3658
F: XTE J1807–294
G: XTE J1751–305
H: XTE J1814–338
I: IGR J00291+5934

Non pulsating NS LMXBs

1: 4U 2129+47
2: KS 1731–260*
3: 4U 1608–522
4: EXO 1745–248
5: 1M 1716–315
6: XTE J1709–267
7: MXB 1659–298*
8: X 1732–304
9: Cen X–4
10: 1H 1905+000
11: 2S 1803–245
12: 4U 1730–22
13: EXO 1747–214
14: XTE J2123–058
15: SAX J1810.8–26

1709.07034

How to determine quiescent luminosity

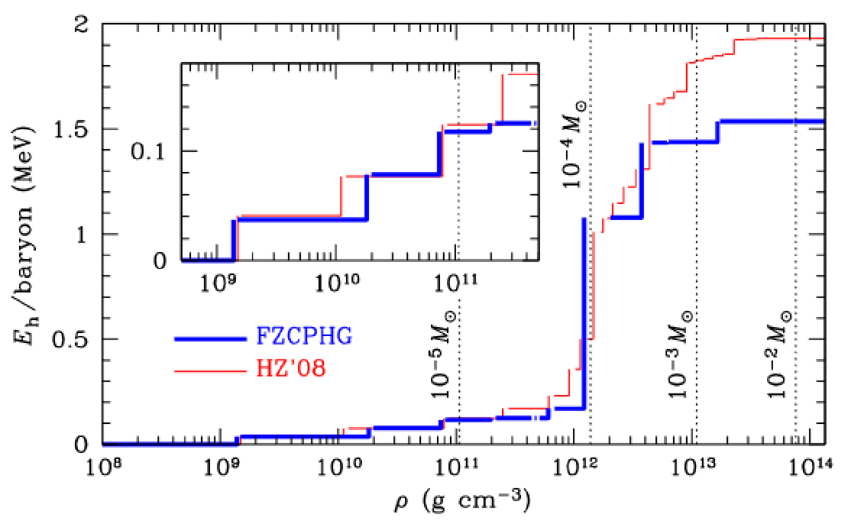


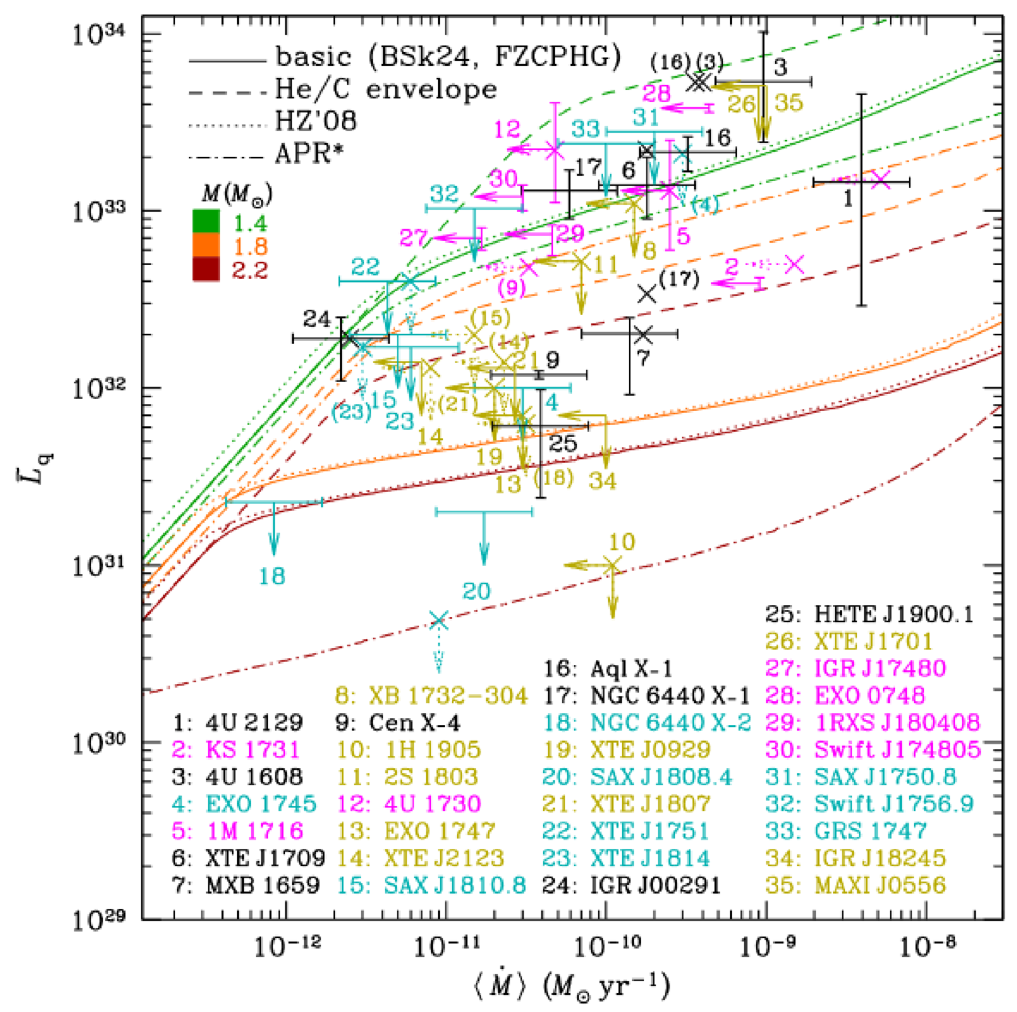
Only reactions in the accreted crust are taken into account.

Quiescent luminosity (black dotted) is equal to the photon luminosity of the cooling curve (blue dot-dashed) at the moment when total cooling luminosity (blue solid) is equal to the heat energy release (black dashed).

1907.08299

Deep crustal heating





1907.08299

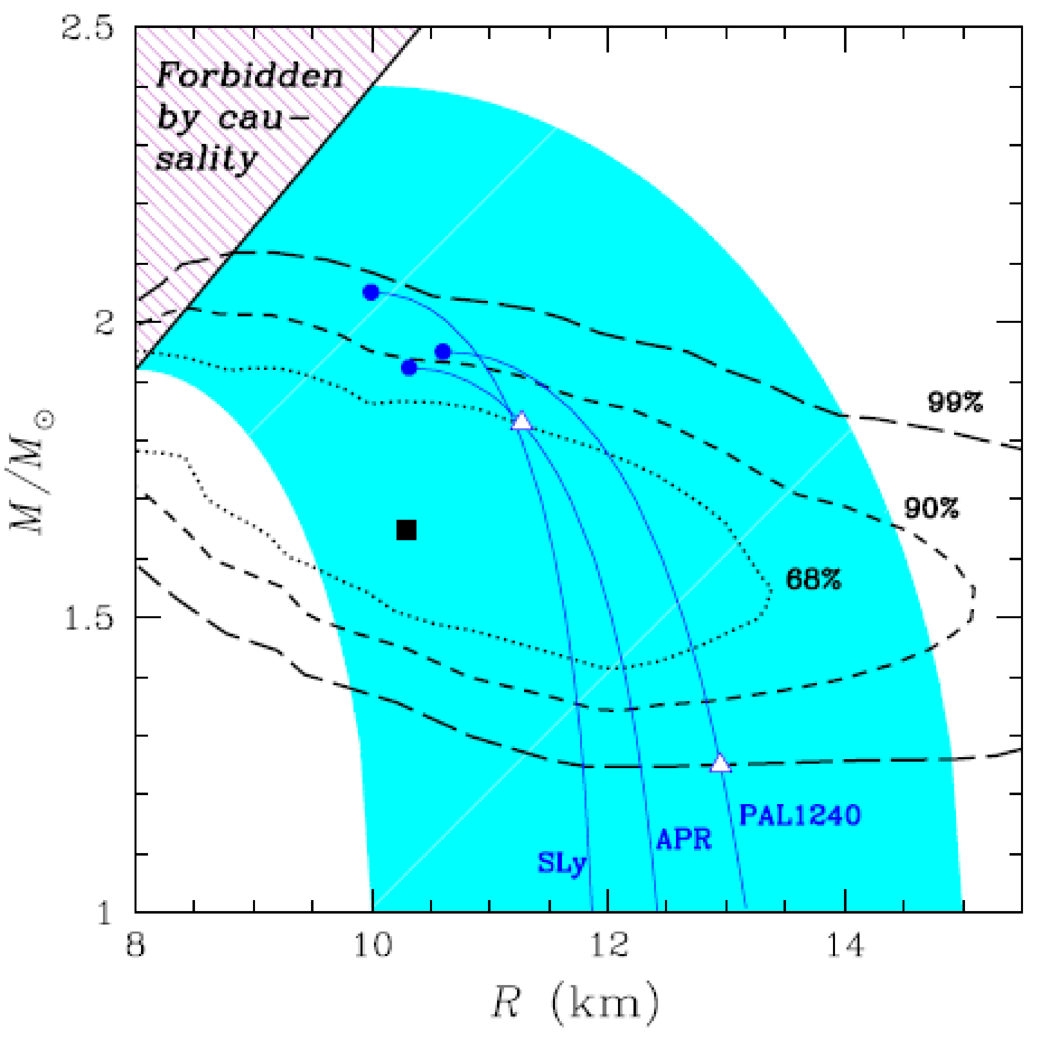
Conclusions

- NSs are born hot, and then cool down at first due to neutrino emission, and after – due to photon emission
- Observations of cooling provide important information about processes at high density at the NS interiors
- Two types of objects are studied:
  - isolated cooling NSs
  - NSs in soft X-ray transients

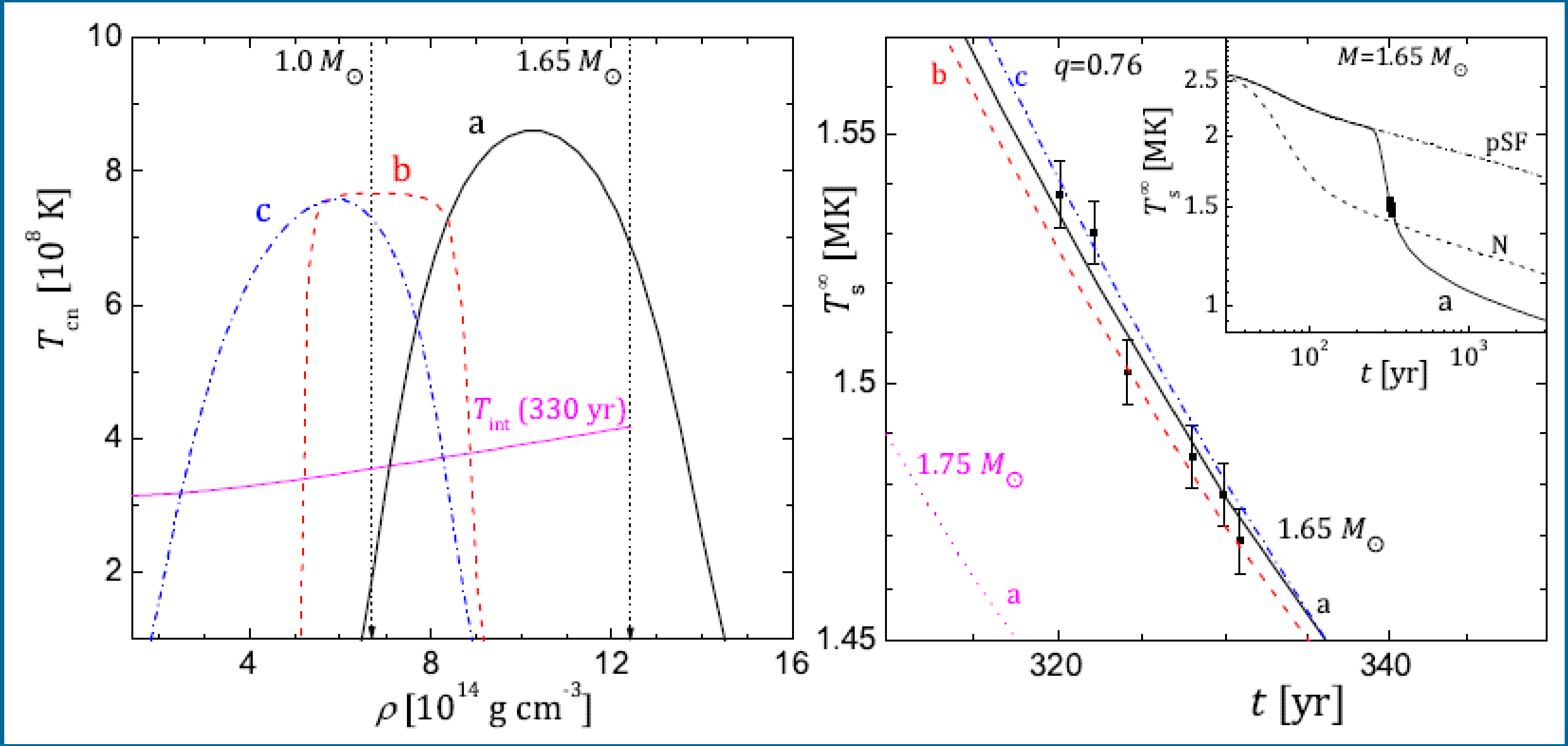
Papers to read

- Or astro-ph/0403657
  Or astro-ph/0508056
  Or astro-ph/0402143
  Or 1507.06186
  Or 2006.15004 ←
- arXiv:astro-ph/9906456 УФН 1999
- 1709.07034 – about cooling of NSs in binaries
- Cooling of magnetars (and general basic of cooling and field evolution) 1911.03095
- Deep crustal heating 1907.08299 ←

M-R from spectral fit

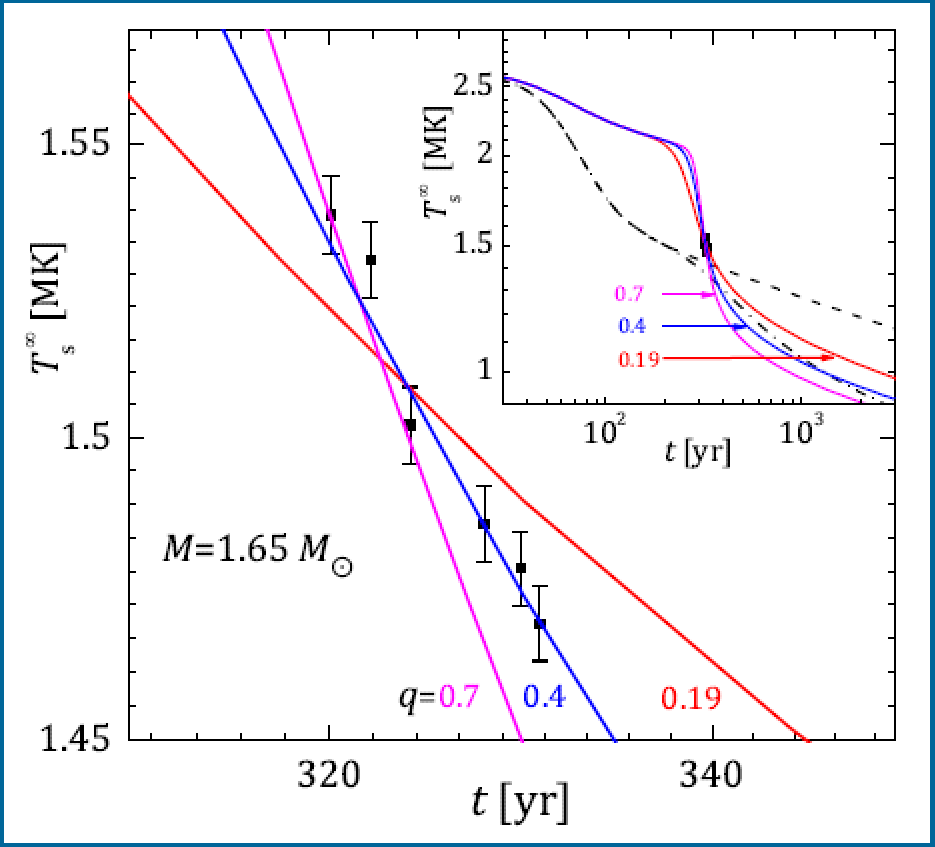


1010.1154


10
8
6
4
2
T_{cn} [10^8 K]
4
8
12
16
ρ [10^{14} g cm^{-3}]
1.0 $M_\odot$
1.65 $M_\odot$
a
b
c
T_{int} (330 yr)
1.55
1.5
1.45
T_s^∞ [MK]
320
340
t [yr]
b
c
q=0.76
1.75 $M_\odot$
a
1.65 $M_\odot$
a
M=1.65 $M_\odot$
2.5
2
1.5
1
T_s^∞ [MK]
10^2
10^3
t [yr]
pSF
N
a

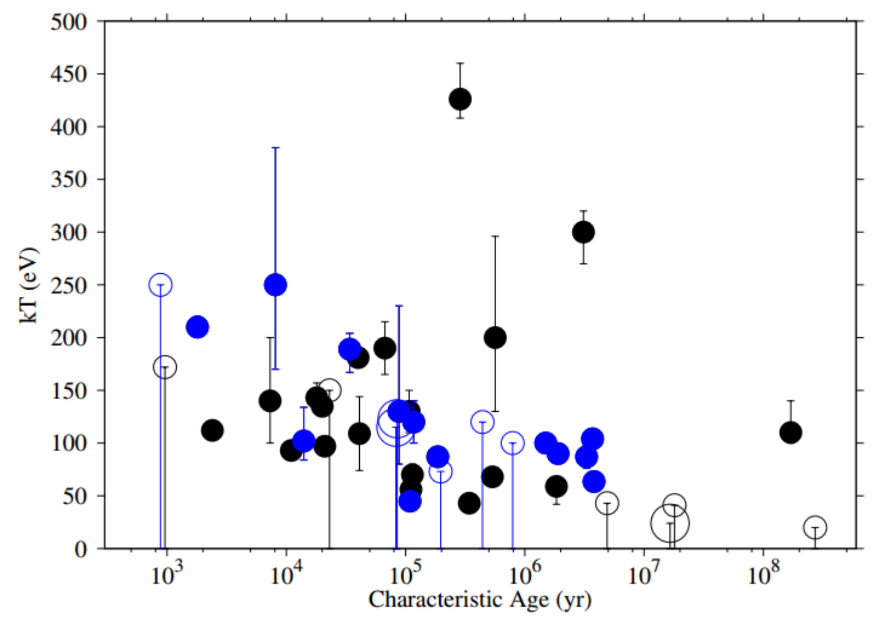
1012.0045

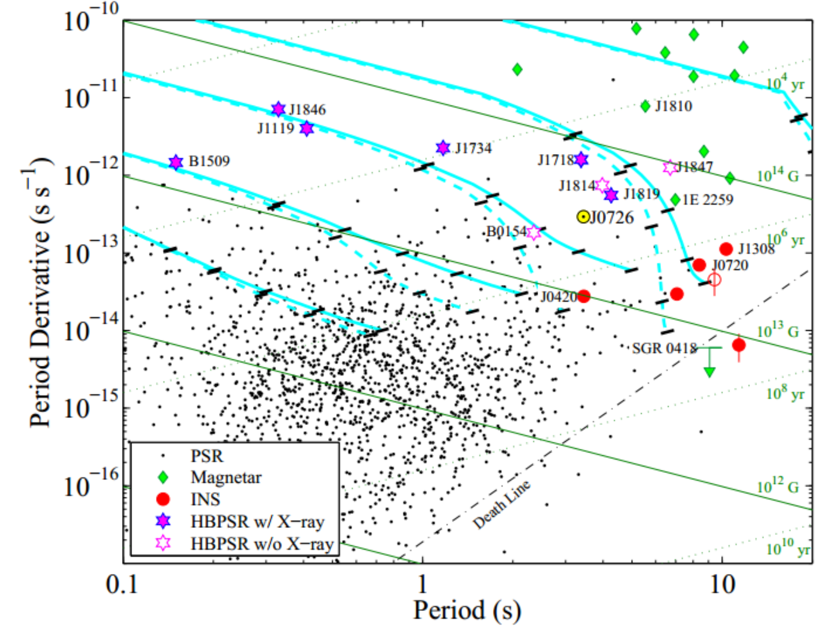
Suppression in the axial-vector channel



1012.0045

Cooling and grand unification for NSs



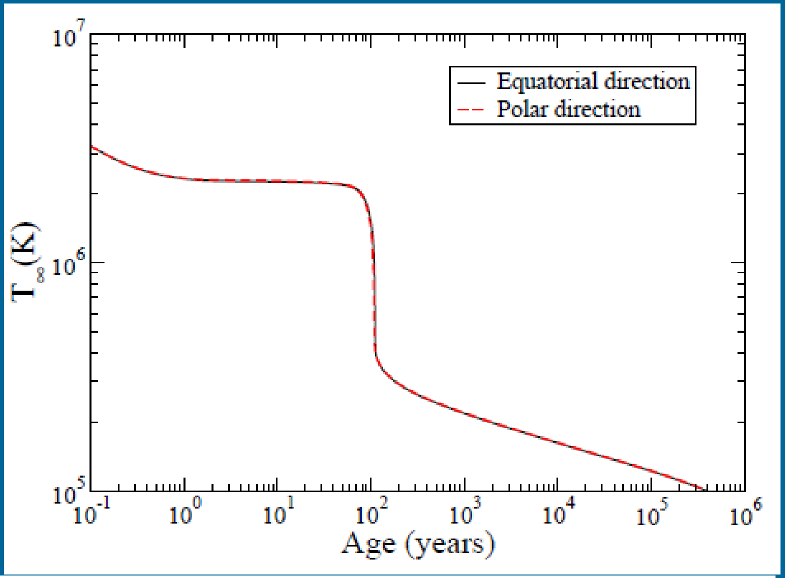


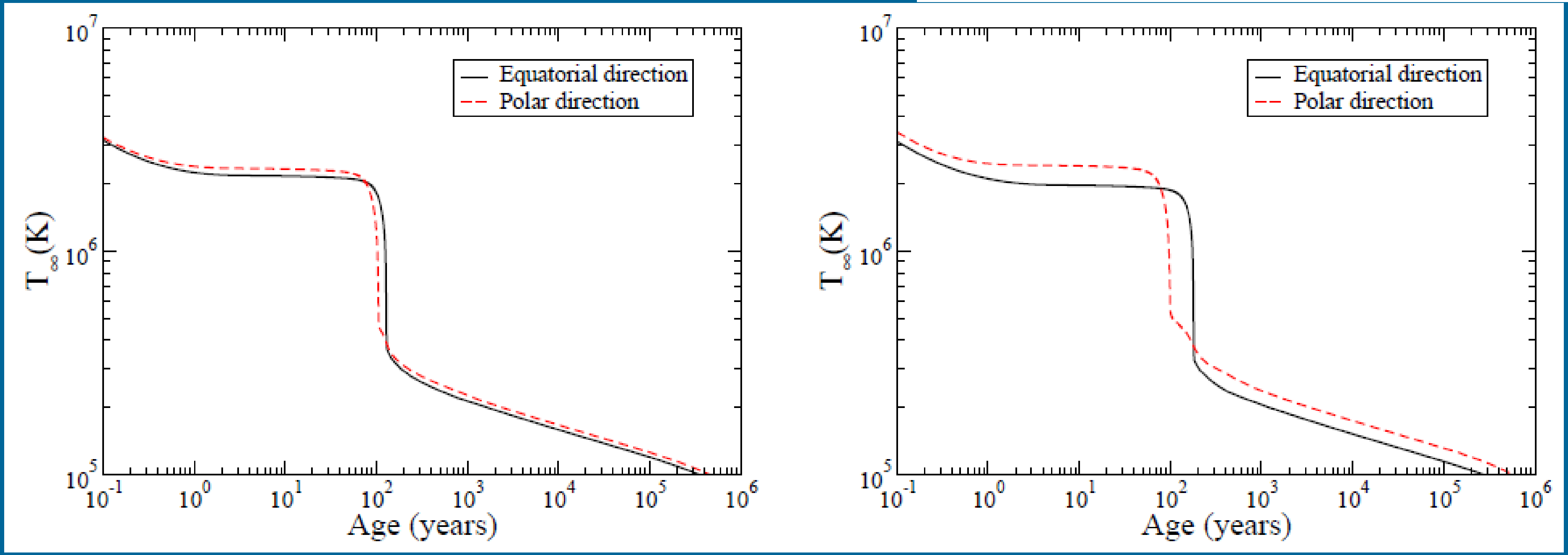
1301.2814

1111.2877

One study shows that highly magnetized NSs can be not hotter than NSs with standard magnetic fields.
Another study demonstrates that some young PSRs with relatively large field are hot, similar to the M7.

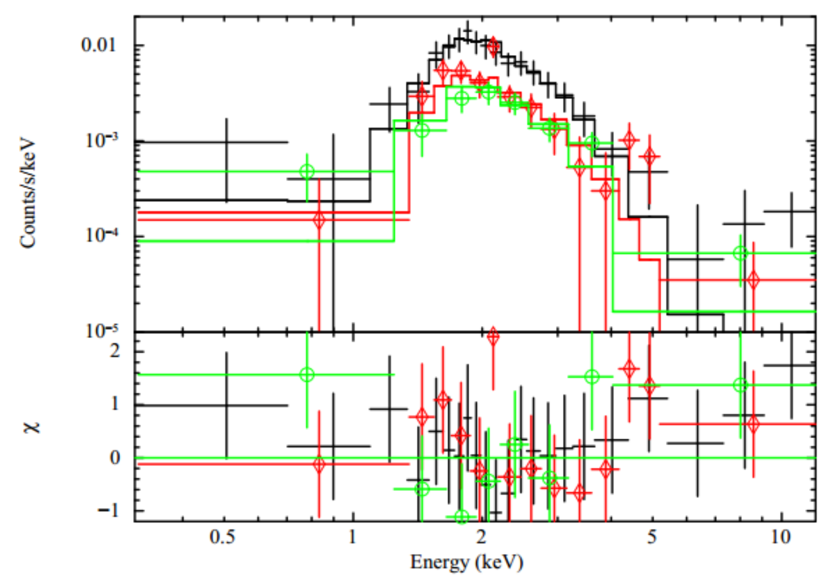
Влияние вращения





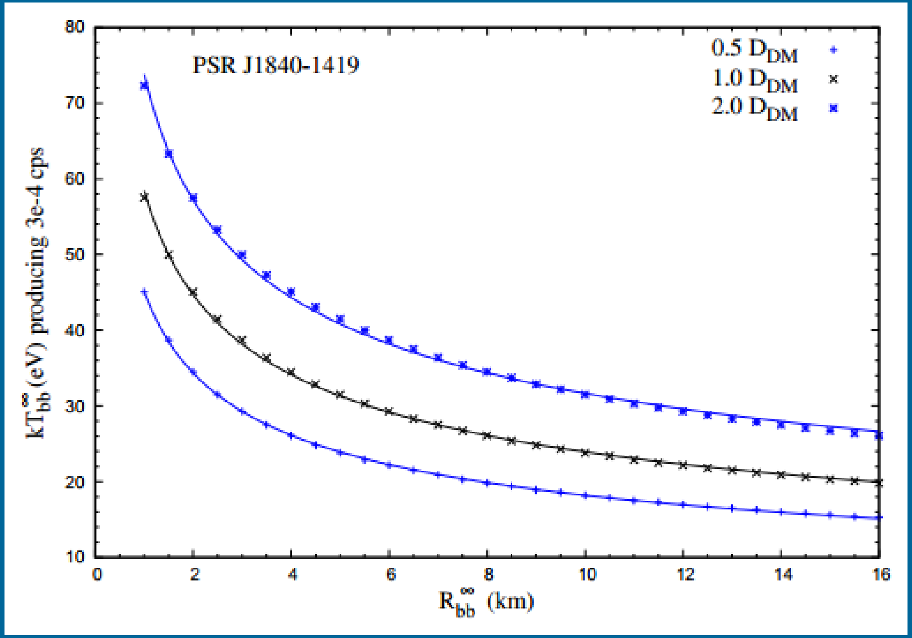
1307.7692

Records

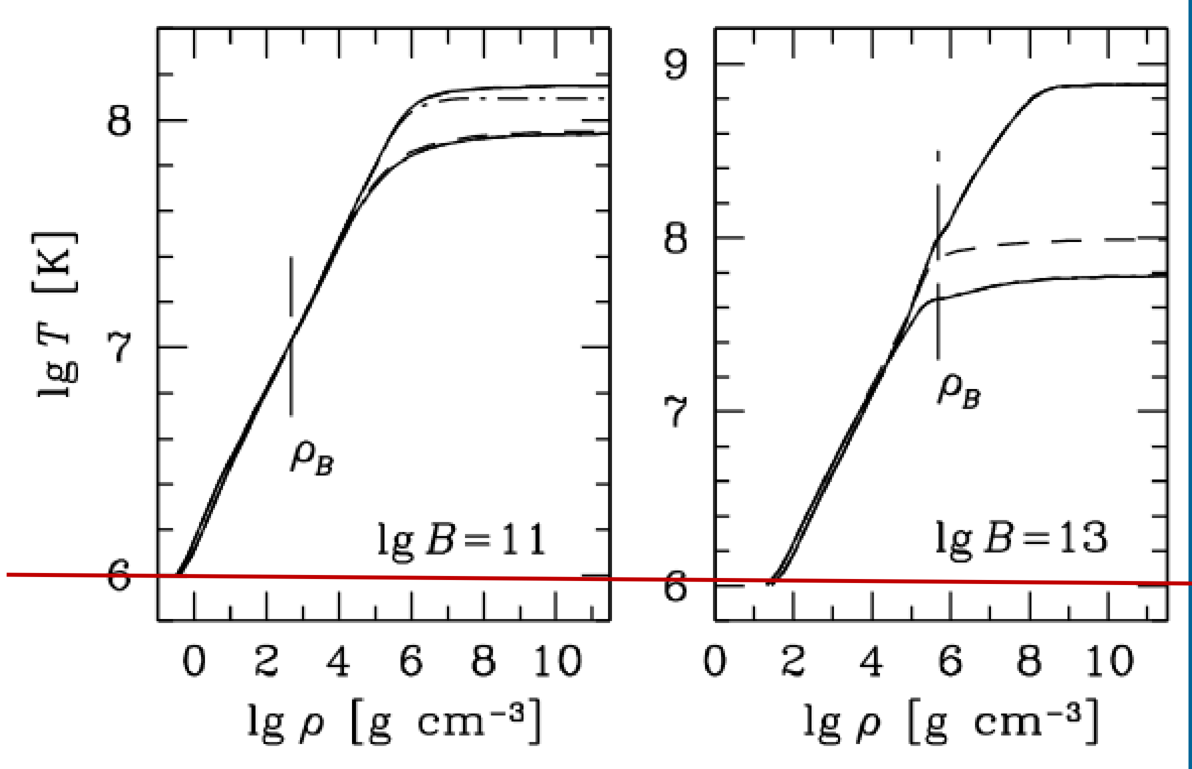


The hottest (in a binary, crustal heating)
SAX J1750.8−2900. T~150 eV.
1202.1531

The coldest. Isolated pulsar. T<30 eV
PSR J18401419
1301.2814



Temperature profile in an envelope



Iron envelope.

Solid line are
new calculations
for poles and equator
(i.e., along and across
magnetic field lines).

Surface
temperature
fixed at 10^6 K.

$$T_s^{(0)} \approx 10^6 \, g_{14}^{1/4} \left[(7\zeta)^{2.25} + (\zeta/3)^{1.25} \right]^{1/4} \quad \mathrm{K}$$

$$\zeta \quad \equiv \quad T_{\mathrm{int},9} \; - \; 0.001 \, g_{14}^{1/4} \sqrt{7 \, T_{\mathrm{int},9}},$$

astro-ph/0105261